



Domain Walls of Low Tension in Cosmology

Holger Bech Nielsen^{*a*,*} and Colin D. Froggatt^{*b*}

^aNiels Bohr Institute,

Blegdamsvej 15 - 21, Copenhagen, Denmark

^b School of Physics and Astronomy, Glasgow University, Glasgow G12 8QQ, Scotland

E-mail: hbech@nbi.dk, Colin.Froggatt@Glasgow.ac.uk

In the present article we put up for discussion the idea of there existing several versions/phases of the vacuum, in the spirit in which we have long worked on this idea, namely the Multiple Point Citicality Principle, which also says that these different vacuum phases have the same energy density. We mention that we indeed predicted the Higgs mass to be (135 ± 10) GeV, which when measured turned out to be 125 GeV, using the assumption of this Multiple Point Criticality Principle. We consider the possibility that there is one type of vacuum in the galaxy clusters (the usual vacuum) and another type of vacuum in the voids. The hope that there could indeed be such a low tension S of the domain wall between these two phases, that it would not totally upset cosmology is based on our dark matter model. In this model dark matter consists of pearls or bubbles of a new vacuum phase, with ordinary matter inside it under very high pressure. The order of magnitude of $S^{1/3} \sim MeV$ or 100 MeV could make such domain walls astronomically viable. We successfully estimate the order of magnitude of the variations in the fine structure constant in different places astronomically, but the similar variations in proton mass over electron mass should have been much bigger than seen experimentally in our model. The Universe's surprisingly early galaxies seen by JWST (James Webb telescope) may agree well with our model. Replacing the usual cosmological constant by domain walls in the standard cosmological model would lead to a cubic root of the tension being $S^{1/3} \sim 30 MeV$.

CORFU2022: 22th Hellenic School and Workshops on Elementary Particle Physics and Gravity Workshop on Tensions in Cosmology Corfu Sep 07 - Sep 12, 2022

*Speaker

[©] Copyright owned by the author(s) under the terms of the Creative Commons Attribution-NonCommercial-NoDerivatives 4.0 International License (CC BY-NC-ND 4.0).

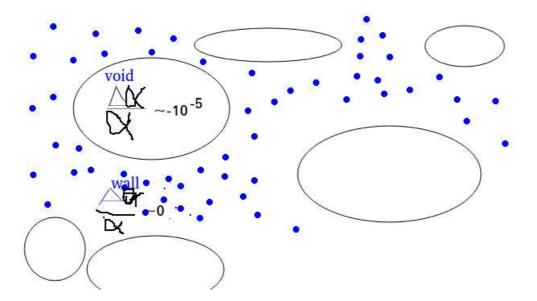


Figure 1: "Artistic impression": Vacua in Universe; Voids one vacuum, clusters of galaxies "our" vacuum. The small dots symbolise galaxies of course. The formulas should illustrate that we imagine that the fine structure constant α should deviate a little bit $\Delta \alpha$ from the usual value, or say the value on Earth, when one measures it in a void. The formula suggests that the deviation is relatively of the order of 10^{-5} in the voids, while assuming the Earth to be in a galaxy dense region there should be no deviation between the Earth-value and the galaxy rich region value.

1. Introduction

Investigating the three dimensional distribution of the galaxies shows that the galaxies cluster into surfaces and line-like structures, so as to leave a lot of voids with relatively much fewer galaxies. The main speculation of the present contribution is that the idea of there being several versions of the vacuum, which we used in the talk [1] in the foregoing workshop "Standard Model and Beyond" as the speculative input in our model for dark matter[2-12], could also mean that there were different vacua in different regions in outer space. For example, one could think that there were one type of vacuum in the voids and another type in the galaxy rich regions. In our model of dark matter, a different vacuum from the usual one keeps some ordinary matter under high pressure in a very concentrated bubble, so that it is essentially hidden and does not interact much with ordinary matter despite the high mass of the bubble. Thus we have a picture of the dark matter as genuinely containing some ordinary matter in it; in this way it has a slight similarity to Khlopov's model [13-16] in which a helium nucleus is part of the dark matter complex. In Figure 1 the voids in the Universe are marked by ellipses in a picture of the galaxy pattern drawn in perspective. If you take these ellipses as symbols for the boundaries between the different vacuum phases, then you could consider this figure to suggestively illustrate how there is speculated to be one phase of the vacuum in the voids and another in the galaxy rich regions.

2. Domain Walls need Low Tension to be Allowed

The tension or in c = 1 notation also the energy per area S of a domain wall between different phases of the vacuum has dimensionality GeV^3 . If an energy scale of the order of the scales at which we look for "new physics" is used, the energy density gets so high that domain walls of cosmological scales of extension would be so heavy as to be totally excluded by the already known energy density (the critical density):

$$\rho_c = \frac{3H^2}{8\pi G} = 1.88 * 10^{-26} h^2 kg/m^3 \tag{1}$$

(where
$$h = H_0 / (100 km / Mpc) = 0.67$$
) (2)

$$= 2.78 * 10^{11} h^2 M_{\odot} M p c^{-2}, \tag{3}$$

$$= 8.5 * 10^{-27} kg/m^3.$$
 (4)

The average energy density that would be produced by even just a single domain wall crossing the visible universe with energy density *S* of the order of the "new physics", meaning say TeV scale to the third power $S \sim TeV^3$, would dominate the critical energy density in the Universe. So we can conclude that "*New physics scale*" walls are excluded.

