

# Measurement of the Z boson production in association with at least two b jets in pp collisions at $\sqrt{s} = 13$ TeV

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The cross sections for Z boson production in association with at least two b jets are measured as a function of various kinematic variables, using pp collisions at  $\sqrt{s} = 13$  TeV recorded by the CMS experiment, corresponding to an integrated luminosity of  $137 \text{ fb}^{-1}$ . Decays of the Z boson to electrons or muons are considered with leading (sub-leading) lepton transverse momentum  $p_T > 35$  (25) GeV and pseudorapidity  $|\eta| < 2.4$ , and the invariant mass within 71 and 111 GeV. Jets are selected with  $p_T > 30$  GeV and  $|\eta| < 2.4$ . The results are compared to various QCD calculations.

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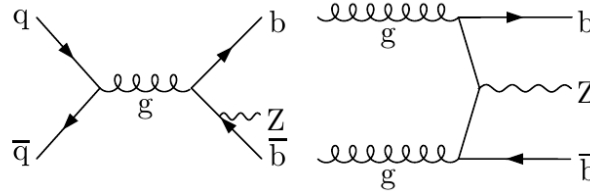
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## 1. Introduction

Studies of Z boson production in association with jets from b quarks provide important tests of perturbative quantum chromodynamics. These processes are also important since they are major backgrounds for many physics processes, including the Higgs boson production in association with a Z boson, ZH ( $H \rightarrow b\bar{b}, c\bar{c}$ ) and searches for new physics phenomena. Figure 1 shows the tree level diagrams of the Z+2b process.



**Figure 1:** Examples of Feynman diagrams of Z+2b production.

## 2. Overview of measurements

The cross section is measured in the leptonic decay channel of Z bosons ( $Z \rightarrow ee, Z \rightarrow \mu\mu$ ) using pp collision data at  $\sqrt{s} = 13$  TeV [1] collected by the CMS experiment [2] during the 2016 ( $35.9 \text{ fb}^{-1}$ ), 2017 ( $41.5 \text{ fb}^{-1}$ ), and 2018 ( $59.7 \text{ fb}^{-1}$ ) data-taking periods, corresponding to an integrated luminosity of  $137 \text{ fb}^{-1}$ .

The collision events are required to pass single lepton triggers. They must have at least one electron (muon) candidate with the minimum transverse momentum  $p_T$  of 27, 32, and 32 GeV (24, 27, and 24 GeV) during 2016, 2017, and 2018 data-taking periods.

Several Monte Carlo (MC) event generators are used to simulate signal process. The Drell–Yan (DY) process is simulated by MADGRAPH5\_aMC@NLO (denoted as MG5\_aMC) [3] calculated at next-to-leading order (NLO). The second sample is simulated by MG5\_aMC calculated at leading order (LO), and the third one is simulated by SHERPA [4]. The MG5\_aMC predictions are with different PDF settings and underlying event tune corresponding to 2016 and 2017–2018 data-taking periods. The 2016 setting uses NNPDF 3.0 PDF set [5] and tune CUETP8M1 [6], while the 2017–2018 setting uses NNPDF 3.1 PDF set [5] and tune CP5 [7]. Background contributions are from top quark-antiquark ( $t\bar{t}$ ), single top quark (s- and t-channel, and  $tW$ ), and diboson (WW, WZ, ZZ) processes.

Electrons and muon are required to be isolated and selected with the kinematic criteria: leading (sub-leading) lepton transverse momentum  $p_T > 35$  (25) GeV and pseudorapidity  $|\eta| < 2.4$ . In the muon channel, Rochester correction [8] is applied to MC samples and data to mitigate the biases occurred during muon  $p_T$  reconstruction. The correction is done by applying the Rochester correction scale factor to the muon  $p_T$ .

Z boson candidates are reconstructed from two isolated, oppositely charged leptons, selected as described before. Afterwards, Z boson candidates are selected with the dilepton invariant mass range  $71 < M_{ll} < 111$  GeV.

