

First measurements from charmless B decays at Belle II

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We report on the first measurements of branching fractions, CP -violating charge-asymmetries, and polarizations in various charmless B decays at Belle II. We use a data sample of electron-positron collisions collected in 2019–2020 at the $\Upsilon(4S)$ resonance from the SuperKEKB collider. The data sample corresponds to an integrated luminosity of 34.6 fb^{-1} . All results are consistent with known values. In addition, these results provide extensive validations of the detector performance and analysis strategies.

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1. Introduction

Charmless B decays are important to search for non-Standard-Model (non-SM) physics in the flavor sector. Many decay channels are governed by ‘penguin’ amplitudes, which are sensitive to non-SM contributions within the loop. Studying them in detail is an important goal of the Belle II experiment [1]. With the largest sample of e^+e^- collisions anticipated in the next decade, Belle II is expected to improve significantly the measurements associated with charmless B decay, such as the determination of the CKM phase α/ϕ_2 [1, 2], the precision test of the $K\pi$ isospin sum rule [1, 3], and the study of local CP-violating asymmetries in the phase space of three-body B decays [1]. In addition, the measurement of decay-time-dependent CP violation in the penguin-dominated $B^0 \rightarrow \phi K^0$ mode, compared with corresponding results from $B^0 \rightarrow J/\psi K^0$ decays, will offer a probe of non-SM physics [1]. Measurements of the longitudinal-polarization fractions (f_L) of decays of B mesons into pairs of vector mesons also probe non-SM dynamics. Previous measurements of f_L with $B^0 \rightarrow J/\psi K^0$ decays showed a sizable contribution from transverse polarization, while most predictions expect the longitudinal component to dominate [4]. More precise f_L measurements may shed light on the issue.

SuperKEKB [5] is an asymmetric e^+e^+ collider, which started collision operations with the Belle II detector [6] in March 2019. We use a data sample of 34.6 fb^{-1} , which was collected at the $\Upsilon(4S)$ resonance up to May 2020. This report presents the measurements of branching fractions (\mathcal{B}), CP-violating charge-asymmetries (\mathcal{A}_{CP}), and f_L based on the following B decays reconstructed in Belle II data: $B^0 \rightarrow K^+\pi^-$, $B^0 \rightarrow \pi^+\pi^-$, $B^+ \rightarrow K^+\pi^0$, $B^+ \rightarrow \pi^+\pi^0$, $B^+ \rightarrow K^0\pi^+$, $B^0 \rightarrow K^0\pi^0$, $B^+ \rightarrow K^+K^-K^+$, $B^+ \rightarrow K^+\pi^-\pi^+$, $B^0 \rightarrow \phi K^0$, $B^+ \rightarrow \phi K^+$, $B^0 \rightarrow \phi K^{*0}$, and $B^{*+} \rightarrow \phi K^{*+}$ [7, 8].

The B reconstruction, event-selection criteria, and background suppression strategy are studied with various simulated signal and background samples. Charged-particle trajectories (tracks) are identified with inner vertex detectors and a central drift chamber with requirements on the displacement from the interaction point to reduce beam-background-induced tracks. The identification of charged particles uses the information from two particle-identification (PID) devices, a time-of-propagation counter in the barrel region and a proximity-focusing aerogel ring-image Cherenkov counter in the forward endcap region. Decays of π^0 candidates are reconstructed by combining two isolated clusters in the electromagnetic calorimeter, with requirements on the helicity angle and the results of a kinematic-fit constrained to the π^0 mass. Decays of K_S^0 candidates are reconstructed from two opposite-charge pion candidates consistent with arising from a common vertex, with additional requirements on their kinematic and topological variables, e.g., momentum, flight distance, distance between pion trajectories, to further reduce the combinatorial background. Decays of ϕ candidates are reconstructed from two opposite-charge kaon candidates. Decays of K^{*0} candidates are reconstructed from one K^+ and one π^- , and decays of K^{*+} candidates are reconstructed from one K_S^0 and one π^+ . In three-body decays, we suppress the relevant peaking backgrounds from charmed or charmonium intermediate states by excluding the corresponding two-body mass ranges.

We use the following two variables to distinguish the signal B events from other backgrounds: the energy difference $\Delta E \equiv E_B^* - \sqrt{s}/2$ between the reconstructed B candidate and half of the collision energy in the $\Upsilon(4S)$ frame, and the beam-energy-constrained mass $M_{\text{bc}} \equiv \sqrt{s/(4c^2) - (p_B^*/c)^2}$, where \sqrt{s} is the collision energy, and E_B^* and p_B^* are the energy and momentum of reconstructed B candidates in the $\Upsilon(4S)$ frame.