3. How high an Energy per Area is possible?

Consider a domain wall having a dimension of the visible Universe length scale:

$$R_{visible} = 13 * 10^9 \text{ light years}$$
(5)

$$= 1.3 * 10^{10} ly * 9.5 * 10^{15} m/ly = 1.2 * 10^{26} m,$$
(6)

which therefore has

$$Area \sim 10^{52} m^2. \tag{7}$$

Now the total visible critical energy is

visible (critical) energy ~ $(1.2 * 10^{26}m)^3 * 8.5 * 10^{-27} kg/m^3$ (8)

$$= 1.5 * 10^{52} kg.$$
(9)

So this would allow the domain wall to have an energy up to

maximal energy per area ~
$$\frac{1.5 * 10^{52} kg}{10^{52} m^2} \sim 1.5 kg/m^2$$
. (10)

Hence the tolerable energy per area for a domain wall of astronomical extension is ~ $1.5kg/m^2$:

Tolerable energy per area =
$$S = 1.5kg/m^2$$
 (11)

$$= \frac{1.5kg/m^2 * (1.2 * 10^{-12} MeVm)^2}{1.79 * 10^{-30} kgc^2/MeV}$$

= 1 * 10⁶ MeV³. (12)

So the energy scale pointed out by such a slightly too big domain wall energy per area is

$$\sqrt[3]{S} = \sqrt[3]{10^6 MeV^3} = 100 MeV.$$
 (13)

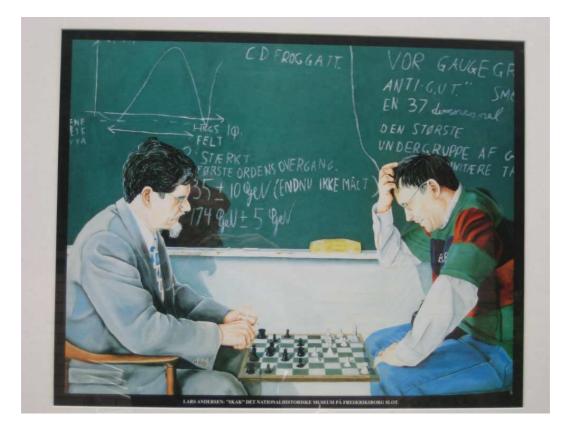


Figure 2: This painting was put up at Frederiksborg Castle near Hillerød *before* the Higgs particle was found in 2012, so nobody knew the mass of the Higgs then. Nevertheless we had already published the predicted Higgs mass 135 ± 10 GeV and except for the 1 hidden by Mogens Lykketofts head you see this number on the painting.

4. Older Achievements

We have worked for a long time on the idea of there being several vacua with the same energy density and this has already led to some achievements:

- *PRE*dicted the Higgs mass *before* it was found [17, 18]. See Figure 2.
- In a rather complicated model with each family of fermions having its own set of gauge fields - like the ones in the Standard Model - one of us (HBN), Niels Brene, Don Bennett, I Picek and others [19–21] fitted the fine structure constants by using the number of families as a parameter and *got 3 before this number was measured*.
- We model dark matter particles or pearls [2-12] as bubbles of a new vacuum containing highly compressed ordinary matter. We get the surface tension *S* of the domain wall between the vacua to have its cubic root of the order

$$S^{1/3} \sim 10 \, MeV.$$
 (14)

An important part of our fitting in this model involves the 3.5 keV X-ray line observed from satellites [22–24].

5. Big Bubbles

The main point of the present talk is to point out that, once you have the suggestion that the tensions S of the domain walls between different vacuum-versions could be as small as the numbers mentioned $(10 MeV)^3$ or $(100 MeV)^3$, then astronomically extended domain walls become a possibility. One concrete idea could be that the domain walls lie around the voids in outer space containing a very low density of galaxies. That is to say we propose the *Idea of Domain Walls around Voids*:

Since our dark matter model has one vacuum inside the dark matter, a large region of such inside the dark matter vacuum would be formally a huge pearl of dark matter, and there could not be any true dark matter inside it. Such a difference would almost certainly have an influence on the number of galaxies formed at different times. But thinking about the likely situation inside such a huge bubble with the inside the dark matter vacuum, we should imagine that the usual dark matter there has been replaced by an equal amount (by mass) of ordinary matter. In such a situation there would be a better chance for the material when contracting to emit radiation and cool down, than in a usual region with both dark matter and (much less) ordinary matter. Thus we would at first expect a faster formation of galaxies and stars in such a region of the "inside dark matter vacuum" than in a region like our own. The matter will namely fall faster together. In turn this could mean that the faster start means a faster burn out and that the galaxies in the "inside dark matter vacuum" regions would 13 milliard years later, in our epoch, send out less light than the usual regions. So today it could be that the galaxies would look sparser, because of being older and giving less light.

In the regions near the domain wall - the outermost part of the "inside dark matter vacuum" region - where the nuclei would probably come in from outside and get speeded up to high temperature, there would be a high temperature and at least in the beginning not much star formation.

So let us investigate as an example the hypothesis that the domain walls follow the borders of the voids - the regions with few or no galaxies - so that these voids really are "formally huge dark matter pearls". This idea is illustrated in Figure 3

6. MPP

We originally proposed our "Multiple Point Criticality Principle" (= MPP) [25–28] to justify that we fitted e.g. the fine structure constants in a somewhat complicated model with the phase transition values taken from a lattice model. That is to say we would generally argue that there should be in Nature some law saying that the values of coupling constants should be made to sit on the phase transition point(s).

