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Charmless B decay measurements at Belle

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In this report, we summarize the recent charmless *B* decay measurements at Belle. The studies are based on the Belle data sample of 711 fb⁻¹ or 121 fb⁻¹ collected at $\Upsilon(4S)$ or $\Upsilon(5S)$ resonance at the KEKB collider. Results of several decay modes are presented. In addition to their branching measurements, the structure in the two-body invariant mass are also investigated for some of the decay modes.

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1. Introduction

Charmless *B* decays are suppressed in Standard Model (SM), and are also sensitive to physics beyond the Standard Model (BSM) within the loop of penguin amplitude. Precise measurements on them could be a good sensitivity test against the prediction by SM. The main experimental challenge is the signal rate which is about 10⁵ times smaller than the $e^+e^- \rightarrow q\bar{q}$ (q = u, d, s, c) continuum processes. Reduction of combinatorial background is hence critical. Using a Belle data sample of 711 fb⁻¹ or 121 fb⁻¹ collected at $\Upsilon(4S)$ or $\Upsilon(5S)$ resonance with the Belle detector [1] at the KEKB asymmetric-energy e^+e^- collider [2], we report the studies of the following *B* and B_s^0 decay modes: $B^0 \rightarrow p\bar{p}\pi^+\pi^-, B^+ \rightarrow p\bar{p}\pi^+\pi^0$ [3], $B^+ \rightarrow K^+K^-\pi^+, B^+ \rightarrow \pi^+\pi^0\pi^0, B_s^0 \rightarrow \eta' X_{s\bar{s}}$ [4], $B_s^0 \rightarrow \eta'\eta$ [5], and $B_s^0 \rightarrow \eta' K_S^{0\dagger}$. Their signal yields (N_{sig}) are measured by one or multi-dimensional extended unbinned likelihood fit on data with different variables, and branching fractions are estimated by $\mathcal{B} = \frac{N_{sig}}{\epsilon \times N_B}$, where ϵ is the signal reconstruction efficiency and N_B is the number of *B* events (772M for B^+ or B^0 , 16.6M for B_s^0).

2. $B^0 \rightarrow p\bar{p}\pi^+\pi^-$ and $B^+ \rightarrow p\bar{p}\pi^+\pi^0$

Baryonic *B* decays have various interesting features, such as the enhancement in the di-baryon low mass threshold [6], and the different angular distributions from different modes e.g. between $B^+ \to p\bar{p}K^+$ and $B^+ \to p\bar{p}\pi^+$ [7]. Signal *B* candidate is identified by the energy difference $\Delta E \equiv E_B - E_{\text{beam}}$ and the beam-energy-constrained mass $M_{\text{bc}} \equiv \sqrt{E_{\text{beam}}^2/c^4 - |p_B/c|^2}$, where E_{beam} is the beam energy, and p_B and E_B are the momentum and energy of the reconstructed *B* meson, respectively. We use 2D fit with ΔE and M_{bc} to extract N_{sig} , and obtain $\mathcal{B}(B^0 \to p\bar{p}\pi^+\pi^-) = (0.83 \pm 0.17(\text{stat.}) \pm 0.17(\text{syst.})) \times 10^{-6}$ and $\mathcal{B}(B^+ \to p\bar{p}\pi^+\pi^0) = (4.58 \pm 1.17(\text{stat.}) \pm 0.67(\text{syst.})) \times 10^{-6}$. The total measured $\mathcal{B}(B^+ \to p\bar{p}\pi^+\pi^0)$ is about a factor of 10 smaller than the predicted $\mathcal{B}(B^+ \to p\bar{p}\rho^+)$ from a theoretical calculation by generalized factorization method [8]. Figure 1 shows the $M_{\pi\pi}$ distribution. A χ^2 fit is perform on $M_{\pi^+\pi^0}$ and we obtain 86 ± 41 events for $B^+ \to p\bar{p}\rho^+$. Table 1 shows the signal yields in $M_{p\bar{p}}$ bins. Branching fraction of $B^0 \to p\bar{p}\pi^+\pi^-$ in the threshold enhancement region (the lowest bin) is estimated as $(0.35 \pm 0.13(\text{stat.}) \pm 0.07(\text{syst.})) \times 10^{-6}$, which is consistent with the LHCb result [9].



Figure 1: $M_{\pi^+\pi^-}$ (left) and $M_{\pi^+\pi^0}$ (right) distributions of $B^0 \to p\bar{p}\pi^+\pi^-$ and $B^+ \to p\bar{p}\pi^+\pi^0$, respectively.