Jets are reconstructed using the anti- $k_T$  clustering algorithm [9] with a distance parameter of 0.4. Jets are required to have  $p_T > 30$  GeV and  $|\eta| < 2.4$ . If jets overlap with  $p_T > 25$  GeV leptons in a cone,  $\Delta R < 0.4$ , they are discarded. Jets from pileup events are excluded.

Missing transverse energy ( $E_T^{\text{miss}}$ ) is used in this analysis to reduce the  $t\bar{t}$  background by requiring  $E_T^{\text{miss}} > 50$  GeV.  $\vec{E}_T^{\text{miss}}$  is defined as the vector sum of the transverse momentum of all observables reconstructed in the event,  $\vec{E}_T^{\text{miss}} = -\sum \vec{p}_T, i$ .

The deep combined secondary vertex (DeepCSV) [10] algorithm is used in this analysis to identify  $b$  jets from light or  $c$  jets. Tight tagging requirement (50%  $b$  quark tagging efficiency with 0.1% misidentification rate for jets originating from gluons or  $u, d, s$  quarks) is applied to reduce  $DY+X$  backgrounds, where  $X = \text{light or } c \text{ jets}$ .  $b$  jets are tagged based on unique properties of  $B$  hadrons inside the jets, such as relatively long lifetime, which causes the secondary vertex displaced from the primary vertex, and the relatively large  $B$  hadron invariant mass.

### 3. Results

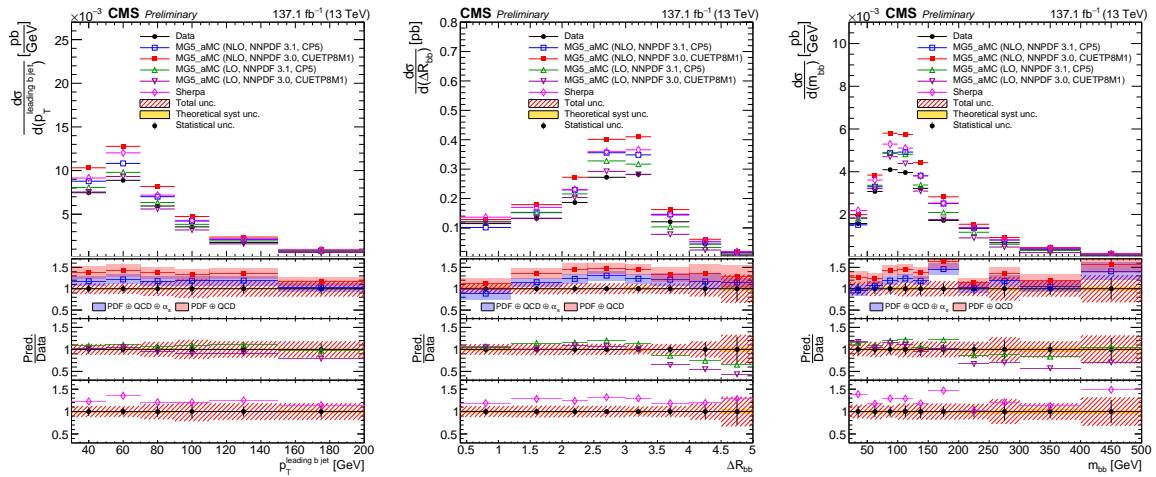
The inclusive cross section for  $Z + \geq 2$   $b$  jets for the combination of channels are listed in Table 1. The measured integrated cross section is  $0.65 \pm 0.03$  (stat)  $\pm 0.07$  (syst)  $\pm 0.02$  (theo) pb. The measured cross section is compared with several predictions: `MG5_aMC` calculated at NLO, `MG5_aMC` calculated at LO and SHERPA. The measured integrated cross section is in good agreement with `MG5_aMC` at LO prediction, while `MG5_aMC` at NLO and SHERPA predictions overestimate the measured cross section value.

