Search for an association between neutrinos and radio-selected blazars with ANTARES

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Recently, evidence for an association between high energy neutrinos detected by IceCube and radio-selected blazars has been found by Plavin et al. (2020, 2021). This result was achieved using an all sky complete sample of 3411 blazars selected on their parsec-scale flux density at 8 GHz higher than 150 mJy. We perform a positional correlation analysis using the same sample of radio-selected blazars, with the latest point source sample of neutrinos extracted from the data collected by the ANTARES detector between January 29, 2007 and February 28, 2020. Preliminary results are presented and discussed.
1. Introduction

Almost ten years after the discovery of a diffuse astrophysical neutrino flux in 2013 by the IceCube Collaboration [1], the question of the origin of this excess is still not resolved. However, the recent studies using IceCube data ([2], [3], [4]), support the idea that extra-galactic sources, and in particular blazars could produce a detectable part of the observed high-energy neutrino flux.

The blazars are of particular interest because of the relativistic Doppler boost enhancing the flux of electromagnetic radiation and potentially associated neutrinos, making these objects detectable up to cosmological distances. Previous searches have been conducted using the gamma-ray selected blazars from Fermi catalogs, with IceCube data [5], [6], [7] and with ANTARES [8].

Recently, a correlation between radio-selected blazars and the IceCube very high energy track-like events has been reported in [9]. Moreover, the association has been confirmed in a second study using the lower-energy neutrinos from the IceCube 7yr public point source sample [10]. In the present study, the possible correlation between these radio-selected blazars and the neutrinos from the latest 13yr ANTARES point source sample is investigated.

2. The blazar catalog

The emission of high-energy neutrinos is a signature of hadronic interactions between accelerated charged particles and ambient matter or radiation. This processes could take place either in the central regions of galaxies: accretion disk, base of the jet; or in the outer and extended parts: kiloparsec-scale jets, blobs, or lobes and hot spots. Considering the distance of these sources, the angular resolution of present and future neutrino telescopes is not sufficient to differentiate between those acceleration regions.

The very long baseline interferometric (VLBI) radio measurement is a unique way to distinguish the emission of the central part of a galaxy from the outer parts, by resolving the central region of AGNs on a parsec scale, up to cosmological distances. The vast majority of the blazars used in the present analysis have a VLBI radio emission that is dominated by the base of the observed jet, and thus allows to test the for a possible neutrino emission from the central regions of AGN.

The 8 GHz VLBI observations have been compiled by authors of [9] to provide a flux-limited complete subsample of AGNs on the full sky. The quantity $S_{8\text{GHz}}$, the flux density integrated over VLBI images at 8 GHz, is used to selected objects by requiring $S_{8\text{GHz}} > 0.15\text{Jy}$. The complete sample contains 3411 objects, with 3027 in the field of view of ANTARES and available in the following link: http://astrogeo.org/rfc/ (version being used: 2020d).

3. The ANTARES data set

The ANTARES neutrino telescope is a water Cherenkov detector that operates since 2007 on the bottom of the Mediterranean Sea, 40-km off-shore Toulon (France). The detection method is based on the observation of the Cherenkov light emitted in water by the relativistic particles produced by interactions of neutrinos in the vicinity of the instrumented volume [11].
The data sample used in the present analysis consists of 10162 track-like events after reconstruction and selection, that have been detected by ANTARES between January 29, 2007 and February 28, 2020, cumulating a total livetime of 3845 days. These events are almost totally induced by charged current interactions of $\nu_\mu$, producing a high energy muon that can propagate on large distances in water.

The event selection has been optimized for the detection of point-like sources, assuming a simple power-law energy spectrum $\propto E^{-2}$. A detailed description of the selection procedure can be found in [12].

For the current analysis, an updated time calibration has been used, the reconstruction and selection algorithms have been re-applied to the full ANTARES data to get the most accurate estimation of neutrino candidates arrival direction. The events contained in the previous 11yr point source sample used in [8] are not all included in this new 13yr sample as their updated reconstructed parameters may not fulfill the selection criteria. The overlap between the two samples is 7287 events, out of the 8754 present in the previous one.