[†]Throughout this paper, inclusion of charge-conjugate decay modes is always implied.

Table 1: Fitted yields of $B^0 \to p\bar{p}\pi^+\pi^-$ (0.6 GeV/ $c^2 < M_{\pi^+\pi^-} < 1.22$ GeV/ c^2) and $B^+ \to p\bar{p}\pi^+\pi^0$ ($M_{\pi^+\pi^0} < 1.3$ GeV/ c^2) in $M_{p\bar{p}}$ bins.

$M_{p\bar{p}} (\text{GeV}/c^2)$	Yield of $B^0 \to p \bar{p} \pi^+ \pi^-$	Yield of $B^+ \to p \bar{p} \pi^+ \pi^0$
$M_{p\bar{p}} < 2.85$	$26.1^{+10.0}_{-9.1}$	$133.5^{+26.6}_{-25.2}$
$2.85 < M_{p\bar{p}} < 3.128$	$19.6^{+10.2}_{-9.3}$	$12.3_{-9.7}^{+10.3}$
$3.128 < M_{p\bar{p}}$	$29.1_{-13.1}^{+16.2}$	$-3.8^{+15.1}_{-13.8}$

3. $B^+ \rightarrow K^+ K^- \pi^+$

Compared to previous measurement by Belle [10], $B^+ \to K^+ K^- \pi^+$ result is updated with a re-optimized binning to study the property of the structure and localized \mathcal{A}_{CP} at low $M_{K^+K^-}$ region which were also observed in BaBar [11] and LHCb [12–14]. Signal yields and \mathcal{A}_{CP} are extracted by using 2D fit with ΔE and M_{bc} within each $M_{K^+K^-}$ bins, and Figure 2 shows the results. The structure at $M_{K^+K^-} < 1.1 \text{ GeV}/c^2$ has an \mathcal{A}_{CP} of $-0.90 \pm 0.17(\text{stat.}) \pm 0.03(\text{syst.})$ with a significance of 4.8σ . Helicity angle (θ_{hel} , defined as the angle between B^+ and K^+ in the K^+K^- rest frame) for signal events within $M_{K^+K^-} < 1.1 \text{ GeV}/c^2$ is shown in Figure 3. The distribution is consistent with a coherent sum of spin-0 and spin-1 the most.



Figure 2: Differential branching fraction (left) and \mathcal{A}_{CP} (right) distributions as a function of $M_{K^+K^-}$ for $B^+ \to K^+K^-\pi^+$.

4. $B^+ \to \pi^+ \pi^0 \pi^0$

The major challenge in the $B^+ \to \pi^+ \pi^0 \pi^0$ measurement is the shower leakage [15] due to two π^0 in the reconstruction, and the correlation between energy and other variables. e.g. between ΔE and $M_{\pi\pi}$. To handle those effects, we require the momentum to be greater 0.5 GeV/ c^2 for all π^0 candidates. By a 3D fit with ΔE , $M_{\rm bc}$, and a Neural-Network [16] output discriminant for continuum suppression [17], we obtain inclusive $\mathcal{B}(B^+ \to \pi^+ \pi^0 \pi^0) = (19.0 \pm 1.5(\text{stat.}) \pm 1.4(\text{syst.})) \times 10^{-6}$ and $\mathcal{A}_{CP} = (9.2 \pm 6.8(\text{stat.}) \pm 0.5(\text{syst.}))\%$. We use the ${}_s\mathcal{P}lot$ technique [18] to isolate signal on $M_{\pi\pi}$ distribution, and perform a 2D binned fit on the histogram to extract the signal model composition as shown in Figure 4. In addition to the $B^+ \to \rho(770)^+\pi^0$ structure at low $M_{\pi^+\pi^0,\min}$



Figure 3: The helicity angle distribution with applying efficiency correction and comparisons to different models, where the LHCb model is from Ref. [13].

region[†], and we also observe a new structure at $M_{\pi^0\pi^0}$ region, which is modeled by an incoherent sum of $f_0(980)$, $f_2(1270)$, and $f_0(500)$. A combined branching fraction for this $\pi^0\pi^0$ structure is measured as $(6.9 \pm 0.9(\text{stat.}) \pm 0.6(\text{syst.})) \times 10^{-6}$, which has a significance of 9.2 σ . A large \mathcal{R}_{CP} is seen at $M_{\pi^0\pi^0} \sim 1.4 \text{ GeV}/c^2$ as shown in Figure 5.