2. Continuum background suppression

One of the main challenges of reconstructing charmless B decays is the large combinatorial background from $e^+e^- \rightarrow q\bar{q}$ ($q = u, d, s, c$) processes. Signal rates 10^5 times smaller than those to produce continuum background and the lack of distinctive final-state features (leptons or intermediate resonances) make the reconstruction of signal hard. Therefore, to discriminate between signal and continuum background, a binary boosted decision-tree (BDT) classifier is used which makes use of a non-linear combination of more than 30 variables. The input variables to the BDT include event topology variables, flavor-tagging information, vertex-fitting information, and kinematic-fit information. All of these variables are required to have little or no correlation with ΔE and M_{bc} .

3. Signal extraction and measurement results

We use unbinned maximum likelihood fits to extract signal yields from the data to calculate various physics observables. In the $B \rightarrow hh$ and $B \rightarrow hhh$ ($h = K$ or π) analyses, only ΔE is fit for events restricted to $M_{bc} > 5.27$ GeV/ c^2 . The fits to the two $B \rightarrow \phi K$ modes use five variables: ΔE , M_{bc} , output of the continuum suppression BDT discriminator (C'_{out}), K^+K^- candidate mass ($m_{K^+K^-}$), and cosine of the ϕ candidate's helicity angle ($\cos\theta_{H,\phi}$). The fits to the two $B \rightarrow \phi K^*$ modes use seven variables: $K^+\pi^-$ candidate mass ($m_{K\pi}$) and cosine of the K^* candidate's helicity angle ($\cos\theta_{H,K^*}$) in addition to those used in $B \rightarrow \phi K$ modes. By fitting the data, we determine the following quantities:

- Branching fractions: $\mathcal{B} = \frac{N}{\epsilon \times 2 \times N_{BB}}$, where N is the signal yield, ϵ is the signal reconstruction efficiency, which is determined from simulation and validated with control samples, and N_{BB} is the number of $B\bar{B}$ events (19.7×10^6 for B^+B^- and 18.7×10^6 for $B^0\bar{B}^0$). N_{BB} is obtained from the measured integrated luminosity, the exclusive $e^+e^- \rightarrow \Upsilon(4S)$ cross section, and $\mathcal{B}(\Upsilon(4S) \rightarrow B^0\bar{B}^0)$ [9].
- CP asymmetries: The raw asymmetries are obtained as $\mathcal{A} = \frac{N(b) - N(\bar{b})}{N(b) + N(\bar{b})}$, where $N(b)$ and $N(\bar{b})$ are the yields of the final-state mesons containing b and \bar{b} flavors, respectively. The CP asymmetry \mathcal{A}_{CP} is obtained by subtracting the instrumental asymmetry \mathcal{A}_{det} from \mathcal{A} , where $\mathcal{A}_{\text{det}}(K^+\pi^-) = -0.010 \pm 0.003$ and $\mathcal{A}_{\text{det}}(K_S^0\pi^+) = -0.007 \pm 0.022$ are measured on large samples of $D^0 \rightarrow K^-\pi^+$ and $D^+ \rightarrow K_S^0\pi^+$ decays with negligible CP violation. Then $\mathcal{A}_{\text{det}}(K^+)$ is determined as $\mathcal{A}_{\text{det}}(K^+) = \mathcal{A}_{\text{det}}(K^+\pi^-) - \mathcal{A}_{\text{det}}(K_S^0\pi^+) + \mathcal{A}_{\text{det}}(K_S^0) = -0.015 \pm 0.022$, where an upper bound on $\mathcal{A}_{\text{det}}(K_S^0)$ is used based on the previous measurements [10].
- Longitudinal polarization fractions: $f_L = \frac{N_L/\epsilon_L}{N_L/\epsilon_L + N_T/\epsilon_T}$, where $N_{L(T)}$ and $\epsilon_{L(T)}$ are the signal yield and signal reconstruction efficiency with longitudinal (transverse) polarization, respectively. The distinctive helicity-angle distributions allow the separation of the two signal components.

