

## Results from the LHCf Run II in proton-proton collisions at $\sqrt{s} = 13$ TeV

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The LHCf experiment, at the Large Hadron Collider (LHC), consists of two small independent calorimeters placed 140 metres away, on opposite sides of the ATLAS interaction point (IP1). LHCf has the capability to measure zero-degree neutral particles, covering the pseudorapidity region above 8.4. By measuring the very-forward particle production rates at the highest energy possible at an accelerator, LHCf aims to improve our understanding of hadronic interactions in air-showers induced by ultra-high-energy cosmic rays in the atmosphere. This contribution will highlight recent results from Run II measurements with p-p collisions at 13 TeV. At first, the advantages of an ATLAS-LHCf combined analysis are discussed and a preliminary energy spectrum of very-forward photons produced in diffractive collisions as tagged by ATLAS is shown. The neutron energy spectrum measurements, for several pseudorapidity regions, are presented and compared with the predictions of various hadronic interaction models. From these measurements the average inelasticity of the collisions, which strongly affects the development of an air-shower, has also been extracted. Finally, the preliminary  $\pi^0$  Feynman-x and transverse momentum spectra, which affect the development of the electromagnetic component of an air-shower, are shown.

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## 1. Introduction

The indirect detection techniques of the Ultra-High-Energy Cosmic Rays (UHECRs, with energies  $\gtrsim 10^{18}$  eV) are based on the observation of the shower of secondary particles produced in the interaction of a cosmic ray with a nucleus of the atmosphere (the so-called “air-showers”). A clear interpretation of the data of ground based experiments is made difficult due to the large systematic uncertainty associated to the disagreement between different hadronic interaction models employed in the simulations of the air-showers. The tuning of these phenomenological models with experimental data is therefore fundamental to reduce the systematic uncertainty of UHECR measurements. The LHC-forward (LHCf [1]) experiment is designed to measure the very-forward neutral particles production at the Large Hadron Collider (LHC). The very-forward region covered by LHCf is crucial for the tuning of hadronic interaction models since it is the region where most of the energy flow of secondary particles is contained. After a brief description of the experimental apparatus in Section 2, the LHCf analysis results for photon, neutron and  $\pi^0$  production in proton-proton collisions at  $\sqrt{s} = 13$  TeV will be presented in Section 3.

## 2. The LHCf experiment

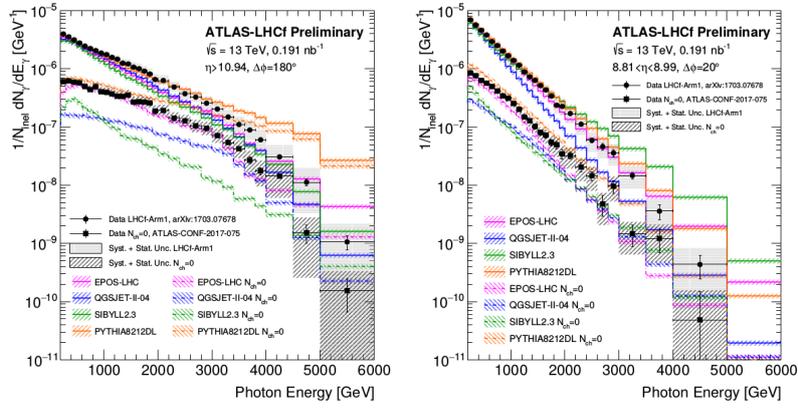
The LHCf experiment is composed of two independent detectors (named “Arm1” and “Arm2”) placed  $\sim 140$  meters away from the Interaction Point 1: The detectors are placed in the Target Absorber Neutral (TAN) instrumentation slot, where the beam pipe from the interaction point turns into two separate beam pipes: this position allows the measurement of particles emitted at a pseudorapidity  $\gtrsim 8.4$ , up to zero-degree. Since charged particles are deflected by the dipole magnets which bend the colliding beams into their own beam pipes, only neutral particles can be measured by the detectors. Each detector is made of two sampling and position sensitive calorimeters (called “towers” hereafter) which use  $\text{Gd}_2\text{SiO}_5$  (GSO) scintillators as active layers and tungsten as absorber. Each tower contains 4 couples of position sensitive layers to measure the transverse position: GSO scintillator bars are used in Arm1 while silicon microstrip detectors are used in Arm2. All calorimeters have a depth of 44 radiation lengths and 1.6 hadronic interaction lengths. The energy resolution is better than 3% for photons above 100 GeV [2] and  $\sim 40\%$  for neutrons. The position resolution of the silicon microstrip detectors for electromagnetic showers is  $40\mu\text{m}$  [3], while for GSO bars is  $200\mu\text{m}$  [2]. Position resolution for neutrons showers is  $\sim 1$  mm.