The energy estimator used in the present analysis relies on the measure of muon energy deposit in water [13]. The selected tracks have estimated neutrino energies ranging from $\sim 100$ GeV to $\sim 1$ PeV. The angular resolution depends on the energy and the topology of each event, but in average it is better than 0.4° above 10 TeV. Due to the detector condition evolving with time (loss of optical modules, ageing of the photo-multipliers), the median angular resolution shows a significant variation over the years. This behavior is taken into account for the likelihood analysis presented in the section 5, by using Monte Carlo simulations [14] to build the pdfs separately for two different time periods: before 2014 and after 2014.

4. Counting method

Following the approach used in [9], a simple method based on counting neutrino-blazar pairs is used as a first analysis. The quantity that is computed is the number of neutrinos that are located at an angular distance $\Psi$ from a blazar less than $x \cdot \beta$, where $\beta$ is the angular uncertainty coming from the neutrino reconstruction, and $x$ is a free parameter that is varied in the intervall [0; 2].

This $x$ parameter is introduced to take into account a possible systematic difference between the output of the reconstruction algorithm and the true (unknown) angular error radius. With the help of Monte-Carlo simulations, it is however possible to study the relation between the error estimate $\beta$ and the true 68% containment radius $\Psi_{68\%}$. For values $\beta \leq 0.5$, and reconstructed energies above $\sim 10$ TeV, the relation between $\beta$ and $\Psi_{68\%}$ is close to what is expected from a two-dimensional Gaussian function $\Psi_{68\%} \approx 2 \cdot \beta$. For higher values of $\beta$ or at lower energies, the value of $\Psi_{68\%}$ starts to deviate from Gaussian behavior, by being significantly higher. The ANTARES point spread function (PSF), that is used in the following likelihood analysis, also shows the same significant difference compared to a Gaussian function.

The use of the $x$ parameter is then a simple empirical way to take into account this complicated behavior while still using the information of the quality of reconstruction of ANTARES neutrinos on an event-by-event basis. As the value of this parameter is scanned, a trial factor correction is applied to get a correct p-value estimation.
The result of the counting analysis applied to the full catalog of blazars and the total ANTARES 13yr point source sample is shown in figure 1 (left). The absolute minimum is found for $x = 0.81$, where 451 pairs are observed in data, while 389 are expected from random simulations (62 pairs in excess). The associated pre-trial p-value is $p = 1.2 \times 10^{-3}$ (3.2 $\sigma$), leading after correction to a post-trial p-value of $P = 2.2 \times 10^{-2}$ (2.3$\sigma$).

In [9], the blazars having associated IceCube neutrinos are found to have a higher than average radio flux density (time-integrated). The radio-neutrino correlation is therefore studied with ANTARES neutrinos by performing an additional scan on the radio flux density $S_{8\text{GHz}}$. Blazars are kept in the sample if they satisfy $S_{8\text{GHz}} > S_{\text{min}}$.

The result of this two dimensional scan ($x, S_{\text{min}}$) is shown in figure 1 (right). The global minimum is found for $x = 0.8$ and $S_{\text{min}} = 0.15$ Jy, with a pre-trial p-value $p = 1.2 \times 10^{-3}$. This minimum corresponds to the previous findings of the one dimensional scan, and is obtained for the lowest value of the flux density cut, meaning that the whole blazar catalog is included. This means that the potential ANTARES neutrinos-blazar association is not eventually dominated by the high flux sources.

However, a second local minimum is observed for $x = 0.45$ and $S_{\text{min}} = 3.7$ Jy, with a pre-trial p-value $p = 3.2 \times 10^{-3}$. This excess is caused by 4 neutrino-blazar pairs, while the expected number of chance coincidences is only 0.5, with blazars brighter than 3.7 Jy at separations below 0.45$\beta$. There are only 20 objects in the catalog that have a flux higher than 3.7 Jy. This motivates a manual inspection of the matching objects, their properties are listed in table 1.

All four neutrino candidate events can be considered notable: (i) two of them have high energies, top 0.3-0.6% of the whole dataset; (ii) two are associated with the same blazar J0538-4405; and (iii) the fourth event is associated with J1743-0350, one out of four likely associations with high-energy IceCube tracks found in [9].