Figures 1 and 2 show the ΔE distributions in data for $B^0 \rightarrow K^+\pi^-$ and $B^+ \rightarrow K^+K^-K^+$, with fit projections overlaid. Figure 3 shows the ΔE , M_{bc} , C'_{out} , $m_{K^+K^-}$, and $\cos\theta_{H,\phi}$ distributions in data

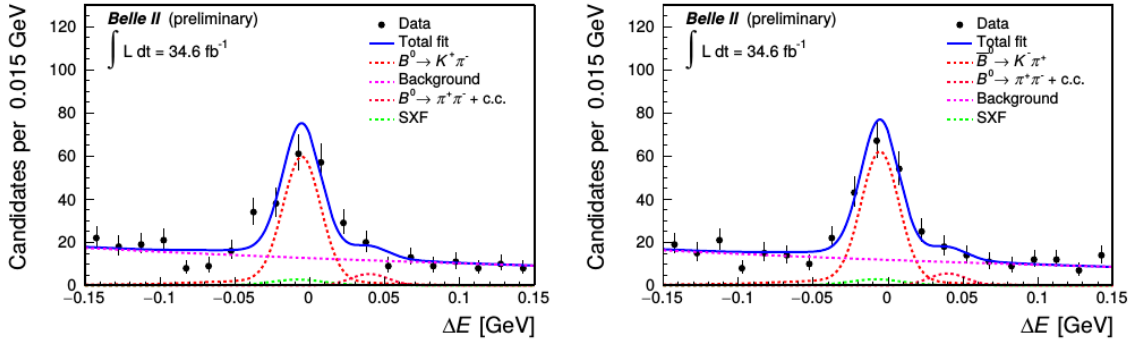


Figure 1: Distribution of ΔE for $B^0 \rightarrow K^+\pi^-$ (left) and $B^0 \rightarrow K^-\pi^+$ (right) decays with fit projections overlaid.

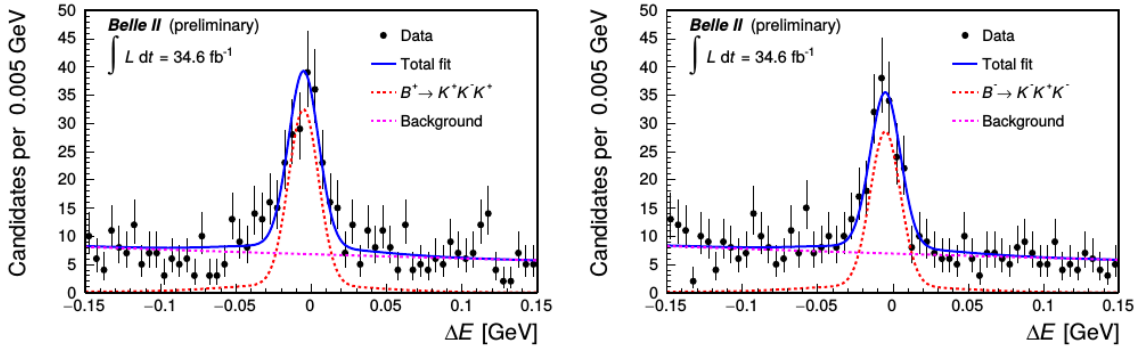


Figure 2: Distribution of ΔE for $B^+ \rightarrow K^+K^-K^+$ (left) and $B^- \rightarrow K^-K^+K^-$ (right) decays with fit projections overlaid.

for $B^+ \rightarrow \phi K^+$, with fit projections overlaid. Figure 4 shows the ΔE , M_{bc} , C'_{out} , $m_{K^+K^-}$, $\cos\theta_{H,\phi}$, $m_{K\pi}$, and $\cos\theta_{H,K^*}$ distributions in data for $B^+ \rightarrow \phi K^{*+}$ decay, with fit projections overlaid. The major systematic uncertainties come from tracking, PID, and fit modelling. All the measurement results are summarized in Table 1.

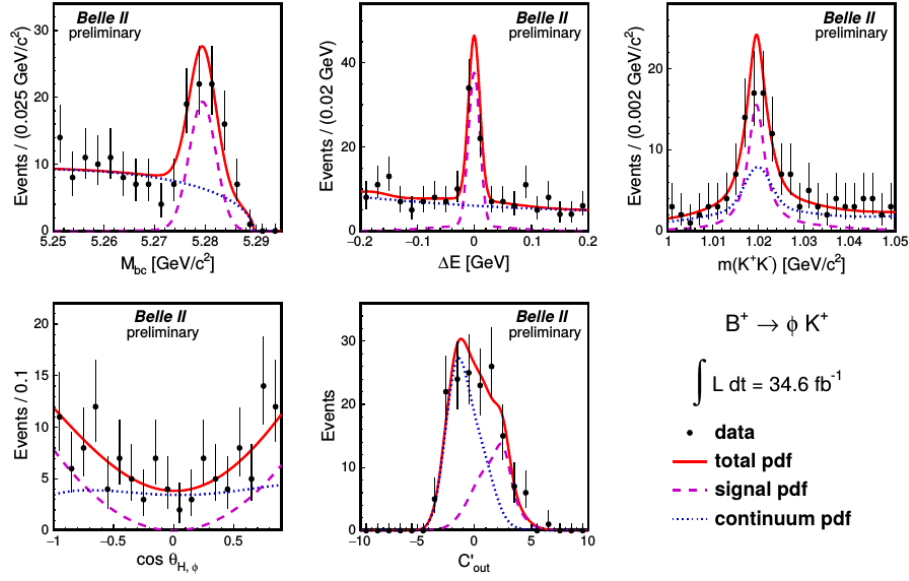


Figure 3: Distribution of ΔE , M_{bc} , C'_{out} , $m_{K^+K^-}$, and $\cos\theta_{H,\phi}$ for $B^+ \rightarrow \phi K^+$ decay with fit projections overlaid.

