

Unveiling generalized parton distributions through the pseudo-distribution approach

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Understanding the intricate three-dimensional internal structure of the nucleon has been a long-standing challenge. The main quantitative tool to map out this structure are the generalized parton distributions (GPDs). In these proceedings, we present our exploratory results from applying the pseudo-distribution approach on the lattice, showing the unpolarized isovector E GPD as an example. We use one ensemble of $N_f = 2 + 1 + 1$ twisted mass fermions with a clover term, at a non-physical pion mass of 260 MeV and a lattice spacing of 0.093 fm.

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1. Introduction

The precise structure of the most fundamental baryon, the nucleon, which is the building block of all matter, still eludes our understanding. Unraveling the composition of hadrons has been a persistent quest for all particle physicists, theoreticians and experimentalists alike, for decades. The partonic description of nucleons in the form of parton distribution functions (PDFs) is relatively well studied and explored, but its higher-dimensional counterpart, generalized parton distributions (GPDs), and tomographic imaging remain largely uncharted. Current and upcoming experiments aim to shed light on this structure [1]. Complementing these experimental pursuits, first-principle investigations on the lattice play a crucial role.

Historically, lattice calculations primarily focused on determining different form factors, specifically moments of partonic distributions. However, addressing the full x -dependence of these distributions proved challenging due to the limitations imposed by the Euclidean spacetime metric. Around a decade ago, the groundbreaking concept of computing quasi-distributions marked a significant leap forward in overcoming these challenges [2]. This innovative approach involves determining these quasi-distributions using the lattice-calculable spatial correlations of a boosted hadron. Said distributions are later translated into relevant, physical (Minkowski) ones. This approach necessitates lattice observables that are not only renormalizable but also share the same infrared structure as their physical counterparts. Substantial theoretical and numerical achievements have been reached utilizing this approach, as detailed in various reviews, see e.g. [3–6].

There exists an alternative approach known as pseudo-distributions [7, 8]. This approach can be used on the same lattice data as quasi-distributions, but generates fundamentally different systematic effects. The crucial disparity arises from the manner in which they factorize into their physical distributions. Quasi-distributions perform this in momentum space, whereas pseudo-distributions do so in coordinate space. Despite this fundamental difference, both methodologies converge to the same physical distributions in the infinite momentum limit, provided all lattice-specific and other systematic effects are accounted for.

Previous work utilizing the pseudo-distribution approach concentrated on PDFs, see e.g. [9–12]. Extension of pseudo-PDFs to pseudo-GPDs was developed by Radyushkin [8, 13] and has not been subject to detailed lattice investigations so far. In these proceedings, we share our results for the first exploration of the GPDs utilizing this approach. As an example, we concentrate on the unpolarized E GPD in the flavor non-singlet case ($u - d$).

2. Theoretical background

2.1 Bare matrix elements

The general form of the Euclidean bare matrix elements (MEs) for the vector case is

$$F_V^\mu(z, P_f, P_i) = \langle N(P_f) | \bar{\psi}(z) \gamma^\mu W(0, z) \psi(0) | N(P_i) \rangle, \quad (1)$$

where P_i and P_f are the initial and final nucleon's four-momenta and $W(0, z)$ is the Wilson line taken in the 3-direction and of length z . These matrix elements are obtained from a suitable ratio of three-point and two-point functions, see [14] for details. We work with one value of the source-sink

separation, $t_s = 10a$, and extract the MEs using the plateau method in the range of operator insertion times $3a \leq \tau \leq 7a$, which has been found to be robust at this level of statistical precision [14, 15].