The radio light curves measured in VLBI at 2 and 8 GHz are shown in figure 2 for J0538-4405. There is a major flare visible around 2011-2012, and a sign of another flare in 2018 (end
ANTARES neutrinos association with radio-selected blazars
Julien Aublin

<table>
<thead>
<tr>
<th>Blazars</th>
<th>Neutrinos</th>
</tr>
</thead>
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<tr>
<td>Source name</td>
<td>Arrival time</td>
</tr>
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<td>J0609-1542</td>
<td>2011-01-30</td>
</tr>
<tr>
<td>J1743-0350</td>
<td>2019-03-08</td>
</tr>
<tr>
<td>J0538-4405</td>
<td>2011-08-09</td>
</tr>
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<td></td>
<td>2018-03-20</td>
</tr>
</tbody>
</table>

Table 1: Properties of the 4 neutrino-blazar pairs found for objects brighter than 3.7 Jy.

of observations). These periods are close to the dates of detected neutrinos, 2011-08-09 and 2018-03-20.

Additionally, J0538-4405 is a γ-bright source (labelled as 4FGL J0538.8-4405 in the Fermi 4FGL catalog [15]) and has a rich Fermi light-curve where strong γ-ray activity is apparent around 2010-2012. Unfortunately, the two ANTARES neutrino candidates have a large angular error β value, that favors an atmospheric muon origin rather than an astrophysical origin.

Figure 2: Left: arrival directions of ANTARES track events in equatorial coordinates around the position of the blazar J0538-4405. The neutrino events are represented by dashed empty circles with a radius equal to the angular error estimate β while the estimated energy is indicated by the color code. The thin black circles centered around the source (red star marker) indicate the 1° and 5° angular distances to the source. The position of other blazars in the catalog are represented by a green star. Right: Radio light-curves of the blazar J0538-4405 measured in VLBI at three different frequency (see color code). The arrival times of the ANTARES neutrino candidates are indicated by a vertical dashed line.

The other two blazars J0609-1542 and J1743-0350, also experience radio flares close to the time of neutrino detection (see presentation). J1743-0350 flares both in 2019 when ANTARES received a neutrino, and in 2011 when the corresponding IceCube high-energy track occurred [9].

5. Likelihood analysis

A complementary likelihood analysis is presented here, making use of more information than the simple counting method about the neutrino candidates via the construction of energy and
antares neutrinos association with radio-selected blazars

JULIEN AUBLIN

declination-dependent probability density functions. A pure power-law energy spectrum $E^{-\gamma}$ is assumed for the potential astrophysical neutrinos, and three values are tested $\gamma = 2.0$, 2.25 and 2.5.

The log-likelihood for both the null hypothesis $H_0$ and the alternative hypothesis $H_1$ is written as:

$$\ln \mathcal{L}(H_0) = \sum_{i}^{N} \ln (\mu_b B_i) - \mu_b,$$

$$\ln \mathcal{L}(H_1) = \sum_{i}^{N} \ln (\mu_s S_i + \mu_b B_i) - \mu_s - \mu_b,$$

where $N$ is the total number of observed neutrino candidates, $S_i$ is the probability density function (PDF) of the signal and $B_i$ is the background PDF. The free parameters are the estimated number of signal $\mu_s$ and background $\mu_b$ events. The test statistic is then computed as: $\lambda = \ln \left( \frac{\max(\mathcal{L}(H_1))}{\max(\mathcal{L}(H_0))} \right)$.

The signal and background PDF are defined as the product of a spatial and an energy term:

$$S_i = f_s(\alpha_i, \delta_i) \cdot g_s(E_i) \quad \text{and} \quad B_i = f_b(\delta_i) \cdot g_b(E_i), \quad (2)$$

where $(\alpha, \delta)$ are the equatorial coordinates and $E$ the estimated energy of the neutrinos. The spatial term for the background is independent of the right ascension, and is estimated via a polynomial fit the real data $\sin \delta$ distribution. The spatial term is obtained by summing over all the sources contributions:

$$f_s(\alpha_i, \delta_i) = \frac{1}{\sum_{j=1}^{N_{sources}} w_j \mathcal{T}_j(\alpha_i, \delta_i)}, \quad w_j = w_j^{model} A(\delta_j), \quad (3)$$

where $\mathcal{T}_j(\alpha_i, \delta_i)$ is the point spread function evaluated at the $j^{th}$ blazar’s position, with weight $w_j$. The weight of the $j^{th}$ source takes into account the declination-dependent acceptance $A(\delta)$ of the neutrino track sample computed for the corresponding $E^{-\gamma}$ energy spectrum. Two different weighting choices are considered: (i) an equal weight $w_j^{model} = 1$, that could be considered as the default assumption, and (ii) a flux weight $w_j^{model} = S_{\nu \mathrm{G} \mathrm{Hz}}$ to test the potential correlation between neutrino emission and radio flux density.

