KM3NeT performance on oscillation and absorption tomography of the Earth

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The KM3NeT neutrino telescope, currently under construction, consists of two detectors in the Mediterranean Sea, ORCA and ARCA, both using arrays of optical modules to detect the Cherenkov light produced by charged particles created in neutrino interactions. Although originally designed for neutrino oscillation and astrophysical research, this experiment also bears unprecedented possibilities for other fields of physics. Here we present its performance for neutrino tomography, i.e. the study of the Earth’s internal structure and composition. Owing to the different energy ranges covered by its two detectors ORCA and ARCA, KM3NeT will be the first experiment to perform both oscillation and absorption neutrino tomography. Resonance effects in the oscillations of GeV neutrinos traversing the Earth will allow KM3NeT/ORCA to measure the electron density along their trajectory, leading to potential constraints of the proton-to-nucleon (Z/A) ratio in the traversed matter. Absorption tomography aims at the detection of neutrinos in the TeV-PeV range with KM3NeT/ARCA. At PeV energies, the Earth is opaque for neutrinos which leads to a reduction of the upgoing neutrino flux at the detector side from which conclusions can be drawn about the density of the inner layers of the Earth. We show here first sensitivity studies of the potential of KM3NeT to address open questions of geophysics concerning the chemical composition and matter distribution in the Earth’s core and mantle through neutrino tomography.

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1. Introduction

The inner structure of the Earth can be simply described as a superposition of different shells of uniform chemical compositions. The planets innermost region is subdivided into a solid core formed of Fe-Ni, with an approximate radius of 1200 km, surrounded by a molten outer core with a radius of 3480 km and slightly lower density. Hereafter comes the mantle, in first order uniformly consisting of pyrolite. The outermost layer is a thin (<100 km) silicate crust. All this knowledge comes from seismic data, geodetic measurements and experimental petrology and led in combination with constraints on radius, total mass and momentum of inertia to the Preliminary Earth Model (PREM) [2]. It is widely used as reference density profile of the Earth and illustrated in Fig. 1.

The upcoming generation of large neutrino telescopes like KM3NeT [4] provide new ways of studying the Earth’s interior, mostly independently of seismic measurements. One the one hand, neutrinos with energy larger than a few tens TeV start to get absorbed during their propagation through the Earth. The absorption probability increases with the neutrino energy and the amount of matter along the neutrino path. Energy- and zenith-dependent measurements of the neutrino rates with KM3NeT/ARCA can therefore give insights about the density of the Earth and its layers. This method has been used recently with a one-year sample of muon neutrinos collected with the IceCube neutrino telescope, confirming the feasibility of absorption tomography of the Earth [1]. On the other hand, KM3NeT/ORCA will detect ~GeV neutrinos whose mean free path is much bigger than the Earth’s diameter. Instead of being absorbed, those lower energetic neutrinos "feel" the presence of matter, i.e electrons, as an extra potential to their propagation Hamiltonian. This leads to resonance effects in their oscillation probability, which depends on the electron density $n_e$ within the traversed matter, where

$$n_e = \frac{N_A}{m_n} \times \frac{Z}{A} \times \rho_m,$$  \hspace{1cm} (1) 

with the Avogadro number $N_A$ and the nucleon mass $m_n$. The effective proton-to-nucleon ratio $Z/A$ varies among chemical elements, e.g. it is 0.4655 for Fe and 1 for H. It can therefore be used to constrain the nature and abundances of light elements in the outer core, which can not be revealed by seismologic measurements only, but whose presence is required to explain the jump in density between inner and outer core [3]. Constraining the inner Earth composition will also improve our understanding of the dominant processes that led to the formation of our planet. Our analysis combines models for the neutrino flux, and cross section, and for the inner structure of the Earth to compute the number of events at the detector site. The neutrino oscillation or absorption probabilities are affected by the Earth parameters, i.e. $Z/A$ and $\rho_m$. The detector response is simulated using a response matrix calculated with MC simulations (discussed in Sec.2). We find that after 10 years of data taking with two KM3NeT/ARCA building blocks, we will be able to profile the Earth density with comparable sensitivity as a similar study by IceCube [1]. With 10 years of KM3NeT/ORCA data we can set limits on the outer cores $Z/A$ and density (see Sec.3).
2. Methodology

2.1 Underlying signal

The input model of the atmospheric neutrino flux is taken from Honda(2015)\(^{\text{1}}\) [7] for energies 1-100 GeV. Between 100 GeV and 100 PeV we use a combination of the conventional flux taken from Honda(2007)\(^{\text{2}}\) [8] and a prompt neutrino flux from Enberg [9]. To simulate the effects of matter on neutrinos traversing the Earth, for both oscillation and absorption, the OscProb software [10] is used.