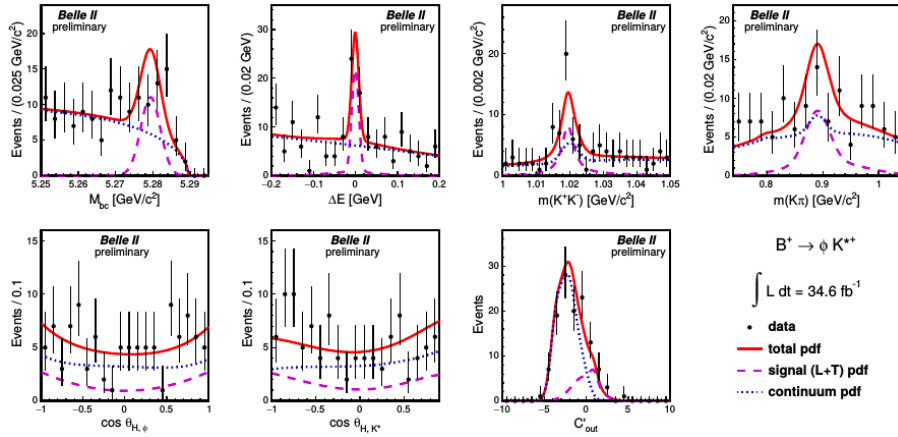


Figure 4: Distribution of ΔE , M_{bc} , C'_{out} , $m_{K^+K^-}$, $\cos\theta_{H,\phi}$, $m_{K\pi}$, and $\cos\theta_{H,K^*}$ for $B^+ \rightarrow \phi K^{*+}$ decay with fit projections overlaid.

Table 1: Summary of measurement results. The first uncertainties are statistical and the second ones are systematic.

Mode	\mathcal{B} (10^{-6})	\mathcal{A}_{CP}	f_L
$B^0 \rightarrow K^+ \pi^-$	$18.9 \pm 1.4 \pm 1.0$	$0.030 \pm 0.064 \pm 0.008$	-
$B^0 \rightarrow \pi^+ \pi^-$	$5.6^{+1.0}_{-0.9} \pm 0.3$	-	-
$B^+ \rightarrow K^+ \pi^0$	$12.7^{+2.2}_{-2.1} \pm 1.1$	$0.052^{+0.121}_{-0.119} \pm 0.022$	-
$B^+ \rightarrow \pi^+ \pi^0$	$5.7 \pm 2.3 \pm 0.5$	$-0.268^{+0.249}_{-0.322} \pm 0.123$	-
$B^+ \rightarrow K^0 \pi^+$	$21.8^{+3.3}_{-3.0} \pm 2.9$	$-0.072^{+0.109}_{-0.114} \pm 0.024$	-
$B^0 \rightarrow K^0 \pi^0$	$10.9^{+2.9}_{-2.6} \pm 1.6$	-	-
$B^+ \rightarrow K^+ K^- K^+$	$32.0 \pm 2.2 \pm 1.4$	$-0.049 \pm 0.063 \pm 0.022$	-
$B^+ \rightarrow K^+ \pi^- \pi^+$	$48.0 \pm 3.8 \pm 3.3$	$-0.063 \pm 0.081 \pm 0.023$	-
$B^0 \rightarrow \phi K^0$	$5.9 \pm 1.8 \pm 0.7$	-	-
$B^+ \rightarrow \phi K^+$	$6.7 \pm 1.1 \pm 0.5$	-	-
$B^0 \rightarrow \phi K^{*0}$	$11.0 \pm 2.1 \pm 1.1$	-	$0.57 \pm 0.20 \pm 0.04$
$B^{*+} \rightarrow \phi K^{*+}$	$21.7 \pm 4.6 \pm 1.9$	-	$0.58 \pm 0.23 \pm 0.02$

4. Summary

Belle II reports its first measurements of charmless B decays with a data sample corresponding to 34.6 fb^{-1} . The measurements include branching fractions, CP asymmetries, and longitudinal polarization fractions. All the results are in agreement with the known values, and provide a good validation of the detector performance and analysis strategies.

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