This is what (often) happens with a microcanonical system - meaning a system in which the energy is fixed - when it splits up and you get more than one phase at the same time so as to achieve the given total energy. For example, you may get as in Figure 4 both ice and fluid water together provided there is a (fixed) given amount of energy, as in the microcanonical ensemble. It should be stressed that even though the fixed energy is *not fine-tuned* you end up with more than one phase existing together, and that can only happen in the case of slush say at one specific (fine tuned) temperature 0^0C . So one got the fine tuning without putting any fine-tuning in! Most often the two

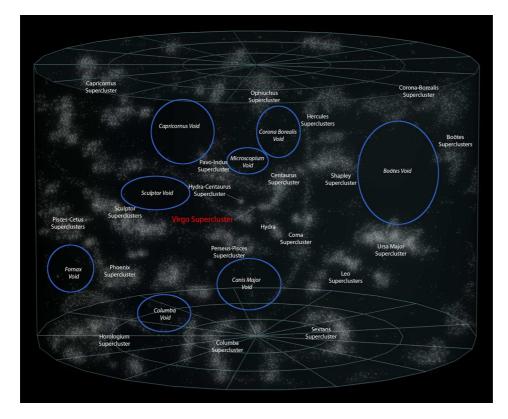


Figure 3: Here is shown in perspective the measured galaxy densities and the voids with very few galaxies are marked with rings and names. Accidentally these rings just made (in Wikipedia) by the author of the figure to point out the voids could be considered our proposal for where to find the domain walls.

or several phases achieved in this way, by fixing *extensive quantities* (like total energy or the total amount of some element or some molecule) will occur (at the end) in comparable amounts.

The suggestion is that it is a similar mechanism to the microcanonical ensemble, that causes the coupling constants in Nature to tend to be precisely at phase transition points - the "multiple point". But then, by analogy to the slush etc., there should be extended amounts of the different phases of *comparable size*. It is these extended phases, which we think of as being different versions of the vacuum. The proposed mechanism for why we should have these different phases of the vacuum predicts that they should occur in order of magnitudewise comparable amounts. So it would not be expected that the one phase would only occur inside some tiny dark matter bubbles, while another phase would be extended over the whole world and fill almost everything. No, with this idea of extensive quantities fixed, the analogy suggests that also the inside the dark matter bubbles phase must occur as astronomically extended huge bubbles, e.g. in the voids.

7. Vary α

We consider here, as a signal for the above suggested existence of several vacuum phases in comparable amounts, the possibility that the natural constants such as the fine structure constant α should be different in the different phases. As it is possible to observe the fine structure constant in



Figure 4: Slush $\Rightarrow 0^0 C$.

different clouds of atoms of hydrogen, such an investigation has indeed been performed although the accuracy has so far only been barely sufficient to give convincing results.

7.1 Expected Deviations from Constant α

But let us first estimate the order of magnitude of the deviation from a constant fine structure α from vacuum phase to vacuum phase, which we would predict using the parameters extracted from our long studied model for the dark matter:

• Cubic root of tension in wall S or of energy density per unit area

 $\sqrt[3]{S} = 10 MeV$ (allowed for cosmological domain walls).

We extracted this tension from the assumption that we could estimate how large the typical size of our bubbles making up the dark matter would be and the pressure inside the bubbles. This pressure was estimated from the frequency of the radiation of X-rays emitted from the dark matter pearls being identified with the mysterious 3.5 keV radiation line observed in X-ray spectra [22–24].

• Lowering of the potential energy for a nucleon by entering into the "new" vacuum (the one inside the dark matter bubbles)

$$\Delta V = 2 M e V. \tag{15}$$

Nucleons 2 per mille lighter in the "new" vacuum.

Actually we found [10] that both the frequency of the 3.5 keV X-ray line and the overall magnitude of this radiation as obtained by Cline and Frey [29] could be fitted with just one parameter combination $\frac{\Delta V}{\xi^{1/4}}$. Here ξ is the ratio of the typical dark matter bubble radius to the critical one below which the nuclei in it get spit out. This was in a model, wherein the X ray emission came from colliding dark matter pearls uniting with each other.

• Energy density in the two vacua exactly the same. (Also a theorem by Gia Dvali).

This was the original main assumption of the "Multiple Point Criticality Principle" (=MPP), which we mentioned gave us the Higgs mass prediction. But it was actually proven by Gia Dvali [30] that it should not be possible to have several vacuum phases unless they have the same energy density.

7.2 Calculation of the Order of Magnitude of the Difference in α from phase to phase.

A physical argument for the change of the (effective) fine structure constant:

Because of self energy diagram corrections, we should think of the photon as being replaced with some amplitude by a proton anti-proton pair virtually. The probability for finding the photon in this virtual state of a proton anti-proton pair is proportional to the fine structure constant, because the coupling causing the transition is proportional to e the charge quantum and the probability for finding the proton anti-proton pair will go as the square of this transition (of course $\alpha \propto e^2$). But it is also inversely proportional to the excess energy of the proton anti-proton state; this is the well-known second order perturbation formula for the energy shift in quantum mechanics.

