



Telescope Array 10-Year Monocular Spectrum Measurement

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Telescope Array (TA) is the largest cosmic ray detector in the northern hemisphere. We present a measurement of the cosmic ray energy spectrum for energies above $10^{17.5}$ eV using fluorescence telescopes in monocular mode. A novel weather classification scheme using machine learning was used to select data parts with good weather to ensure the quality of the fluorescence data. The data from the Black Rock Mesa (BR) and Long Ridge (LR) fluorescence telescope stations were analysed separately in monocular mode, with the calculated fluxes combined into a single spectrum. The 10-year monocular combined cosmic ray energy spectrum is in excellent agreement with previous measurements from the northern hemisphere. We present fits of the combined spectrum to a series of broken power law models. A thrice-broken power law is observed to be the best fit considering the Poisson deviance per degrees of freedom. The three breaks suggest an additional feature of the spectrum between the previously observed Ankle at $10^{18.7}$ eV and the GZK suppression at $10^{19.8}$ eV.

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1. Introduction

The Telescope Array (TA) Cosmic Ray Observatory is the largest cosmic ray detector operating in the northern hemisphere[1]. To better understand the nature of cosmic ray populations at the highest energies and the mechanisms that influence their propagation through the universe, we continue to measure the cosmic ray energy spectrum.

TA has passed 11 years of operation. As a followup to the previous TA FD monocular measurements with 3.5 years[2] and 7.5 years[3] of data and the latest measurement made by the TA SDs[4], we preform a spectrum calculation using nearly 10 years of data from the BR and LR FD stations in monocular mode. While the FDs operate with 11% duty cycle compared to the nearly 100% duty cycle of the SDs, the FDs observe the longitudinal development in the Extensive Air Showers (EAS) and have less hadronic-model dependency for event reconstruction.

2. Monocular Reconstruction

The analysis of events begins by calculating the shower-detector plane (SDP). The normal to the shower detector plane is determined by minimizing

$$\chi^2 = \sum_{i=1}^{N_{\text{good}}} (\hat{n} \cdot \hat{v}_i)^2 N_{\text{pe},i}$$
(1)

where \hat{n} is the normal to the SDP, \hat{v}_i the pointing direction of a good PMT, and $N_{\text{p.e.},i}$ is the number of photoelectrons is a good PMT.

Next, the geometry of the shower within the SDP is reconstructed by fitting the time vs. angle of observed signals. The fitted form is

$$t_i = t_0 + \frac{R_P}{c} \tan\left(\frac{\pi - \psi - \chi_i}{2}\right) \tag{2}$$

where t_i the time of a good tube, t_0 time when the first light can reach to the detector, R_P is the impact parameter, ψ is the angle of the shower with respect to the ground within the SDP, and χ_i is the angle viewed by a good tube from the horizon within the SDP. In this fit, R_P , ψ , and t_0 are the floating parameters and there is typically considerable correlation between them, requiring tracks with considerable angular extent to allow a good fit.

Once the geometry is determined, the shower profile is measured using the flux of photons observed by the detector, the atmospheric attenuation given the atmospheric properties, and the average fluorescence yield. The shower profile is fit to a Gaisser-Hillas[5] profile as a function of shower slant depth X:

$$N_{\rm ch} = N_{\rm max} \left(\frac{X - X_{\rm max}}{X_0 - X_{\rm max}}\right)^{\frac{X - X_{\rm max}}{\Lambda}} \exp\left(\frac{X_{\rm max} - X}{\Lambda}\right)$$
(3)

where N_{max} is the size of the shower at maximum, X_{max} is the slant depth of the maximum of the shower (measured in g/cm²), X_0 is an unobservable "beginning" of the shower and Λ is the decay parameter of the shower. The calorimetric energy of the shower is determined by integrating the energy deposited by the shower particles over the full extent of the shower

$$E_{\rm cal} = \int_{X_0}^{\infty} \left. \frac{dE_{\rm dep}}{dX} \right|_X dX \tag{4}$$