Measured	$0.65 \pm 0.03$ (stat) $\pm 0.07$ (syst) $\pm 0.02$ (theo)
<code>MG5_aMC</code> (LO) (NNPDF 3.1, CP5)	0.71
<code>MG5_aMC</code> (LO) (NNPDF 3.0, CUETP8M1)	0.63
<code>MG5_aMC</code> (NLO) (NNPDF 3.1, CP5)	$0.77 \pm 0.07$
<code>MG5_aMC</code> (NLO) (NNPDF 3.0, CUETP8M1)	$0.90 \pm 0.09$
SHERPA	0.84

**Table 1:** Measured and predicted integrated cross sections (in pb) for  $Z + \geq 2$   $b$  jets final states.

The observed distributions at reconstruction-level are unfolded to particle-level by applying corrections for detector effects represented by a response matrix. The response matrix is constructed from reconstruction-level leptons and  $b$  jet objects which are spatially matched to the corresponding particle-level objects within  $\Delta R < 0.3$ .

The unfolded differential cross section distributions as functions of leading  $b$  jet  $p_T$ , angular separation of two  $b$  jets ( $\Delta R_{bb}$ ), and the invariant mass of the two  $b$  jets ( $m_{bb}$ ) are shown in Fig. 2. From the leading  $b$  jet  $p_T$  distribution in Fig. 2, we can see that `MG5_aMC` at LO predictions are in good agreement with data, while `MG5_aMC` at NLO and SHERPA predictions overestimate the data from 20%–50%. In the  $\Delta R_{bb}$  distribution in Fig. 2, `MG5_aMC` at NLO and SHERPA predictions overestimate the data at high  $\Delta R_{bb}$  region, while `MG5_aMC` at LO predictions have the opposite trend. For the  $m_{bb}$  distribution in Fig. 2, none of the predictions agrees with data.



**Figure 2:** Differential cross section distribution as functions of the leading  $b$  jet  $p_T$  (left),  $\Delta R_{bb}$  (middle), and  $m_{bb}$  (right). The lower panels show the ratio of each MC prediction and data.

#### 4. Summary

We measured the integrated cross section for  $Z + \geq 2$   $b$  jets production, and predictions of three different MC simulations (MG5\_aMC at LO, MG5\_aMC at NLO, SHERPA) have been used to compare with data. The MG5\_aMC at LO prediction gives a good description of the measured integrated cross section. For the differential cross sections, the shapes in the data are well described by the predictions from MG5\_aMC at NLO and SHERPA.

#### References

- [1] CMS Collaboration, “Measurement of  $Z+b$  jets cross section in proton-proton collisions at  $\sqrt{s} = 13$  TeV”, CMS-PAS-SMP-20-015.
- [2] CMS Collaboration, “The CMS experiment at the CERN LHC”, *JINST* **3** (2008) S08004.
- [3] J. Alwall et al., “The automated computation of tree-level and next-to-leading order differential cross sections, and their matching to parton shower simulations”, *JHEP* **07** (2014) 079 [arXiv:1405.0301].
- [4] E. Bothmann et al., “Event generation with Sherpa 2.2”, *SciPost Phys* **7** (2019) 034 [arXiv:1905.09127].
- [5] The NNPDF collaboration., Ball, R.D., Bertone, V. et al., “Parton distributions for the LHC run II”, *JHEP* **04** (2015) 040 [arXiv:1410.8849].
- [6] CMS Collaboration, “Event generator tunes obtained from underlying event and multiparton scattering measurements”, *Eur. Phys. J. C* **76** (2016) 155 [arXiv:1512.00815].
- [7] CMS Collaboration, “Extraction and validation of a new set of CMS PYTHIA8 tunes from underlying-event measurements”, *Eur. Phys. J. C* **80** (2020) 4 [arXiv:1903.12179].

- [8] A. Bodek et al., “Extracting muon momentum scale corrections for hadron collider experiments”, *Eur. Phys. J. C* **72** (2012) 2194 [arXiv:1208.3710].
- [9] M. Cacciari, G. P. Salam, and G. Soyez, “The anti- $k_t$  jet clustering algorithm”, *JHEP* **04** (2008) 063 [arXiv:0802.1189].
- [10] CMS Collaboration, “Identification of heavy-flavour jets with the CMS detector in  $pp$  collisions at 13 TeV”, *JINST* **13** (2018) P05011 [arXiv:1712.07158].