The effect of such a partial virtual replacement of the bare photon by the proton anti-proton pair is to diminish the coupling of the bare photon, because it is part of the "time" (really probability) replaced by the not directly coupling pair. The bigger the chance of finding the pair the smaller the coupling of the dressed photon propagator. This probability for finding the pair, instead of the bare photon itself, increases with lower energy or mass of the pair. So the coupling α is the smaller in the case of the lighter proton and anti-proton. This lighter pair situation is in our model for dark matter supposed to happen in the inside the dark matter phase. So we predict that the smaller fine structure constant should be found in the inside the dark matter phase. But since we have dark matter around in our galaxy, we must exist in a phase of outside the dark matter vacuum. So we should ourselves be in a phase with slightly bigger α . Thus we predict that the deviating α clouds should have a lower α than the value here.

The order of magnitude is expected to be proportional to the ratio $\frac{\Delta V}{m_p}$ where m_p is the proton mass and ΔV is our estimated relative potential energy for a nucleon in the two phases, which we claim to be of the order of $\Delta V \approx 2MeV$.

Let us now consider the proton loop self energy diagram for the photon in the two phases:

• In the approximation that one could use non-relativistic second order perturbation theory for the sign, we expect that the (bare) photon is mixed more with the proton anti-proton pair in the phase with the smallest nucleon masses. Thus the photon in this (m_p small) phase couples weaker than in the m_p big phase.

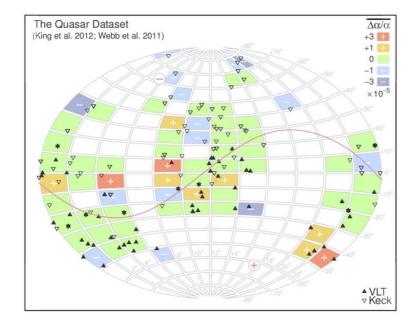


Figure 5: A sky graph [33] giving the small regions in which the various deviations for the fine structure constant have been measured, because appropriate clouds of atomic gas had been radiated from behind by quasars and the generated absorption lines observed. The agreement with our prediction of the sign, that $\frac{\Delta \alpha}{\alpha} \leq 0$ is really not well satisfied, in as far as it should mean on this plot that there were no red nor yellow squares, but only the no deviation green squares and the blue ones which have positive deviations. But the order of magnitude is o.k. Strictly speaking we would predict only two different values, but clearly there are huge statistical fluctuations (indeed the uncertainty would allow for no deviations at all).

• The relative correction to the fine structure constant α , i.e. $\frac{\Delta \alpha}{\alpha}$, has one factor α in addition to the 2/940, the relative difference in nucleon mass between the two phases:

$$\frac{\Delta \alpha}{\alpha} \approx -\alpha * \frac{2}{940} = -1.4 * 10^{-5}$$
 (16)

• The deviations found (see Figure 5) are of this 10^{-5} order of magnitude, but in the literature they are fitted by a dipole variation in space [31, 32] rather than with voids versus clusters.

More detailed observations are needed before a proper check of our model can be made, but in Figure 5 we show a sky-graph from reference [33] with the results found so far. It is necessary to investigate whether the clouds studied were in voids or in galaxy rich regions. We predict that deviations $\frac{\Delta \alpha}{\alpha}$ shall be of order 10^{-5} giving smaller values of α in the voids than in the galaxy rich regions. But whether you have a void or a galaxy rich region not only depends on the direction but also on the distance, which in practice has to be taken as the red shift z.

8. m_p/m_e ratio not varying at 10^{-6} level is problematic for our model

Searches have also been made to detect cosmological variations in the proton to electron mass ratio $\mu = \frac{m_p}{m_e}$ using molecular absorption lines and μ was found not to vary by more than one part in a million.

The wavelengths of the Lyman and Werner transitions of the H_2 molecule are sensitive to μ and King et al [34] used the spectra from 3 quasars, viewed through molecular clouds, to limit the variation in μ relative to the present day laboratory measured value to $\frac{\Delta\mu}{\mu} = (2.6 \pm 3.0) \times 10^{-6}$. Also the inversion transitions of ammonia are particularly sensitive to μ when compared to molecular rotational transitions. Murphy et al [35] used this technique to limit relative deviations from the laboratory value to $\frac{\Delta\mu}{\mu} < 1.8 \times 10^{-6}$. Later Kanekar [36] improved this limit to obtain $\frac{\Delta\mu}{\mu} < 3.6 \times 10^{-7}$. Methanol transitions were then used [37] to find the more stringent limit $\frac{\Delta\mu}{\mu} < 1.8 \times 10^{-7}$.

In our model the proton mass varies on the per mille level ($\Delta V \approx 2MeV$) from phase to phase. So, of course, we would predict that μ should vary on the per mille level between one vacuum phase and another one, in strong conflict with the above limits. Our only explanation could be: *There are very few molecular clouds in the voids, so they were not measured.*

9. Our picture of domain walls pushed around by nucleon pressure

In our dark matter model there is an interaction between nucleons and the domain wall, in as far as the nucleons have a lower potential of the order of a few MeV in the "inside-dark-matter" vacuum phase. This means that big bubbles of this "inside" type will tend to be blown further up in size by the nucleons crossing the wall from the "outside"-side and coming into the "inside"-side with dramatically higher velocity. The consequent recoil of the domain wall will cause the "inside" vacuum to become even bigger. In fact a balance may be achieved in which a temperature difference between the two sides counteracts the pressure connected with the curvature of the wall. That should mean the temperature should be higher in the "inside the dark matter" vacuum, and thus the formation of stars and galaxies in the present epoch would be made more difficult. In fact very hot gas, which rather becomes plasma, is basically invisible, since the interaction of light with just nuclei and electrons corresponds to Compton scattering with a very broad spectrum and does not give a very characteristic signal.