The missing energy, energy carried off by neutrinos and muons which do not deposit most of their energy in the atmosphere is added back in by a correction calculated from CORSIKA[6]

$$\frac{E_{\rm cal}}{E_0} = -0.5717 + 0.1416 \log_{10}(E_{\rm cal}/\text{eV}) - 0.003328 \left[\log_{10}(E_{\rm cal}/\text{eV})\right]^2$$
(5)

We use a new variable $R_{X_{\text{max}}}$ which measures the distance to the brightest part of the shower, as shown to the right. This variable is more relevant for defining the aperture than the impact parameter X_{max} which may refer to a part of the shower far away from where light is being emitted. See Figure 1.



Figure 1: Definition of $R_{X_{\text{max}}}$ and its relation to the impact parameter X_{max} .

3. Event Selection

From triggered events, a selection was made by various cuts. Cuts at the geometry reconstruction stage are given in Table 1. Cuts at the profile reconstruction phase are given in Table 2.

Good weather run segements were chose according to at novel neutral-network weather classification scheme. Using the fact that BRM and LR PMTs are DC connected and thus can provide a measurement of the background light level, one can make a movie of the night sky with one-minute long frames, and look for the passage of stars and/or clouds across the field-of-view. A neutralnetwork was designed and trained to recognize data parts which were either clear or cloudy. The data presented here comes only from the clear weather set. The good weather run time, including deadtime estimates is given in Table 3. Note that time when both detectors were operating and able to view events in tandem is also given.

With these cuts and taking data from good weather days results in the event energy distribution shown in Figure 2.

4. Aperture and Exposure Calculation

The instantaneous aperture of the detector of the detector was determined by Monte Carlo simulation. The showers thrown in the simulation were drawn from a shower library of CORSIKA-simulated proton and iron showers. The ratio of proton and iron showers at a particular energy were chosen according to the HiRes/MIA composition fit[7]. The energy was chosen according to the TA SD spectrum presented at the 2015 ICRC[8]. The detector response was simulated according

| Event Geometry Reconstruction Cuts | |
|---|--|
| Good PMT Fraction | $N_{\text{good PMT}}/N_{\text{PMT}} \ge 3.5\%$ |
| Number of good PMTs | $N_{\text{good PMT}} \ge 6$ |
| NPE per Degree | $N_{\rm pe}/\Delta\theta > 25~{ m deg}^{-1}$ |
| Pseudo-distance (inverse angular speed) | l > 1.5 km |
| SDP Angle (w.r.t. vertical) | $\zeta < 80^{\circ}$ |
| Impact Parameter | $R_P \ge 0.5 \text{ km}$ |
| Shower angle in SDP | $\psi < 130^{\circ}$ |
| Shower angle uncertainty | $\sigma_\psi < 36^\circ$ |
| Timing Fit | Successful |
| | $\chi^2/\mathrm{ndf} < 10$ |
| Track Length | $\Delta \theta > 10^{\circ}$ |
| | $\Delta \theta > 7^{\circ}$ (only in lower mirror) |
| Zenith Angle | $\theta_{\rm zen} < 70^{\circ}$ |
| Trigger Time w/i Frame | $t_0 < 25.6 \ \mu s$ |
| Crossing Time | $\Delta t < 6 \ \mu s$ for $R_P < 5 \ km$ |

Event Geometry Reconstruction Cuts

 Table 1: Event selection cuts applied at the geometry reconstruction phase.

| Profile Fit | Successful |
|-----------------------------|--|
| First Slant Depth Observed | $150 \text{ g/cm}^2 \le X_1 \le 1200 \text{ g/cm}^2$ |
| Slant Depth Extent Observed | $\Delta X \ge 150 \text{ g/cm}^2$ |
| Bracketing | $X_1 \leq X_{\max} \leq X_{\text{last}}$ |

Table 2: Event selection cuts applied at the profile reconstruction phase.