## 3. LHCf Run II results

The data set used in the following analyses has been taken during the LHCf operation in proton-proton collisions at  $\sqrt{s} = 13$  TeV in LHC fill #3855. A dedicated low-luminosity beam setup was used with a luminosity of  $0.4 \div 1.4 \times 10^{29} \text{ cm}^{-2} \text{ s}^{-1}$ . The photon and neutron analysis results shown in Section 3.1 and 3.2 used an integrated luminosity of  $0.19 \text{ nb}^{-1}$ , while the  $\pi^0$  analysis used a larger data sample of  $2.1 \text{ nb}^{-1}$  and  $0.8 \text{ nb}^{-1}$  for Type I and Type II  $\pi^0$  events, respectively (see Section 3.3 for the explanation of the event types).

### 3.1 LHCf-ATLAS combined photon analysis

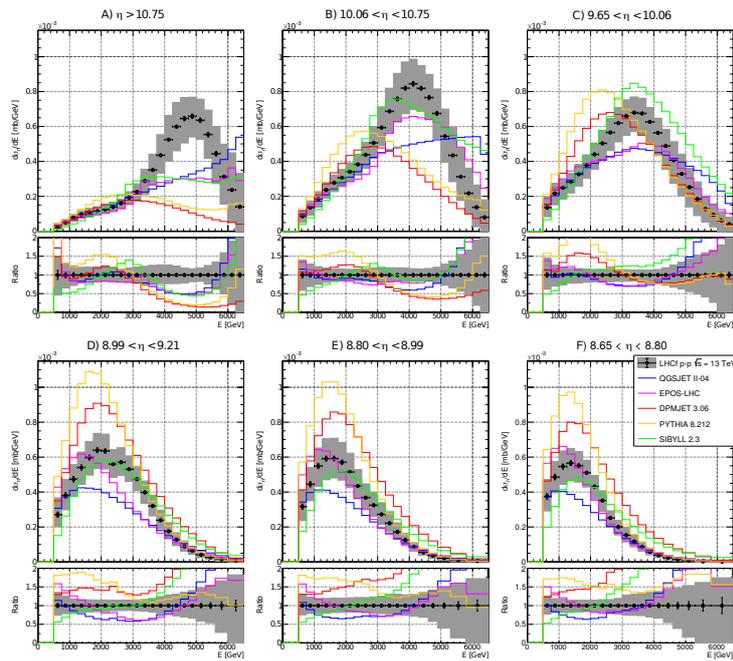
LHCf photon measurements in p-p collisions at 13 TeV have already been published [4]. However, the LHCf experiment alone does not have the capability to distinguish the type of collision occurred at the interaction point. Since 2013 LHCf and ATLAS [5] experiments exchange the trigger signals during LHCf dedicated runs, so it is possible to perform a combined analysis with the event information from both experiments. The number of tracks recorded by ATLAS detectors in the central region can be used to discriminate diffractive events from non-diffractive ones: selecting events with no charged particles in the region  $|\eta| < 2.5$  a pure sample of low-mass ( $M_X < 20$  GeV) diffraction events can be selected [6]. The preliminary results of the combined analysis are presented in Figure 1 [7], where the photon energy spectrum is shown for both the inclusive and the low-mass diffraction component. The diffraction spectrum of EPOS model has a good agreement with data in the  $\eta > 10.94$  region, while PYTHIA model has a better agreement for  $8.81 < \eta < 8.99$ .



**Figure 1:** Photon energy spectrum for pseudorapidity regions  $\eta > 10.94$  (left) and  $8.81 < \eta < 8.99$  (right) in p-p collisions at  $\sqrt{s} = 13$  TeV. Measured inclusive and low-mass diffractive spectra are represented as black circles and squares, respectively. Predictions from several hadronic interaction models for the inclusive and diffractive spectra are represented as solid and dashed lines, respectively. Hatched areas show statistical+systematic errors for data and statistical errors for simulations.