Domain wall universes have been studied [38–40], but we shall not go much into that in the present article.

10. After James Webb Observation Addendum

Recently the James Webb telescope discovered galaxies that had been produced much earlier than expected in the development of the Universe [41–44], even some within the first half a milliard years after the Big Bang. We should like to present a possible explanation for such surprisingly early stars and galaxies in terms of the model discussed above with its domain walls and picture of dark matter.

Let us therefore give a very short survey of how our model would influence or not influence the early cosmology:

• From the very beginning we should expect our domain walls to be present due to thermal fluctuations to the degree that they make sense at all. In detail we have in mind as our best guess that our domain walls are the result of QCD with the Nambu JonaLasinio spontaneous break down of chiral symmetry being in more than one phase. It could be said that QCD in

another phase was not quite in the usual confinement phase but that the color symmetry was rather say Higgsed - spontaneously broken - so as to be aligned along the breaking direction of the Gellman SU(3), if such an alignment could make any sense. At least one has two SU(3) groups at hand. But whatever the phase distinction might be, we speculate from the hoped for order of magnitude of the domain wall tension that it is involved with physics at the energy scale of MeV's or up to hundreds of MeV. Then it would mean that for higher temperatures than that, we would have the domain walls all over the Universe as a result of thermal fluctuations. This could give an extra, otherwise unexpected, contribution to the effective cosmological constant in say the first second after the Big Bang.

- As temperature falls the domain walls, starting from being distributed all over the Universe, begin to contract and two types of region occur corresponding to the two vacua. Now the type of vacuum ending up as the inside the dark matter pearls one is supposed to attract the nucleons, which thus collect in that phase. For reasons which we would have to calculate or just speculate about the phase going to be the one outside the dark matter pearls wins and becomes at first the dominant phase. Because of the higher pressure inside the dark matter pearls or proto-pearls, the Big Bang Nucleosynthesis is supposed to take place earlier inside these proto-pearls. We have previously speculated [2, 3] that an explosive fusion of helium into say carbon occurs inside the pearls, where almost all the nucleons are collected. Hereby a fraction of the nuclei is emitted and becomes the ordinary matter, while the remaining pearls when contracted make up the dark matter.
- A very few bubbles are supposed accidentally to be exceptionally large, and thus the pressure needed to keep them from collapsing is appreciably smaller. Now, because of the attraction of the nucleons, the nucleons are effectively pulled into such a large bubble making it even larger. The nucleons hitting the domain wall from the usual vacuum side will namely be pulled in with a force from the potential difference for nucleons on the two sides. This will give a reaction force pulling the domain wall from the inside dark matter vacuum towards the usual vacuum. Once the bubble has got sufficiently large the oppositely pointing force due to the tension in the domain wall will be so weak that the bubble of inside the dark matter vacuum they will unite with them, and the material in the hit pearls will be spit into the big bubble. So this will further contribute to making the very big bubble even larger. In this way some very few bubbles passing a critical limit can be very large. It would be such exceptionally large bubbles, which have passed the critical limit beyond which they grow instead of contracting, that could be the ones filling say the big voids.
- If the big bubbles are produced in the just described way, by growing up from pearls only a bit larger than the usual dark matter pearls, they will say in the era of Big Bang Nucleosynthesis for the ordinary matter not yet be so dramatically large. Thus, since there are very few of them, they would probably have very little significance in this Big Bang Nucleosynthesis Era, and we should not see much effect of them in the fitting of the abundances of the light elements like helium, lithium and the hydrogen isotopes. It is of course important for the viability of our model that it does not make much change in the Big Bang Nucleosynthesis,

which works very well without any extra ordinary matter. It is rather crucial that our dark matter model, which genuinely consists of ordinary matter, has it so strongly packed into the pearls that the packed in part cannot make itself felt in the Big Bang Nucleosynthesis.

- But as time develops and the exceptionally large bubbles grow and grow they can come to cover significant parts of the space in the Universe. Now you then have two different types of region in the Universe:
 - There is the region with the usual vacuum, which is the outside the dark pearls one, covered with a swarm of dark matter pearls which are really spheres of the other type of vacuum filled with highly compressed ordinary matter inside such as say carbon. Interpreting the highly pumped up pearls as dark matter, this part is exactly the usual picture of the cosmological content.
 - There are also supposedly comparable amounts of space with the vacuum of the inside the dark matter pearl type. In these regions there can be no dark matter. In a formal sense such a whole huge region is a piece of dark matter. However this is only formally so, because the density of ordinary matter inside is only about 6 times the density of ordinary matter in the other type of region. But now there cannot be any true dark matter in this region. One could think of having small pearls of the outside the dark pearls type, but that does not attract the nucleons, it rather repels them. So one cannot make heavy pearls from that inside the dark matter vacuum. It would just collapse, unless one could find something else with which to pump them up.

