Measurement of collective dynamics in small and large systems with the ATLAS detector

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Investigating collectivity in proton-proton (pp) and proton-nucleus (p + A) collisions is essential for understanding the mechanism of formation of Quark Gluon Plasma (QGP). In pp collisions at $\sqrt{s_{NN}} = 13$ TeV, ATLAS has reported no impact on flow coefficients when excluding low-$p_T$ jets, indicating collectivity in small systems does not originate from hard processes. Comparing flow decorrelation in Xe+Xe and pp collisions at $\sqrt{s_{NN}} = 5.44$ TeV and 5 TeV respectively, underscores the role of sub-nucleonic fluctuations in determining longitudinal energy deposition. ATLAS has also reported a significant correlation between event-wise flow and transverse momentum fluctuations in Xe+Xe and Pb+Pb collisions at $\sqrt{s_{NN}} = 5.44$ TeV and 5 TeV respectively, emphasizing the role of nuclear structure in accurately describing bulk observables in heavy-ion collisions and provides the first experimental evidence from a high-energy experiment for a significant triaxiality in $^{129}$Xe.
1. Introduction and Measurements

Heavy-ion collisions at the Large Hadron Collider produce QGP, whose space-time evolution is well described by relativistic viscous hydrodynamics. Driven by the large pressure gradients, the QGP expands rapidly in the transverse plane and converts the spatial anisotropy in the initial state into momentum anisotropy in the final state. The collective expansion in each event is quantified by a Fourier expansion of particle distribution in azimuth given by $\frac{dN}{d\phi} = \frac{N}{2\pi} (1 + 2 \sum_{n=1}^{\infty} v_n \cos n(\phi - \Phi_n))$, where $v_n$ and $\Phi_n$ represent the amplitude and phase of the $n$th-order azimuthal flow vector $V_n = v_n e^{i\Phi_n}$. Collectivity in this context means that the produced particles exhibit a common property, such as a common velocity field or a common direction [1]. Model calculations show that the $V_n$ is approximately proportional to initial state eccentricity $E_n$ for $n = 2$ and 3, as well as for $n = 4$ in central collisions [2]. In addition, the fluctuations in the size of the overlap area in the initial state give rise to fluctuations in radial flow, which, in turn, lead to event-by-event fluctuation of the average transverse momentum ($p_T$).

The most intriguing feature of the azimuthal anisotropies is the “ridge” behavior, which is an enhancement in the production of particles with small azimuthal angle ($\phi$) separation extending over a large range of pseudorapidity ($\eta$) separation [3]. The presence of the ridge in heavy-ion collisions is typically attributed to hydrodynamic flow and is viewed as a signature of collective behavior and the formation of QGP. In smaller collision systems such as $p+A$ or $p+p$, the presence of the ridge challenges the notion that collective behavior is exclusive to QGP formation and has led to several phenomenological models attributing the origin of long-range behavior to hard or semi-hard processes. In these processes, multi-particle correlations between outgoing partons may arise due to saturation of the parton configurations in the incident hadrons [4]. If the long-range correlations arise due to hard or semi-hard processes, removing particles associated with jets from the analysis would weaken the long-range correlation.

To investigate the influence of hard processes on the origin of collectivity in small systems, ATLAS has probed the role of semi-hard and hard processes on collective behavior observed in small systems [5]. The events were systematically categorized into five distinct groups based on their two-particle correlations, each representing varying contributions from particles produced by jets, all falling within a $p_T$-range of 4 GeV. These five categories are: $h - h$, which denotes the analysis of all particles; $h^{UE} - h^{U}$, indicating correlations between one particle originating from a jet and another from the underlying event; and three classes pertaining to 2-Particle Correlations between particles from the underlying event: AllEvents, where the analysis includes tracks within one unit in $\eta$ from any jet above the chosen threshold of 10 GeV, which are excluded from the 2-Particle Correlation analysis; NoJet, where the analysis is carried out on events that do not possess a single jet with transverse momentum ($p_T$) greater than the chosen threshold, thereby encompassing events primarily governed by soft processes; and WithJet, representing the analysis performed on events featuring at least one jet with $p_T$ greater than the chosen $p_T$ threshold.

Figure 1 compares measured $v_2$ with varying contributions from jets. It is evident that both integrated $v_2$ and differential $v_2(p_T)$ remain unaltered by the presence or absence of jets. Consequently, this measurement led to the conclusion that hard scatterings do not play a role in generating long-range correlations, effectively ruling out their contribution to the ridge phenomenon.

