

# Transverse momentum spectra and Nuclear Modification factor in Xe-Xe collisions at 5.44 TeV under HYDJET++ framework

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Transverse momentum  $p_T$  spectra of charged hadrons at mid-pseudorapidity in deformed Xe-Xe collisions at 5.44 TeV center-of-mass energy under the Monte Carlo HYDJET++ model (HYDrodynamics plus JETs) framework is reported. Results have been presented in  $|\eta| < 0.8$  kinematic and (0-60)% centrality range. The nuclear modification factor in Xe-Xe collisions is calculated for most central, semi-central, semi-peripheral, and most peripheral collision centralities. Transverse momentum spectra and nuclear modification factor  $R_{AA}$  show strong  $p_T$ , and centrality dependence. Average transverse momentum  $\langle p_T \rangle$  as a function of collision centrality is presented. The results have been compared with ALICE experimental data.

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## 1. Introduction

High transverse momentum quarks and gluons (jets), created from early stage hard scatterings are very useful probes of such highly excited nuclear matter [1]. Transverse momentum  $p_T$  spectra of charged hadrons is a suitable observable for studying these hard partonic jets [2]. One of the major signals of QGP formation is the strong suppression of these high- $p_T$  particles (called as jet quenching) observed in heavy ion collisions at RHIC and LHC. The suppression of particle spectra in A+A collisions is measured by the nuclear modification factor  $R_{AA}$ .

## 2. Model Formalism

HYDJET++ is a Monte Carlo model of relativistic heavy ion collisions which simultaneously simulates two independent components: soft hydro-type (also called as hydro part) state and the hard (also called as jet part) state resulting from the medium-modified multiparton fragmentation. The details of the model and the corresponding simulation procedure can be found in the paper [3] and the references therein. The hard state of an event in HYDJET++ is treated using Pythia Quenched (PYQUEN) model [4] which modifies a jet event produced by PYTHIA by producing nucleonic collision vertices using Glauber model at a certain impact parameter [5–8]. The hard state in HYDJET++ is separated from the soft state by a free parameter named  $p_T^{min}$ . The soft state of a HYDJET++ event is a thermal hadronic state created on the chemical and thermal freeze-out hypersurfaces derived from the parameterization of relativistic hydrodynamics with preset freeze-out conditions [9, 10].



Figure 1: Transverse momentum spectra of all charged paricles with and without jet part for minimum bias collisions over centrality along with ALICE experimental data for comparison[11]. Here we have used  $|\eta| < 0.8$  pseudorapidity cut.

Xenon is a moderately deformed nucleus. The deformation is incorporated via the modification of the Woods-Saxon Nuclear density profile function as explained in references [12, 14]. The values of different parameters have been taken from the reference [15]. The deformed structure permits us to have many collision geometries such as minimum bias ( $\theta$  =random (0 to  $2\pi$ )), body-body ( $\theta = \pi/2$ ), and tip-tip ( $\theta = 0$ ) collisions with respect to the collision axis.

### 3. Results and Discussion

We find a suitable match of model  $p_T$ -spectra (see figure 1) including jet part with ALICE experimental results [11] while without jet part the model completely underpredicts the data. Below 3.0 GeV/c, the collective flow hydrodynamics dominates the spectra but beyond that, hard scatterings come into play.  $p_T$  spectra decreases as we move from most-central to most-peripheral results. Without jet part, the transverse momentum distribution is very small and completely underestimates the experimental data.



**Figure 2:** Variation of average transverse momentum  $\langle p_T \rangle$  with respect to collision centrality at  $|\eta| < 0.8$  for minimum bias, body-body and tip-tip collisions[12]. The results show average  $p_T$  for hydro, jet and hydro+jet part separately comparing them with ALICE experimental data [11] and AMPT model results in string melting version[13].



**Figure 3:** Nuclear Modification Factor  $R_{AA}$  of charged hadrons with respect to  $p_T$  for minimum bias, bodybody and tip-tip collisions over centrality along with ALICE experimental data for comparison [11]. Here we have used  $|\eta| < 0.8$  pseudorapidity cut.

Minimum bias HYDJET++ results for total and hydro part (see figure 2) show suitable match with ALICE experimental data whereas jet part shows opposite behaviour. At a given collision centrality,  $\langle p_T \rangle$  is higher for tip-tip collisions than for body-body collisions. In the AMPT model [13] (string-melting version),  $\langle p_T \rangle$  shows a weak centrality dependence compared to HYDJET++ model. The minimum bias nuclear modification factor  $R_{AA}$  from HYDJET++ model (in figure 3) shows a suitable match with ALICE experimental results [11]. Tip-tip  $R_{AA}$  is higher than body-body  $R_{AA}$ . The suppression increases at low  $p_T$ , reaches its maximum at around  $p_T \simeq 2$  GeV/c. Nuclear modification factor  $R_{AA}$  or suppression increases with collision centrality and then decreases in most-peripheral collisions.

#### 4. Conclusions

Minimum biased transverse momentum distribution of charged hadrons show suitable match with ALICE experimental yield from low to high  $p_T$  region.  $\langle p_T \rangle$  shows strong dependence on centrality with and without jet part. HYDJET++ model shows a suitable match with the experimental data within error bars. However,  $\langle p_T \rangle$  for only jet part is somewhat centrality independent. This behaviour is in good agreement with results from the AMPT model in string-melting version. Minimum bias  $R_{AA}$  of charged hadrons shows a suitable match with the ALICE experimental data, tip-tip  $R_{AA}$  being higher than body-body  $R_{AA}$ .

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### References

- [1] J. Mulligan (ALICE), PoS ICHEP2020, 556 (2021).
- [2] G.-Y. Qin, Nuclear Physics A 931, 165 (2014).
- [3] I. Lokhtin, et. al., Computer Physics Communications 180, 779 (2009).
- [4] I. P. Lokhtin and A. M. Snigirev, The European Physical Journal C-Particles and Fields 45, 211 (2006).
- [5] J. D. Bjorken, Phys. Rev. D 27, 140 (1983).
- [6] E. Braaten and M. H. Thoma, Phys. Rev. D 44, R2625 (1991).
- [7] R. Baier, et. al., Phys. Rev. C 60, 064902 (1999).
- [8] R. Baier, et. al., Phys. Rev. C 64, 057902 (2001).
- [9] N. S. Amelin, et. al., Phys. Rev. C 74, 064901 (2006).
- [10] N. S. Amelin, et. al., Phys. Rev. C 77, 014903 (2008).
- [11] S. Acharya et al. (ALICE), Phys. Lett. B 788, 166 (2019).
- [12] S. Pandey, et. al., Phys. Rev. C 103, 014903 (2021).
- [13] S. Kundu, et. al., EPJ A 55, 157 (2019).
- [14] S. Pandey and B. K. Singh, Journal of Physics G: Nuclear & Particles 49, 095001 (2022).
- [15] P. Möller, et. al., Atomic Data and Nuclear Data Tables 109-110, 1 (2016).