But in these regions you would have rather closely the same amount of the ordinary matter as one has ordinary plus dark matter in the other regions. This would mean that the formation of the clumps or fluctuations would be very much the same for both types of region, as long as it is only the question of their non-relativistic nature that counts.

• However if at some point the interaction with e.g. photons becomes important, then the two regions above could behave very differently. The dark matter pearls should be so heavy compared to the number of electrons they would likely take up or lose that, even if they were hot, they would not be able to emit so much light as the atoms that can form from the nuclei like hydrogen and helium. When a clump of either dark matter or ordinary matter contracts, it heats up, but then comes the difference: The dark matter interacts only very little with light and, in first approximation can not cool down but rather stops the contraction by the increased pressure due to the heating. However, for the ordinary matter which contains atoms that can emit light and get cooled, the contraction can go further heating the ordinary matter but it can continue to cool. So the ordinary matter contracts by the help of the cooling by light emission, while the dark matter with a very slow variation of the density in space, and a much smaller clump of ordinary matter in the middle of the dark clump.

But in a region in which the vacuum is of the inside the dark pearls type we have only ordinary matter - and about 6 times the density of the other region - but no genuine dark matter (although formally the whole region is a dark matter bubble). Now it can all contract

as fast as ordinary matter by emission of light. So a clump of a certain geometrical size would in such a region with high ordinary matter density contract without leaving a much larger dark matter halo around. A size that should have given a small galaxy could when it is 6 times as dense give an appreciably more massive galaxy.

This was the surprise from the James Webb early galaxy observations: There were at an early time more massive galaxies - with mass as large as the Milky Way say - than one should expect at the time 500 million to 700 million years after the Big Bang [44]. It should, according to the theoretical expectation, have been the dark age with essentially no stars and thus only rather little light.

If this explanation of ours for the too early star and galaxy formation should be the right one, then from the above description you may notice that we had two regions, and the early stars and galaxies only formed in one of the regions. In the region with the vacuum which is outside the dark pearls the prediction of the astronomers that there should be a dark age should hold just as predicted. So we should expect that the mysterious too early galaxies should clump into some regions, and not be present outside these regions.

Roughly you could say one type of region starts earlier than the other type.

10.1 Compact Galaxies

For the six galaxies reported in reference [44] to appear old and massive means they were forming hundreds of stars a year shortly after the Big Bang. In comparison, the Milky Way only forms around one to two new stars every year. Further, these potential galaxies are said to be about 30 times more compact in size than ours despite having as many stars. "Compact" must mean geometrically of less extent than the Milky Way.

If it were true that they were formed in a vacuum of the type which is usually inside the dark matter, there would be about 6 times as much ordinary matter in the region of the recently found very early galaxies than around normal galaxies. Furthermore, contrary to galaxies in our region of the Universe, their contraction would not be provided by dark matter but rather by ordinary matter. Since ordinary matter tends to contract much more than the dark matter, this would mean that these early galaxies should be expected to be geometrically smaller than usual galaxies. So it would be an immediate qualitative success of our two vacuum phase story that the surprising early galaxies are indeed more compact.

11. Chance of getting rid of the Dark Energy!

We should call attention to the fact that the idea of having extended domain walls of the unexpectedly low order of magnitude of the tension, which we propose, gives a chance of setting the cosmological constant to zero in fits to the cosmological data, in spite of the famous increasing speed of Hubble expansion. This idea is contained in the work of Richard A. Battaye et al [38]. The point is that replacing the energy density from the cosmological constant (the dark energy) by the average energy density from a system of domain walls, spread relatively smoothly in space on a large scale, will give 2/3 of the negative pressure of the corresponding energy of a cosmological constant. This means that if one can squeeze into the fit just 3/2 times as much domain wall energy

as is in the present fit assigned to the cosmological constant, then one can get the domain walls do the job of explaining the increasing Hubble expansion in much the same way as the cosmological constant. We would say that replacing the cosmological constant by domain walls is an attractive idea, if possible, although we must admit that it is a great mystery why the cosmological constant is so small as it is and that it would still be a mystery or perhaps even a bigger one, if it were even smaller.

In section 3 we obtained an upper limit of $S^{1/3} < 100$ MeV for the tension of a single domain wall having a dimension of the visible Universe length scale. Here we will estimate the tension required if the domain walls lie along the borders of the voids and their energy density could replace the dark energy density.

We take the voids to have sizes of the order of 100 million light years, so that the typical distance between domain walls is

"distance between walls" = $(30 \ to \ 300)$ million light years (17)

$$\approx$$
 100 million light years. (18)

In fact we can crudely ignore whether the dark energy is 75% or 100% of the energy density of the Universe, as we are only making order of magnitude estimates. So let us take the energy density of the dark energy, which we want to replace by domain walls in our cosmological model, to be

$$\text{'energy density''} = 10^{-26} kg/m^3 \tag{19}$$

and "distance between domain walls" =
$$10^8$$
 light years = $10^{24}m$. (20)

Then

"wall surface energy density needed" =
$$10^{24}m * 10^{-26}kg/m^3 = 10^{-2}kg/m^2$$
 (21)

$$(30 \, MeV)^3$$
. (22)

Thus the domain wall tension or energy per unit area needed for replacing dark energy has the cubic root:

$$S^{1/3} = 30 \, MeV. \tag{23}$$

In our model of dark matter bubbles of radius R, the pressure provided by the surface tension S of the surrounding domain wall has to balance the pressure due to the highly relativistic degenerate electron gas inside the bubble having a Fermi momentum p_f :

$$\frac{2S}{R} = \frac{1}{12\pi^2} p_f^4.$$
 (24)