To explore the impact of sub-nucleonic fluctuations on collectivity, ATLAS has measured
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Figure 1: (Left) $p_T$ dependence of the $v_2$ obtained for the 40–150 multiplicity interval. (Right) $v_2$ as a function of the (efficiency corrected) multiplicity. The data-points for the Inclusive sample are drawn at the nominal values, AllEvents and NoJet points are shifted slightly for clarity. Data-points for the WithJet sample are evaluated over coarser multiplicity intervals to reduce statistical uncertainties. The error bars and shaded bands correspond to statistical and systematic uncertainties, respectively [5].

decorrelation in smaller systems. Decorrelation refers to the reduction in the strength of flow-correlation with increasing $\eta$ gaps between correlated particles. In this context, peripheral Xe+Xe and $pp$ collisions provide simplified scenarios for understanding the longitudinal structure of QGP [6]. Decorrelation arises from the fact that in heavy ion collisions, owing to initial state fluctuations, deposited energy and the transverse shape of the fireball is not boost invariant. Model studies have shown that preferential emission by forward and backward-moving participant nucleons produce particles preferably in the forward direction, leading to an event-by-event torqued fireball. This leads to an asymmetry of flow magnitude as well as a “Twist” of the symmetry plane in the longitudinal direction on an event-by-event basis [7, 8]. Previous measurements conducted by the ATLAS collaboration had revealed that state-of-the-art hydrodynamic models are inadequate in describing the longitudinal structure of initial energy deposition[9].

Figure 2 shows the comparison of the decorrelation signal ($F_2$) for $pp$ and Xe+Xe collisions [6]. At the lowest multiplicities, the decorrelations observed in $pp$ collisions are close to those in Xe+Xe collisions, implying that Xe+Xe events with low multiplicities are predominantly characterized by single nucleon-nucleon configurations. Conversely, as multiplicities increase, the observed $F_2$ values in Xe+Xe collisions diminish, signifying a distinct mechanism for additional particle production in these two systems. As a result, this study concludes that the correlation between the geometry of the initial state and overall particle production exhibits variations at sub-nucleonic scales compared to nucleonic scales. Furthermore, the predictions from the AMPT model, which involves two-color strings spanning a significant longitudinal extent in $pp$ collisions, do not align well with the data. This suggests a need for further refinement and fine-tuning of longitudinal modeling to describe the data accurately.

In addition to generating anisotropic flow, the collective response to the overall transverse size ($R$) in the initial state also leads to a large “radial flow”, reflected by an increase of the $[p_T]$. Any correlated fluctuations between the $E_n$ and $R$ in the initial state are expected to generate a dynamical correlation between $v_n$ and $[p_T]$ in the final state. A three-particle correlator has been proposed to
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Figure 2: A comparison of AMPT theory calculations to the 5 TeV \( p\bar{p} \) (left) and 5.44 TeV Xe+Xe (right) results. In data, the raw \( F_2 \) values are shown as a function of \( N_{\text{rec}} \). Additionally, the subtracted \( F_2 \) values are shown. The statistical uncertainties are drawn as vertical lines. The systematic uncertainties are plotted as bands [6].

quantify this correlation [10]:

\[
\rho_n = \frac{\langle \langle v_n^2 \delta p_T \rangle \rangle}{\sqrt{\langle \langle v_n^2 \rangle \rangle} \sqrt{\langle \langle \delta p_T \delta p_T \rangle \rangle}},
\]

where the averages are over events with similar particle multiplicity.

An investigation on the system-size dependence of \( v_n - [p_T] \) correlation is performed in \(^{129}\text{Xe} + ^{129}\text{Xe} \) collisions and comparing them with \(^{208}\text{Pb} + ^{208}\text{Pb} \) collisions at ATLAS [11]. Recent measurements show that the \( \rho_n \) exhibits significant differences between these two systems, especially in the central collisions due to the deformation of nuclear structure in \(^{129}\text{Xe} \). Deformations in nuclear structure is generally parametrized using \( R(\theta, \phi) = R_0 (1 + \beta [\cos \gamma Y_{2,0} + \sin \gamma Y_{2,2}] ) \) where \( R_0 \) is the nuclear radius, \( Y_{l,m} \) are spherical harmonics, \( \beta \), and \( \gamma \) are quadrupole deformation parameters. \( \beta \) is the magnitude of overall deformation, while \( \gamma (0 \leq \gamma \leq 60^\circ) \) describes the length imbalance between the three axes of the spheroid, called “Triaxiality”. Recent models indicate \( \rho_2 \approx a + b \cos(3\gamma) \beta^3 \) [12], where, \( a \) and \( b \) are values for spherical nuclei. In addition to these initial-state-driven long-range global correlations, the \( v_n - [p_T] \) correlation measurement may have short-range “non-flow” correlations from resonance decays and jets. The non-flow correlation is suppressed by requiring correlation between particles from different subevents separated in \( \eta \) [13].

In peripheral centralities, a double sign change in \( \rho_2 \) is predicted as evidence of the existence of initial momentum anisotropy [14].

Figure 3 displays \( \rho_2 \) for Pb+Pb collisions versus \( N_{\text{ch}}^{\text{rec}} \) compared with different models incorporating initial and final state dynamics. In the 0–10% centrality interval, where the effects of nuclear deformation are important, all models generally show reasonable agreement with each other and with the data. In particular, the Trajectum model quantitatively reproduces the ordering between \( 0.5 < p_T < 2 \text{ GeV} \) and \( 0.5 < p_T < 5 \text{ GeV} \). In the peripheral collisions, all model predictions for \( \rho_2 \) show a sharp decrease and a sign-change, qualitatively consistent with the ATLAS data.