Figure 4: Projection of the fit result to ${}_{s}Weights M_{\pi^{+}\pi^{0},\min}$ -vs- $M_{\pi^{0}\pi^{0}}$ histogram.



Figure 5: ${}_{s}Weights \mathcal{A}_{CP}$ as a function of $M_{\pi^{0}\pi^{0}}$ for $M_{\pi^{+}\pi^{0},\min} > 1.9 \text{ GeV}/c^{2}$. The first few bins are combined due to low number of events.

[‡] $M_{\pi^+\pi^0,\min}$ refers to the smaller of two $M_{\pi^+\pi^0}$ values in a reconstructed *B* candidate.

5. $B_s^0 \to \eta' X_{s\bar{s}}, B_s^0 \to \eta' \eta$, and $B_s^0 \to \eta' K_s^0$

As *B* decays with η' in the final state have been observed firstly by CLEO [19, 20], we have found some special properties in this particle and decays involving it. η' mass is higher than the expectation from symmetry considerations [21]. Measurements of $\mathcal{B}(B \to \eta' X_s)$ [23–25] also show unexpected enhancement compared with SM prediction [22]. Any new observation on decays with η' could provide further information to understand its property.

We report the first measurement on $B_s^0 \to \eta' X_{s\bar{s}}$ based on a semi-inclusive method [4]. Simulation of $X_{s\bar{s}}$ fragmentation is performed with PYTHIA 6 [26] with a flat mass distribution. $X_{s\bar{s}}$ candidates are reconstructed with two kaons $(K^+K^- \text{ or } K^\pm K_S^0 \text{ with } K_S^0 \to \pi^+\pi^-)$ and up to four pions with at most one π^0 . η' candidates are reconstructed with $\pi^+\pi^-\eta$ and $\eta \to \gamma\gamma$. N_{sig} is extracted by a 1D fit on M_{bc} with $-0.12 \text{ GeV} < \Delta E < 0.05 \text{ GeV}$ in $M_{X_{s\bar{s}}}$ bins. As none of the bins shows significant yield, we set an upper limit on $\mathcal{B}(B_s^0 \to \eta' X_{s\bar{s}})$ as 1.4×10^{-3} at 90% confidence level (C.L.).

Branching fraction and *CP* asymmetry of $B_s^0 \to \eta' \eta$ decay could be affected by various BSM scenarios [28]. Along with the results of other $B_{d,s}^0 \to \eta \eta, \eta' \eta, \eta' \eta'$ modes, measurement on $B_s^0 \to \eta' \eta$ is helpful to extract *CP*-violating parameters from SU(3)/U(3) symmetry [29]. N_{sig} of $B_s^0 \to \eta' \eta$ is extracted by 3D fit with ΔE , M_{bc} , and $M_{\eta'}$. We obtain 2.7 ± 2.5 events and set upper limits of $f_s \times \mathcal{B}(B_s^0 \to \eta' \eta)$ and $\mathcal{B}(B_s^0 \to \eta' \eta)$ as 1.3×10^{-5} and 6.5×10^{-5} at 90% C.L., respectively, where f_s is the fraction of $B_s^{(*)0} \bar{B}_s^{(*)0}$ in $b\bar{b}$ events and its world average is 0.201 ± 0.031 [27].

 $B_s^0 \to \eta' K_S^0$ decay contains contributions from gluonic and electroweak penguin amplitudes, such that it is sensitive to BSM physics [28] which can affect both decay rate and *CP* asymmetries. N_{sig} of $B_s^0 \to \eta' K_S^0$ is extracted by 3D fit with ΔE , M_{bc} , and $M_{\eta'}$. We obtain -3.21 ± 1.85 events and set upper limits of $f_s \times \mathcal{B}(B_s^0 \to \eta' K_S^0)$ and $\mathcal{B}(B_s^0 \to \eta' K_S^0)$ as 1.64×10^{-5} and 8.16×10^{-5} at 90% C.L., respectively.

6. Summary

We report the results of several charmless *B* decays using Belle data collected at $\Upsilon(4S)$ or $\Upsilon(5S)$ resonance. In addition to branching fraction measurement, we also look into the distribution of two-body invariant mass of $B^0 \to p\bar{p}\pi^+\pi^-$, $B^+ \to p\bar{p}\pi^+\pi^0$, $B^+ \to K^+K^-\pi^+$, and $B^+ \to \pi^+\pi^0\pi^0$ to study their decay structure. We do not observe significant signal yield for $B_s^0 \to \eta' X_{s\bar{s}}$, $B_s^0 \to \eta' \eta$, and $B_s^0 \to \eta' K_s^0$, and upper limits on branching fraction are estimated at 90% C.L.. In near future, larger data set from Belle II [30] can further improve the measurement on these decay modes.