For a given incident angle \(\theta_z\), the trajectory of a neutrino through the Earth is modelled along the corresponding baseline with a sequence of steps of constant matter and \(Z/A\) and thus the electron density (according to Eq. 1). The latter induces an effective potential for charged-current interactions \(V_{CC} = \pm \sqrt{2} G_F N_e(x)\) (where \(\pm\) applies respectively to \(\nu_e\) and \(\bar{\nu}_e\)), which modifies the neutrino oscillation probabilities. These are calculated from the quantum evolution equations for the neutrino states for each shell along the neutrino path. The absorption probability within each shell is calculated using the simple exponential function

\[
P_{Abs}(E, \Theta) = \Pi_i e^{-l_i(\theta) \rho_i \frac{\sigma_i}{\rho_i} \sigma(E)}
\]

with \(l\) the path length through the respective shell \(i\), \(\rho\) its density, \(u\) the atomic mass number and \(\sigma\) the neutrino-nucleon cross-section weighted for water molecules, obtained with the GENIE Monte Carlo generator [11].

The latter is also used to compute the rate of events interacting in the detector. This differential rate (corresponding to the number of neutrino interactions of given flavour, energy \(E\) and zenith \(\theta\) occurring in the detector per unit exposure\(^{\text{2}}\)) is computed for each (anti-)neutrino flavour \(\beta\) as:

\[
\frac{\text{d}N_{\beta}^{\text{int}}(E, \theta)}{\text{d}E \text{d}\theta} = \sigma_{\nu\beta}(E) \cdot \sum_{\nu_\alpha} P_{\nu_\alpha \rightarrow \nu_\beta}(E, \theta) \cdot (1 - P_{Abs,\beta}(E, \theta)) \cdot \frac{\text{d}\Phi_{\nu_\alpha}}{\text{d}E \text{d}\theta}(E, \theta)
\]

\(^{\text{1}}\)We have used the tables for Gran Sasso site without mountain, azimuth-averaged, at solar minimum.

\(^{\text{2}}\)The term exposure corresponds to the product of the detector lifetime and its effective mass [Mton years]
where $\sigma_{\nu\nu}$ is the total interaction cross-section for neutrino type $\nu_\beta$, $P_{a\to\beta}$ is the $\nu_a \to \nu_\beta$ oscillation probability, and $d\Phi_{\nu_a}/dE\,d\theta$ is the unoscillated differential flux of atmospheric neutrinos at the detector location. For studies with ARCA we neglect oscillation effects, i.e. $P_{a\to\beta} = \delta_{a\beta}$, for studies with ORCA we neglect absorption effects, i.e. $P_{\text{Abs}} = 0$.

2.2 Detector response modelling

The KM3NeT collaboration is currently building a network of water Cherenkov neutrino telescopes in the Mediterranean sea. The two detectors ORCA and ARCA (Oscillation/Astrophysical research of cosmosics in the abyss[4]) are using the same technical components, 3D arrays of digital optical modules (DOM) hosting 31 small photomultipliers each to detect Cherenkov light from particle showers of neutrino-nuclei interactions, distributed along slender strings, anchored at the sea bottom at a depth of about 2500 and 3500 m respectively. The KM3NeT DOMs detect the Cherenkov light induced by charged particles that are generated in the interaction of a neutrino with the matter surrounding the array. The detector layout, i.e. horizontal and vertical spacing of the DOMs, determines the threshold energy that is needed to trigger an event in the detector. KM3NeT/ORCA, with an instrumented volume equivalent to $\sim 7$ Mtons, will cover the 1-100 GeV energy range, while KM3NeT/ARCA, with $\sim 1$ Gton fiducial mass, detects neutrinos at higher energies. One distinguishes between two event topologies, track- and shower-like. A $\nu_\mu$ that interact with a nucleus via a charged current (CC) create a $\mu$ that will propagate through the detector, and inducing the emission of Cherenkov light along its trajectory. All $\nu_e$ and neutral-current (NC) interactions induce an approximately spherical light emission around the interaction point and are thus referred to as shower-like events. A special case are $\nu_\tau$-CC, where the produced $\tau$-lepton can either decay into a $e$ or a $\mu$. Note that we did not simulate $\nu_\tau$-events for ARCA, because their total contribution is negligible (only produced by charmed meson decays).

Different approaches to modelling the KM3NeT detector response have been used so far in the literature. The analysis presented in Bourret (2019) [12] used parametrized response
functions that allowed to perform sensitivity studies for different detector types. In the present work, we use instead the same approach as in Bourret (2017), based on a full event-by-event Monte Carlo simulation that reproduces more accurately the specific response of the KM3NeT detectors [13]. Based on these simulations, we create a response matrix \( R \) with a binning in five dimensions \((E_{\text{True}}, \theta_{\text{True}}, E_{\text{Reco}}, \theta_{\text{Reco}}, f)\) where \( f \) encrypts the interaction channel and reconstructed class (track- or shower-like). For a given event with true characteristics \((E_{\text{True}}, \theta_{\text{True}})\), the response function \( R \) provides the probability for this event to be classified in a specific class and \((E_{\text{Reco}}, \theta_{\text{Reco}})\) bin.