Assuming there are roughly equal amounts of protons and neutrons, there must be about twice as many nucleons as electrons inside the bubble. Then it follows [4] that

$$S^{1/3} = \left(\frac{M}{24\pi^5 m_N}\right)^{1/9} \frac{p_f}{2},\tag{25}$$

where *M* is the mass of the dark matter particle and m_N is the nucleon mass. The Fermi momentum is given in terms of the change in potential energy ΔV of a nucleon when it passes through the domain wall by

$$p_f = \frac{2\Delta V}{\xi^{1/4}},\tag{26}$$

where $\xi = R/R_{crit}$ and R_{crit} is the "critical size" of the bubble below which it would collapse.

From our fits [9] to the frequency and intensity of the 3.5 keV X-ray line observed from galaxy clusters, Andromeda and the Milky Way Center, we found that the parameter $\frac{\Delta V}{\xi^{1/4}} \approx 2 MeV$. So equation 25 can be written in the form:

$$\frac{M}{1.24 * 10^{-23} kg} = \left(\frac{S^{1/3}}{2MeV}\right)^9.$$
 (27)

Then using the tension $S^{1/3} = 30 MeV$ from equation 23 for domain walls that can replace the dark energy, we obtain the mass of the corresponding dark matter pearl to be

$$M = 15^9 * 1.24 * 10^{-23} kg \tag{28}$$

$$= 4.8 * 10^{-13} kg. \tag{29}$$

The radius of such a dark matter pearl is given by

$$R = \frac{1}{p_f} \left(\frac{9\pi M}{8m_N}\right)^{1/3}$$
(30)

$$= (4 M eV)^{-1} * (10^{15})^{1/3}$$
(31)

$$= 5 * 10^{-9} m. \tag{32}$$

This size is an order of magnitude smaller than the size $10^{-7}m$ of a typical intergalactic dust grain, so that our considerations that the dust grain typically will be larger than the bubble with the radius *R* remains true.

12. Conclusion

We have worked for a long time on a model for dark matter entirely within the Standard Model, speculating that there is more than one vacuum phase and that the different phases have the same energy density. Our dark matter particles have a size of order a few nanometre and a surface tension of roughly $S = (10 MeV)^3$. In this talk we have mainly considered possible cosmological implications of our model and we list them below.

- Most likely the different vacuum-phases should have similar order of magnitudes of spacetime extension. (Microcanonical ensemble).
- Different vacuum-phases have different "constants" of nature, e.g. the fine structure constant α . Such variation has possibly been observed, but not yet tested for void versus cluster dependence.

- A similar search for variations from phase to phase of the proton to nucleon mass ratio turns out rather catastrophically for our model, in as far as the observations show that the deviation if any is less than of the order of one in ten million, while we would expect deviations at the per mille level for this ratio.
- The very recent observations by the James Webb telescope of very early galaxies may be considered a success for our model, since they could appear in the regions with our "inside the dark matter phase vacuum" where the dark matter would have been replaced by ordinary matter. Stars and galaxies would most likely form earlier in these regions than in regions with the usual mixture of dark and ordinary matter.
- The lack of galaxies today in the regions of the inside the dark matter phase, identified with the voids, might find an explanation by the heating up by nuclei crossing the domain wall from the outside the dark matter phase into this phase as it expands into "our phase" (= the one we live in).
- We have proposed a physical idea in QCD with spontaneous breaking of chiral symmetry: Once you have a mass for the quarks, you can imagine that
 - Either the breaking is genuinely broken by the quark mass, so that there is no genuine spontaneous breaking, or
 - the quark mass is only guiding the direction of breaking in a genuine spontaneously broken chiral symmetry.

We propose a phase transition between such phases. With several quark masses a breaking in several steps may be possible.

References

- H. B. Nielsen and C. D. Froggatt, "Dusty Dark Matter Bubbles of a New Vacuum stopped in Earth and Radiating 3.5 keV X-rays" in "Standard Model and Beyond" in Corfu Summer Institute 2022 [arXiv:2305.12291].
- [2] C. D. Froggatt and H. B. Nielsen, Phys. Rev. Lett. 95 231301 (2005) [arXiv:astro-ph/0508513].
- [3] C.D. Froggatt and H.B. Nielsen, Proceedings of Conference: C05-07-19.3 (Bled 2005); arXiv:astro-ph/0512454.
- [4] C. D. Froggatt and H. B. Nielsen, Int. J. Mod. Phys. A 30 no.13, 1550066 (2015) [arXiv:1403.7177].
- [5] C. D. Froggatt and H. B. Nielsen, Mod. Phys. Lett. A30 no.36, 1550195 (2015) [arXiv:1503.01089].
- [6] H.B. Nielsen, C.D. Froggatt and D. Jurman, PoS(CORFU2017)075.
- [7] H.B. Nielsen and C.D. Froggatt, PoS(CORFU2019)049.