The point-spread function $\mathcal{T}_{E, \beta, \delta}$, defined as the probability density for the event direction to fall within a given angular distance from the true source direction, is obtained from Monte Carlo simulations, and has been parametrized for different energy, $\beta$ and declination intervals. This procedure is performed separately for events before and after RP1, to take into account the better detector performance in the early period.

The results of the likelihood analysis using the catalog of VLBI-selected blazars are summarized in table 2. For the full sample, the $p$-value obtained with the likelihood is between 5% and 10%, depending on the assumed spectral index, the minimum being obtained for the softer energy spectrum and with a weight proportional to the radio flux density.

For comparison, those $p$-values are a factor $\sim 2.5 - 5$ higher than the post-trial $p$-value found with the counting method, while for soft spectral indexes the number of fitted signal events $n_s$ is comparable to the 62 neutrino-blazar pairs in excess previously reported. When considering only blazars with a flux density $S_{\nu \mathrm{G} \mathrm{Hz}} > 3.7 \mathrm{Jy}$, the $p$-values found with the likelihood are between $\sim (2 - 5) \times 10^{-3}$, in agreement with the $p$-value $p = 3.2 \times 10^{-3}$ found with the counting method.
We note that differences in the resulting p-values between the two methods are to be expected, as the likelihood makes use of the additional energy information, and is based on Monte Carlo predictions for the point spread function.

<table>
<thead>
<tr>
<th>Sample</th>
<th>Spectral index</th>
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<tbody>
<tr>
<td></td>
<td></td>
<td>$n_s$</td>
<td>$\lambda$</td>
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<tr>
<td>Full VLBI</td>
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<td>4.33</td>
</tr>
<tr>
<td></td>
<td>$E^{-2.25}$</td>
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<td></td>
<td>$E^{-2.5}$</td>
<td>186</td>
<td>9.76</td>
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<tr>
<td>$S_{\text{GHz}} &gt; 3.7$ Jy</td>
<td>$E^{-2}$</td>
<td>8</td>
<td>4.84</td>
</tr>
<tr>
<td></td>
<td>$E^{-2.25}$</td>
<td>10</td>
<td>5.16</td>
</tr>
<tr>
<td></td>
<td>$E^{-2.5}$</td>
<td>11</td>
<td>4.84</td>
</tr>
</tbody>
</table>

Table 2: Summary of the results obtained with the likelihood analysis: the fitted number of signal events $n_s$, the test statistic $\lambda$ and the p-value $p$ are reported for each combination of sample, spectral index and weighting choice.

6. Discussion and conclusion

A search for an association between VLBI radio-selected blazars and the arrival directions of the track-like events detected by ANTARES in 13 years of operation has been performed with two different methods, a simple counting technique and a more refined likelihood analysis. When searching on the full VLBI catalog, the two methods find an excess with p-values $p = 0.02$ and $p \in [(2 - 5) \times 0.05 - 0.1]$ for the counting and the likelihood respectively.

An additional scan in radio-flux density shows that the bulk of the blazars contributing to this excess do not have a higher than average flux density (time-integrated). However, a second excess at very high flux is found with p-values $p = (2 - 5) \times 10^{-3}$, dominated by three blazars 0609-1542, J1743-0350 and J0538-4405, that are examined individually. The radio light-curves of these blazars show signs of flaring activity around the arrival time of neutrinos, especially for J0538-4405 that experiences a major augmentation of its radio flux over $\sim 4$ years.

An estimation of the p-value of the space and time coincidence between ANTARES neutrinos and radio emission has not yet been performed for those three sources, and will be addressed in a future study.

We advertise the special case of the blazar J0242+1101 that is presented in these conference [16], showing a hint of a spatial and temporal association between ANTARES neutrinos and a long radio flare, together with a gamma-ray flare observed by Fermi. The particular case will require a dedicated analysis to estimate the chance probability of the association between the VLBI, gamma-ray and neutrinos observations.

References


ANTARES neutrinos association with radio-selected blazars

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