### 2.3 Statistical methods

With the methods described above we can now simulate experiments based on given input models, in our case changing density and \( Z/A \) values for different layers of the Earth. We compute then the log-likelihood ratio for two Asimov datasets, i.e. pseudo-experiments where the data corresponds to the mean of the expected results given a set of input parameters. Assuming that the number of events per bin follows Poisson statistics, one can define

\[
\chi^2_{\text{hyp}}(\text{data}) = \sum_{b=1}^{\text{Nh}_{\text{bins}}} 2\left( (\mu_{\text{hyp}})_b - n_b + n_b \ln \left( \frac{n_b}{(\mu_{\text{hyp}})_b} \right) \right)
\]

which corresponds to the likelihood to measure the simulated \textit{data} with an underlying true hypothesis \textit{hyp}. The resulting Asimov sensitivity \( S_{\text{Asimov}} = \sqrt{\Delta \chi^2} \) is in good agreement with the sensitivity obtained from the log likelihood ratio (with some caveats (see [13])), where

\[
\Delta \chi^2 = \chi^2_{\text{prem}}(\text{data}_{\text{prem}}) - \chi^2_{\text{hyp}}(\text{data}_{\text{prem}}) = -\chi^2_{\text{hyp}}(\text{data}_{\text{prem}}).
\]

In this case \( S_{\text{Asimov}} \) corresponds to the confidence level to which the hypothetical Earth model can be distinguished from PREM.

### 3. Results and discussion

#### 3.1 Absorption tomography

Analogously to the analysis reported in [1], performed with 1-year data of IceCube, we use an Earth model with 5 layers with radii of \([1242, 2372, 3502, 4935, 6368]\)(km), such that we have one layer for the inner core and two layers for outer core and mantle, respectively.

We fit simultaneously the densities of the 5 layers to our simulated data and determine the 1\(\sigma\) uncertainties with \textsc{root}Minuit2 [14] without taking into account any systematic uncertainty.

The result shown in Fig. 3 is comparable with the prediction for 10 years of IceCube data presented in [1] (note that our first bin \([0-1242 \text{ km})\] appears to be larger due to the logarithmic \( y \)-scale). The width of the error bands decreases with the total volume of the respective layer. While only a fraction of the detected neutrinos traversed the inner core, all of them pass through the mantle and thus the statistical uncertainty for the density fit of the mantle is reduced. Although finer features of the Earth’s density profile like the density step between inner and outer core can not be resolved, ARCA is able to distinguish the core from the mantle region.
**3.2 Oscillation tomography**

Our goal for the oscillation tomography is to constrain the electron density $n_e$ in the outer core and find hints on its chemical composition. Because $n_e$ is the direct product of the matter density $\rho_m$ and the proton-to-nucleon ratio $Z/A$ (see Eq. 1), a fit of $n_e$ is analogous to either a fit of the $Z/A$ value or a fit of a normalisation factor $\tilde{\rho}$ for the density as illustrated in Fig. 4 (left), while keeping the respective other parameter constant. We fit the $Z/A$ value assuming the PREM density (i.e. $\tilde{\rho} = 1$) and determine its error intervals, from which we can draw the contour lines according to equal values of $n_e$ (Fig. 4 (right)).

We used no priors for the fit to remain completely independent of seismological measurements. PREM uncertainties on the outer core density are smaller than a few percent and would therefore heavily constrain the possible combinations of $Z/A$ and $\rho$. In the future, the constraints from absorption tomography on the outer core density could help to narrow down the contour bands.
Realistic models of the chemical composition of the outer core provide values between 0.466 (iron nickel alloy as in the inner core) and 0.4714 ("exotic" model containing 1 wt% $^3$H [15]) and will not be distinguishable with the current generation of neutrino detectors. The result of oscillation studies presented in the analysis relies on the global best fit of oscillation parameters assuming normal mass ordering [16] and does not include systematic uncertainties. ORCA and other neutrino oscillation experiments will most likely determine the true mass ordering in a shorter time scale than 10 years which is the time scale considered in this simulation and also other oscillation parameters will be determined with a higher precision. An estimation of systematic effects showed that events only traversing the mantle region of the Earth will limit these effects while investigating only the outer core.

3.3 Summary

We showed here that KM3NeT, with its two detectors ARCA and ORCA, is an experiment capable of performing both oscillation and absorption tomography of the Earth. Using a Monte Carlo based model to simulate the detector response, we were able to calculate the sensitivities for measurements of the $Z/A$ value of the outer core, which cannot be measured with seismic measurements, as well as of the Earth’s density profile. Despite the uncertainties related to neutrino tomography are still much bigger than those provided by geoscience, this new approach to explore the Earth interior can contribute to constrain future models of the planet.

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$^3$Weight percent’, meaning one percent of the total mass comes from hydrogen.
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