

Simulations of charged hadron and charmed meson production in Pb+Pb collisions at $\sqrt{s_{NN}} = 5.02$ TeV with HYDJET++ generator

J. Bielčík,^{*a*} L. Bravina,^{*b*} G. Eyyubova,^{*c*} V. Korotkikh,^{*c*} I. Lokhtin,^{*c*} S. Petrushanko,^{*c*} A. Snigirev,^{*c*} J. Štorek^{*a*,*} and E. Zabrodin^{*b*,*c*}

^bDepartment of Physics, University of Oslo, PB 1048 Blindern, N-0316 Oslo, Norway

^c Skobeltsyn Institute of Nuclear Physics, M. V. Lomonosov Moscow State University, Leninskie gory, Moscow 119991, Russian Federation E-mail: jaroslav.bielcik@fjfi.cvut.cz, larissa.bravina@fys.uio.no, Gyulnara.Eyyubova@cern.ch, Vladimir.Korotkikh@cern.ch, igor.lokhtin@cern.ch, Serguei.Petrouchanko@cern.ch, Alexandre.Sniguirev@cern.ch, storejar@fjfi.cvut.cz, zabrodin@fys.uio.no

HYDJET++ is a Monte Carlo event generator merging parametrized soft part inspired by hydrodynamics with hard part containing jets. It has been successful to describe particle production in Pb+Pb collisions at $\sqrt{s_{NN}} = 2.76$ TeV energies. In this contribution, particle spectra and collective flow for the top LHC energy $\sqrt{s_{NN}} = 5.02$ TeV Pb+Pb collisions are presented for the first time. Specifically, the HYDJET++ model version 2.4 has been used to simulate spectra of charged particles, D^0 and J/ψ mesons and related v_2 and v_3 azimuthal flow harmonics. The particle spectra and flow haromines are studied in different centrality bins ranging from 0–10% up to 30–40% centrality in midrapidity region for charged particles and D^0 mesons and in forward rapidity in case of J/ψ mesons. The simulated results have been compared with the LHC data to tune HYDJET++ parameters.

*** The European Physical Society Conference on High Energy Physics (EPS-HEP2021), *** *** 26-30 July 2021 ***

*** Online conference, jointly organized by Universität Hamburg and the research center DESY ***

© Copyright owned by the author(s) under the terms of the Creative Commons Attribution-NonCommercial-NoDerivatives 4.0 International License (CC BY-NC-ND 4.0).

^a Faculty of Nuclear Sciences and Nuclear Engineering, Czech Technical University in Prague, Břehová 7, Prague, Czech Republic

^{*}Speaker

1. Introduction

In central ultra-relativistic heavy ion collisions, extreme energy densities are reached such that quark-gluon plasma (QGP), a novel phase of matter where quarks and gluons are deconfined, can be observed [1]. Different effects, e.g. modified yields of particle species with regards to a proton-proton collision, collective behavior of particles produced in the collision or jet quenching¹, can be recognized as phenomena of the QGP. Different physical models are used to predict the outcome of an experiment and correctly understand the main physical processes.

2. HYDJET++

HYDJET++ is a Monte Carlo (MC) generator for simulation of relativistic heavy ion collisions and merges hydro-inspired blast wave parameterization (soft) with jet quenching (hard) [2, 3]. In the soft part, hadrons are generated at chemical freeze-out hypersurface and thermal equilibrium is assumed during the thermal emission. The hard part is based on PYQUEN (PYthia QUENched) partonic energy loss model [4] which employs jet quenching in PYTHIA [5] generated jet events.

In former studies, production of charged hadrons and charmed mesons was successfully described in Au+Au collisions at $\sqrt{s_{NN}} = 200 \text{ GeV} [2, 6]$ and in Pb+Pb collisions at $\sqrt{s_{NN}} = 2.76 \text{ TeV}$ [7]. It has been found that different simulation parameters need to be used for a correct description of different particle species at different collisions energies. The temperature at thermal freeze-out T_{th} needs to be the same for charged hadrons and D mesons ($T_{\text{th}} = 105 \text{ MeV}$) and different for J/ψ meson ($T_{\text{th}} = 165 \text{ MeV}$) at $\sqrt{s_{NN}} = 2.76 \text{ TeV}$ energy [7]. The values of T_{th} at $\sqrt{s_{NN}} = 5.02 \text{ TeV}$ energy are studied in this proceedings using the most recent version 2.4 of the HYDJET++.



Figure 1: *Left:* Comparison of the HYDJET++ simulated elliptic and triangular flow coefficients of all charged hadrons h^{\pm} with the ATLAS experimental data [8] in 20–30% centrality bin. *Right:* Comparison of the HYDJET++ simulated $p_{\rm T}$ distribution histogram of the J/ψ meson yield to the ALICE experimental data [9].