Figure 3 also compares \( \rho_2 \) (0–20% centrality) with Trento model to constrain \( \gamma_{\text{Xe}} \) [15]. Due to large quadrupole deformation in \(^{129}\text{Xe} (\beta_{\text{Xe}} \sim 0.2) \), \( \rho_2 \) is sensitive to \( \gamma_{\text{Xe}} \). The Trento model
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Figure 3: (Left) $\rho_2$ values in Pb+Pb collisions in two $p_T$-ranges and $|\eta| < 2.5$ compared with various models: Trento and Trajectum models in solid lines and v-USPhydro and IP-Glasma+MUSIC hydrodynamic models in shaded bands, which represent the statistical uncertainties of the model calculations. (Center) Comparison of $\rho_2$ in Xe+Xe and Pb+Pb collisions with the Trento model for various quadrupole deformation parameter values in $0.5 < p_T < 2$ GeV as a function of centrality. (Right) Comparison of $\rho_2$ ratios, $\rho_2, \text{Xe}+\text{Xe}/\rho_2, \text{Pb}+\text{Pb}$, with the Trento model for various quadrupole deformation parameter values in two $p_T$-ranges. The Trento model results are connected by lines for better visualization [11].

Predictions agree with the $\rho_2$ measured for $0.5 < p_T < 2$ GeV for both Pb+Pb and Xe+Xe collisions for appropriate values of deformation parameters. To nullify $p_T$ dependence, ratios of $\rho_2$ between Xe+Xe and Pb+Pb are computed for two $p_T$-ranges and compared in Figure 3. In 10–20% centrality, where triaxiality has a minor impact, the model and data ratio closely match. In 0–10% centrality, where $\rho_2$ strongly depends on triaxiality, the comparison favors $\gamma_{\text{Xe}} \sim 30^\circ$. Thus, ATLAS provided the first experimental evidence from the high-energy collision of a significant triaxiality in $^{129}\text{Xe}$ using $\rho_2$. A recent model study underscores the accuracy of extracting $\gamma$ only when nuclear fluctuations in $\gamma$ are small [16]. Consequently, the $\rho_{(2, \text{Xe})}/\rho_{(2, \text{Pb})}$ ratio supports the triaxial shape of $^{129}\text{Xe}$, but only under the condition that the fluctuations in $\gamma_{\text{Xe}}$ are not extensive.

Figure 4: The centrality dependence of $\rho_2$ in Pb+Pb collisions in the peripheral region of 60–84% for the standard method (left), two-subevent method (middle) and three-subevent method (right), compared between the $N_{\text{rec, ch}}$-based and $\Sigma E_T$-based event-averaging procedures and two $\eta$-ranges. The error bars and shaded boxes represent statistical and systematic uncertainties, respectively [11].

Figure 4 illustrates $\rho_2$ values in Pb+Pb collisions for particles originating from various $\eta$ regions and employing different subevent methods. Additionally, it demonstrates measurements
as a function of particle production in central versus forward rapidities, denoted as $N_{\text{ch}}^{\text{rec}}$ and $\Sigma E_T$ respectively. While some hints of a double sign change in $\rho_2$ in peripheral centralities are observed, it is inconclusive due to the limited range of centrality under study. The extent of the sign change depends on factors like centrality determination methods based on $N_{\text{ch}}^{\text{rec}}$ or $\Sigma E_T$ and the presence of short-range correlations like jets. Future measurements in smaller systems may help isolate these effects and support the search for evidence of initial momentum anisotropy in heavy-ion collisions.

2. Summary

In summary, ATLAS has demonstrated using $pp$ collisions at $\sqrt{s_{\text{NN}}} = 13$ TeV, that hard scatterings do not contribute to long-range correlations, ruling out their role in ridge behavior. Decorrelation measurements in $pp$ collisions at $\sqrt{s_{\text{NN}}} = 5$ TeV and Xe+Xe collisions at $\sqrt{s_{\text{NN}}} = 5.44$ TeV show that the correlation between the initial-state geometry and overall particle production differs at sub-nucleonic scales compared to nucleonic scales. Through examinations of Pb+Pb collisions at $\sqrt{s_{\text{NN}}} = 5$ TeV and Xe+Xe collisions at $\sqrt{s_{\text{NN}}} = 5.44$ TeV, ATLAS has provided the first experimental evidence from the high-energy collision of a significant triaxiality in $^{129}$Xe using $\rho_2$ and constraints $\gamma_{\text{Xe}} \sim 30^\circ$. In addition, the $\rho_2$ also shows a hint of double sign change behavior in peripheral centralities similar to those expected from the existence of initial state momentum anisotropy. Further investigation in smaller systems is necessary for more conclusive evidence owing to the existence of additional effects like centrality fluctuations and non-flow. These comprehensive findings provide crucial contributions to constrain the modeling of heavy-ion collisions and to understand the origin of collective behavior in the QGP medium.

References