References

- A. Abashian *et al.* (Belle Collaboration), Nucl. Instrum. Methods Phys. Res., Sect. A **479**, 117 (2002); also see detector section in J. Brodzicka *et al.*, Prog. Theor. Exp. Phys. **2012**, 04D001 (2012).
- [2] S. Kurokawa and E. Kikutani, Nucl. Instrum. Methods Phys. Res., Sect. A 499, 1 (2003), and other papers included in this volume; T.Abe *et al.*, Prog. Theor. Exp. Phys. 2013, 03A001 (2013) and references therein.

- [3] K. Chu et al. (Belle Collaboration), Phys. Rev. D 101, 052012 (2020).
- [4] S. Dubey et al. (Belle Collaboration), Phys. Rev. D 104, 012007 (2021).
- [5] N.K. Nisar et al. (Belle Collaboration), Phys. Rev. D 104, L031101 (2021).
- [6] K. Abe et al. (Belle Collaboration), Phys. Rev. Lett. 88, 181803 (2002).
- [7] J. Wei et al. (Belle Collaboration), Phys. Lett. B 659, 80 (2008).
- [8] C. Q. Geng, Y. K. Hsiao, and J. N. Ng, Phys. Rev. D 75, 094013 (2007).
- [9] R. Aaij et al. (LHCb Collaboration), Phys. Rev. D 96, 051103 (2017).
- [10] C.-L. Hsu et al. (Belle Collaboration), Phys. Red. D 96, 031101(R) (2017).
- [11] B. Aubert et al. (BaBar Collaboration), Phys. Rev. Lett. 99, 221801 (2017).
- [12] R. Aaij et al. (LHCb Collaboration), Phys. Rev. Lett. 112, 011801 (2014).
- [13] R. Aaij et al. (LHCb Collaboration), Phys. Rev. Lett. 123, 231802 (2019).
- [14] R. Aaij et al. (LHCb Collaboration), Phys. Rev. Lett. 112, 011801 (2014).
- [15] K. Miyabayashi, Nucl. Instrum. Methods Phys. Res., Sect. A 494, 298 (2002).
- [16] M. Feindt and U. Kerzel, Nucl. Instrum. Methods Phys. Res., Sect. A 559, 190 (2006).
- [17] G. C. Fox and S. Wolfram, Phys. Rev. Lett. 41, 1581 (1978). The modified moments used in this paper are described in S. H. Lee *et al.* (Belle Collaboration), Phys. Rev. Lett. 91, 261801 (2003).
- [18] M. Pivk and F. R. Le Diberder, Nucl. Instrum. Methods Phys. Res., Sect. A 555, 356 (2005).
- [19] T.E. Browder et al. (CLEO Collaboration), Phys. Rev. Lett., 81, 1786, 1998.
- [20] G. Bonvicini et al. (CLEO Collaboration), Phys. Rev. D, 68, 011101, 2003.
- [21] K. Ottnad and C. Urbach (ETM Collaboration), Phys. Rev. D 97, 054508 (2018).
- [22] A. Datta, X.-G. He, and S. Pakvasa, Phys. Lett. B 419, 369 (1998).
- [23] G. Bonvicini et al. (CLEO Collaboration), Phys. Rev. D 68, 011101 (2003).
- [24] B. Aubert et al. (BABAR Collaboration), Phys. Rev. Lett. 93, 061801 (2004).
- [25] K. Nishimura et al. (Belle Collaboration), Phys. Rev. Lett. 105, 191803 (2010).
- [26] T. Sjostrand, S. Mrenna, and P. Skands, J. High Energy Phys. 06 026 (2006).
- [27] P. A. Zyla et al. (Particle Data Group), Prog. Theor. Exp. Phys. 2020, 083C01 (2020).
- [28] E. Kou et al., Prog. Theor. Exp. Phys. 123C01 (2019).
- [29] Y.-K. Hsiao, C.-F. Chang, and X.-G. He, Phys. Rev. D 93, 114002 (2016).
- [30] T. Abe et al. (Belle II Collaboration), arXiv:1011.0352 [physics.ins-det], (2010).