- [8] C. D. Froggatt, H. B. Nielsen, "The 3.5 keV line from non-perturbative Standard Model dark matter balls", arXiv:2003.05018.
- [9] H. B. Nielsen (speaker) and C.D. Froggatt, "Dark Matter Macroscopic Pearls, 3.5 keV -ray Line, How Big?", 23rd Bled Workshop on What comes beyond the Standard Models (2020), arXiv:2012.00445.
- [10] C. D. Froggatt and H.B.Nielsen, "Atomic Size Dark Matter Pearls, Electron Signal", 24th Bled Workshop on What comes beyond the Standard Models (2021), arXiv:2111.10879.
- [11] C. D. Froggatt and H.B. Nielsen, "Atomic Size Pearls being Dark Matter giving Electron Signal", arXiv:2203.02779.
- [12] H. B. Nielsen and C. D. Froggatt, PoS(CORFU2021) 095 [arXiv:2205.08871]
- [13] M. Khlopov, Int. J. Mod. Phys. A28 1330042 (2013) [arXiv:1311.2468].
- [14] D. Fargion, M. Khlopov and C. Stephan, Class. Quantum Grav. 23 7305 (2006) [arXiv:astroph/0511789].
- [15] M. Khlopov and C. Kouvaris, Phys. Rev. D77 065002 (2008) [arXiv:0710.2189].
- [16] M. Khlopov, Front. in Phys. 7 4 (2019).
- [17] C. D. Froggatt and H. B. Nielsen, Phys. Lett. B368 96 (1996) [arXiv:hep-ph/9511371].
- [18] H.B. Nielsen (Speaker) and C.D. Froggatt, Presented at Conference: C95-09-03.1 (Corfu 1995); arXiv:hep-ph/9607375.
- [19] D. Bennett, H.B. Nielsen, and I Picek. Phys. Lett. B208 275 (1988).
- [20] H. B. Nielsen, "Random Dynamics and relations between the number of fermion generations and the fine structure constants", Acta Physica Polonica Series B(1), 1989. Crakow School of Theoretical Physics, Zakopane, Poland.
- [21] H.B. Nielsen and N. Brene, Gauge Glass, Proc. of the XVIII International Symposium on the Theory of Elementary Particles, Ahrenshoop, 1985 (Institut fur Hochenergi- physik, Akad. der Wissenschaften der DDR, Berlin-Zeuthen, 1985); D.L. Bennett, N. Brene, L. Mizrachi and H.B. Nielsen, Phys. Lett. B178 (1986) 179.
- [22] E. Bulbul, M. Markevitch, A. Foster et al., ApJ. 789, 13 (2014) [arXiv:1402.2301].
- [23] A. Boyarsky, O. Ruchayskiy, D. Iakubovskyi and J. Franse, Phys. Rev. Lett. 113, 251301 (2014) [arXiv:1402.4119].
- [24] A. Boyarsky, J. Franse, D. Iakubovskyi and O. Ruchayskiy, Phys. Rev. Lett. 115, 161301 (2015) [arXiv:1408.2503].
- [25] D. L. Bennett, C. D. Froggatt and H. B. Nielsen, NBI-HE-94-44, GUTPA-94-09-3, Presented at Conference: C94-07-20 (ICHEP 1994), p.0557-560.

- [26] D. L. Bennett, C. D. Froggatt and H. B. Nielsen, NBI-95-15, GUTPA-95-04-1, Presented at Conference: C94-09-13 (Adriatic Meeting 1994), p.0255-279 [arXiv:hep-ph/9504294].
- [27] D. L. Bennett and H. B. Nielsen, Int. J. Mod. Phys. A9 5155 (1994).
- [28] D. L. Bennett, C. D. Froggatt and H. B. Nielsen, NBI-HE-95-07, Presented at Conference: C94-08-30 (Wendisch-Rietz) p.394-412.
- [29] J. M. Cline and A. R. Frey, Phys. Rev. D90, 123537 (2014) [arXiv:1410.7766].
- [30] Gia Dvali, arXiv:1107.0956.
- [31] J. K. Webb et al, Phys. Rev. Lett. 107 191101 (2011) [arXiv:1008.3907].
- [32] J. A. King et al, Mon. Not. R. Astron. Soc. 422 3370 (2012) [arXiv:1202.4758].
- [33] E. Cameron and T. Pettitt, arXiv:1207.6223.
- [34] J. A. King et al., Phys. Rev. Lett. 101 251304 (2008) [arXiv:0807.4366]
- [35] M. T. Murphy, V. V. Flambaum, S. Muller and C. Henkel, Science 320 Issue 5883, 1611 (2008) []arXiv:0806.3081]
- [36] N. Kanekar, ApJL 728 L12 (2011) [arXiv:1101.4029]
- [37] N. Kanekar et al, MNRAS 448 L104 (2015) [arXiv:1412.7757]
- [38] R. A. Battye, M. Bucher and D. Spergel, arXiv:astro-ph/9908047.
- [39] M. Cvetic and J. Wang, Phys. Rev. D61 124020 (2000) [arXiv:hep-th/9912187]
- [40] P. Avelino, C. Martins, and J. Oliveira, Phys. Rev. D72 083506 (2005) [arXiv:hep-ph/0507272]
- [41] R. P. Naidu et al., ApJL 940 L14 (2022) [arXiv:2207.09434]
- [42] N. J. Adams et al., MNRAS 518 4755 (2023) [arXiv:2207.11217]
- [43] H. Yan et al., ApJL 942 L9 (2023) [arXiv:2207.11558]
- [44] I. Labbé et al., Nature **616** 266 (2023).