| | Operating | Duty | Ontime | Deadtime | Deadtime | Livetime |
|----------------------------|-----------|--------|---------|----------|----------|----------|
| | Nights | Cycle | [h] | [h] | Fraction | [h] |
| BRM | 1593 | 10.97% | 9288.19 | 574.16 | 6.18% | 8714.03 |
| LR | 1493 | 9.03% | 7647.41 | 555.79 | 7.26% | 7091.63 |
| $\text{BRM}\cap \text{LR}$ | 1378 | | 6876.40 | | 13.00% | 5982.47 |

Table 3: Runtime statistics for the two detectors and their overlap.

to daily calibration measurements and an average atmospheric model. The total exposures were fit to the form

$$\xi = p_1 \left(1 - \exp\left[-\frac{\log_{10}(E/eV - p_2)}{p_3} \right] \right)$$
(6)

for most of the aperture to smooth out the effect of MC statistics at the highest energies. The resulting exposures are shown in Figure 3

To verify the accuracy of the exposure calculation we present three comparisons of simulated data distributions to the same distribution in the actual data. We chose three that are most relevant to the exposure: $R_{X_{\text{max}}}$, ψ , and the SDP angle. These comparisons are shown in Figure 4



Figure 2: The energy distribution of events included in this analysis. BR and LR events are seen by each sight singly. BRM \cap LR events are seen by both sights and plotted as the geometric mean of the two measured energies.

5. Spectrum Calculation

Using the data and exposure presented, we calculate the following spectrum of UHECRs. We present it in Figure 5 along with the TA SD spectrum from the 2019 ICRC[4] for comparison.

We have fit the spectrum to a series of once, twice and thrice-broken power law models. The results of the fits are shown in Figure 6. For the two-break point fits, two separate solutions were found with very nearly the same chi-squared value, so we present both. The quality of the three break point fit indicates the presence of a feature between the Ankle and the GZK suppression, which has been observed in other analyses.

| | Single | Double A | Double B | Triple |
|----------------------------|------------------|----------------|------------------|----------------|
| $J_{18}/10^{-30}$ | 2.22 ± 0.01 | 2.22 ± 0.01 | 2.22 ± 0.01 | 2.22 ± 0.01 |
| $(eV m^2 sr s)^{-1}$ | | | | |
| γ_1 | -3.29 ± 0.01 | -3.29 ± 0.01 | -3.29 ± 0.01 | -3.29 ± 0.01 |
| $\log_{10}(E_1/\text{eV})$ | 18.68 ± 0.04 | 18.74 ± 0.03 | 18.72 ± 0.05 | 18.78 ± 0.04 |
| γ_2 | -2.79 ± 0.05 | -2.62 ± 0.06 | -2.70 ± 0.05 | -2.49 ± 0.15 |
| $\log_{10}(E_2/\text{eV})$ | | 19.46 ± 0.10 | 19.83 ± 0.04 | 19.20 ± 0.11 |
| γ_3 | | -3.8 ± 0.4 | -8.0 ± 2.7 | -3.04 ± 0.19 |
| $\log_{10}(E_3/\text{eV})$ | | | | 19.85 ± 0.01 |
| γ_3 | | | | -7.7 ± 2.4 |
| D/ndf | 43.5/26 | 22.69/24 | 23.29/24 | 17.74/22 |

Table 4: Parameters resulting from broken-powerlaw fits to the monocular spectrum.



Figure 3: The exposure of the three components of the detector: BRM and LR solely and BRM&LR together. The sum of the three exposures is also shown.

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Figure 4: Data/MC Comparisons to verify the quality of the exposure calculation. Shown are comparisons of $R_{X_{\text{max}}}$, ψ , and the SDP angle distributions.

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Figure 5: The spectrum of cosmic rays measured in this result. Shown for comparison with the TA SD spectrum presented at the 2019 ICRC.



Figure 6: Four power-law fits to the measured spectrum. The parameters of the fits are given in Table 4.

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