¹Jet quenching is the modification of a jet caused by the QGP medium.

3. Charged hadrons h^{\pm}

The transverse momentum, $p_{\rm T}$, and pseudorapidity η distributions of charged hadrons² can be described well by HYDJET++ in 0–40% centrality range. HYDJET++ also correctly reproduces elliptic and triangular flow distributions v_2 , v_3 calculated by experiment adapted scalar product method in 10–30% semi-central events in 0 < $p_{\rm T}$ < 4 GeV/c region, however, it underestimates the data in 4 < $p_{\rm T}$ < 10 GeV/c region as shown in left graph in Fig. 1. Thermal freeze-out temperature $T_{\rm th} = 105$ MeV has been set for the simulation.

4. J/ψ meson

In the right graph in Fig. 1, HYDJET++ J/ψ transverse momentum p_T distribution is compared to the ALICE experimental data in 0–10% centrality bin. Parameter γ_c , which accounts for deviations of charm multiplicity from the complete thermal equilibrium value, has been set to $\gamma_c = 15$ for a correct description. One can see an underestimation of the experimental data in the $4 < p_T < 6$ GeV/c region. The mismatch can be slightly eliminated by tuning the maximal fluid flow transverse rapidity at thermal freeze-out ρ_{max} parameter. Nevertheless, no significant impact on elliptic flow of the J/ψ is observed in central collisions and $v_2^{J/\psi}$ is well described in the left graph in Fig. 2 up to $p_T < 6$ GeV/c. Thermal freeze-out temperature $T_{th} = 165$ MeV has been used for the J/ψ simulations.



Figure 2: Left: Comparison of the HYDJET++ simulated J/ψ elliptic flow coefficient v_2 with ALICE experimental data [10] in 0–10% centrality bin for two ρ_{max} values. Right: Comparison of the HYDJET++ simulated D^0 elliptic flow coefficient v_2 with the CMS experimental data [11] in 10–30% centrality bin.

5. D^0 mesons

Charm production enhancement parameter γ_c obtained from J/ψ simulations has been also used for the $D^0 p_T$ distribution resulting in a very good match between the HYDJET++ simulation and ALICE experimental data in 3 < p_T < 14 GeV/c region. Elliptic flow of the D^0 meson has been also studied in 10–30% centrality as can be seen in the right graph in Fig. 2. HYDJET++

²Inclusive charged hadrons h^{\pm} are π^{\pm} , K^{\pm} , protons p and antiprotons \bar{p} .

generally follows the trend of the experimental data but overestimation in $4 < p_T < 6$ GeV/c region is observed.

The same thermal freeze-out temperature $T_{\text{th}} = 105$ MeV as for the charged hadrons has been used for HYDJET++ D^0 simulations.

6. Conclusion

For all the studied distributions, a good description of the LHC data has been achieved by the HYDJET++ model. It appears that raising Pb+Pb collision energy from $\sqrt{s_{NN}} = 2.76$ TeV to $\sqrt{s_{NN}} = 5.02$ TeV does not have a significant impact on the thermal freeze-out temperature T_{th} which is the same value for h^{\pm} and D^0 meson and different value for J/ψ . Correct description of charm meson spectra has been achieved by tuning charm enhancement parameter $\gamma_c = 15$ and the maximal fluid flow transverse rapidity at thermal freeze-out ρ_{max} parameter is found to have only a small impact on the J/ψ elliptic flow in central events.

7. Acknowledgement

The work was supported from European Regional Development Fund-Project "Center of Advanced Applied Science" No. CZ.02.1.01/0.0/0.0/16_019/0000778 and by the grant LTT18002 of Ministry of Education, Youth and Sports of the Czech Republic.

References

- [1] R. Pasechnik and M. Šumbera, Universe 3 (2017) 7
- [2] I. P. Lokhtin et al., Comput. Phys. Commun., 180 (2009) 779
- [3] http://lokhtin.web.cern.ch/lokhtin/hydjet++
- [4] I. P. Lokhtin, A. M. Snigirev, Eur. Phys. J. C, 45 (2006) 211
- [5] T. Sjöstrand et al., J. High Energy Phys., 05 (2006) 26
- [6] I. P. Lokhtin et al., J. Phys. Conf. Ser., 270 (2011) 012060
- [7] I. P. Lokhtin et al., J. Phys. G: Nucl. Part. Phys., 43 (2016) 125104
- [8] M. Aaboud et al., Eur. Phys. J. C, 78 (2018) 997
- [9] S. Acharya et al., J. High Energy Phys., 2 (2020) 41
- [10] S. Acharya et al., J. High Energy Phys., 10 (2020) 141
- [11] A. M. Sirunyan et al., Phys. Rev. Lett., 120 (2018) 202301