### 3.2 Neutron energy spectrum and average inelasticity

The neutron energy spectrum in six pseudorapidity bins is shown in Figure 2 [8]. All the distributions exhibit a peak structure, whose position progressively moves to lower energy as pseudorapidity decreases. As already noted [9], the largest disagreement between LHCf data and hadronic interaction models is found in the most forward region ( $\eta > 10.75$ ): the peak structure of the data is not reproduced by any model and all models underestimate the neutron yield by at least 20%. In the other pseudorapidity regions either SIBYLL or EPOS models have a better agreement with respect to other models. The distribution of the elasticity of the collision ( $k_n \equiv 2E_n/\sqrt{s}$ , where  $E_n$  is the energy of the neutron) is shown in the top panel of Figure 3, while the average inelasticity ( $1 - k_n$ ) is shown in the bottom panel. Even if no model well reproduces the elasticity distribution, the predicted average inelasticity of all models has a good agreement with the measured one.



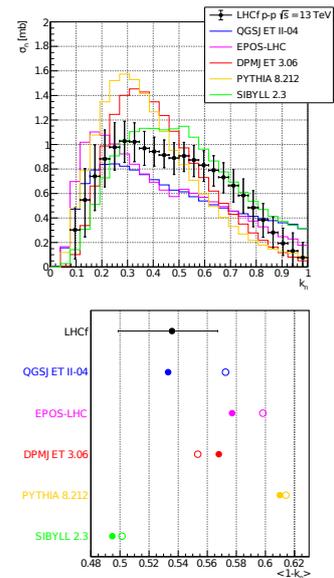
**Figure 2:** Measured neutron energy spectrum (black points) in different pseudorapidity bins compared with model predictions (coloured histograms). Grey area represents statistical+systematic uncertainty.

### 3.3 Neutral pion $p_T$ vs $X_F$ spectrum

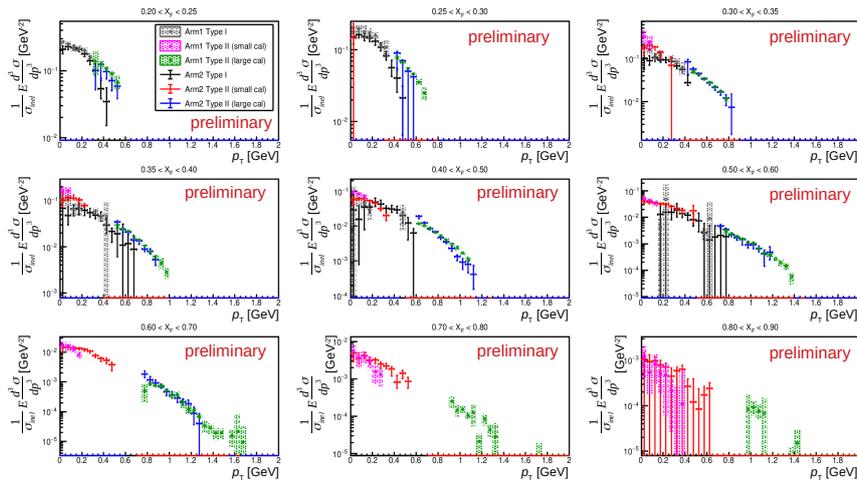
The  $\pi^0$  transverse momentum ( $p_T$ ) and Feynman-X ( $X_F \equiv 2p_Z/\sqrt{s}$ ) are reconstructed from the energy and impact position of the decay photons into the LHCf towers. Two typology of events can be recorded: “Type I” events where one photon hits each tower, and “Type II” events where both photons hit the same tower. The preliminary  $p_T$  spectra in several  $X_F$  bins are shown in Figure 4. Type I and Type II analyses provide results in good agreement in the regions where they overlap, and Arm1 and Arm2 data also appear to be consistent with each other. The consistency of the data between different detectors and event types permits to fully take advantage of the different  $p_T$  and  $X_F$  coverage of Arm1 and Arm2, and of Type I and Type II events. Indeed, thanks to the larger dimension of its large tower, Arm1 is able to reach a larger  $p_T$  with respect to Arm2. Conversely, Arm2 close the gaps of Arm1 acceptance for  $X_F < 0.6$  and extend the low- $p_T$  coverage for  $X_F > 0.6$  thanks to the larger dimension of its small tower.

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**Figure 3:** Top: neutron elasticity distribution for data (black points) and model predictions (coloured histograms). Bottom: average inelasticity for neutron events (full circles) and all event (open circles).



**Figure 4:** Transverse momentum spectra in different  $X_F$  bins measured with the Arm1 and Arm2 detectors of LHCf. Results from different detectors and  $\pi^0$  event types are shown with different colors. Errors bars represent the total estimated uncertainty (statistical + systematic).

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