2.2 GPDs in coordinate space

We calculate our bare MEs in asymmetric frames of reference, in which $P_f = (0, 0, P_f^3)$ and $P_i = (-\Delta^1, -\Delta^2, P_i^3)$, where Δ^μ is the momentum transfer and $P_i^3 = P_f^3 = P^3$ (zero-skewness case, $\Delta^3 = 0$). The extraction of GPDs from such asymmetric frames was discussed extensively in [16], where Lorentz-invariant amplitudes were introduced to parametrize MEs. We recall that there are eight such linearly-independent amplitudes,

$$F_V^\mu(z, P_f, P_i) = \bar{u}(P_f, \lambda') \left[\frac{P^\mu}{m} A_1 + m z^\mu A_2 + \frac{\Delta^\mu}{m} A_3 + i m \sigma^{\mu z} A_4 + \frac{i \sigma^{\mu \Delta}}{m} A_5 + \frac{i P^\mu \sigma^{z \Delta}}{m} A_6 + i m z^\mu \sigma^{z \Delta} A_7 + \frac{i \Delta^\mu \sigma^{z \Delta}}{m} A_8 \right] u(P_i, \lambda). \quad (2)$$

Additionally, we use four parity projectors Γ_κ for each of the four gamma insertions γ_μ (see [16] for details), obtaining 16 different matrix elements. Eight of these are independent upon averaging equivalent contributions with reversed roles of Δ^1 and Δ^2 . Thus, we obtain uniquely determined eight amplitudes. Out of these, one can construct the Lorentz-invariant definition [16] of the E GPD in coordinate space that we choose for this work,

$$E = -A_1 + 2A_5 + 2zP^3 A_6. \quad (3)$$

2.3 Renormalization and matching

Once the coordinate-space Euclidean GPDs (pseudo-GPDs), $E(\nu, z)$ (where $\nu = P^3 z$ is the Ioffe time), have been constructed from the amplitudes, our next step is to renormalize them. Bare MEs are riddled with Wilson-line-induced power divergences and standard logarithmic ones. Due to these being multiplicatively renormalizable [17], they can be eliminated using a ratio scheme, which is created using a combination of zero-momentum and local PDFs [9]. We employ the double ratio, defined as

$$\mathcal{E}(\nu, z) = \frac{E(\nu, z)/M(\nu, 0)}{M(0, z)/M(0, 0)}, \quad (4)$$

where $M(\nu, z)$ is the unpolarized PDF. The double ratios are commonly referred to as reduced Ioffe-time distributions (ITDs). These reduced ITDs are Euclidean objects calculated at different scales $1/z$ and hence, they have to be evolved to a common scale and translated to Minkowski spacetime and the $\overline{\text{MS}}$ renormalization scheme. We employ the one-loop matching [18–20],

$$Q_E(\nu, \mu^2) = \mathcal{E}(\nu, \mu^2) - \frac{\alpha_s}{\pi} \int_0^1 du C(u, z, \mu^2) \left(\mathcal{E}(u\nu, \mu^2) - \mathcal{E}(\nu, \mu^2) \right), \quad (5)$$

where $C(u, z, \mu^2) = (C_F/2) (L(u) + B(u) \ln(z^2 \mu^2 e^{2\gamma_E+1})/4)$. The function $B(u) = \frac{1+u^2}{u-1}$ handles the evolution of the ITD to a common scale μ (taken to be 2 GeV), and the function $L(u) = 4 \frac{\ln(1-u)}{u-1} - 2(u-1)$ matches the ITD to the light-cone and performs the scheme conversion. The resulting quantity, $Q_E(\nu, \mu^2)$, is the light-cone $\overline{\text{MS}}$ -renormalized E GPD in coordinate space.

2.4 Reconstruction of the x -dependence

The matched ITDs are related to the momentum-space GPDs, $q(x, \mu^2)$, by a Fourier transform, $q(x, \mu^2) = \int_{-\infty}^{\infty} dv e^{-i\nu x} Q_E(\nu, \mu^2)$. However, the Fourier transform assumes an infinite range of continuous ITDs at all Ioffe times, while the lattice data are fundamentally discrete. This creates an inverse problem, which one can tackle in various ways [21].

In this work, we first consider the real and imaginary part of the ITDs separately,

$$\text{Re } Q_E(\nu, \mu^2) = \int_0^1 dx \cos(\nu x) q_\nu(x, \mu^2) \quad (6)$$

$$\text{Im } Q_E(\nu, \mu^2) = \int_0^1 dx \sin(\nu x) q_{\nu 2s}(x, \mu^2), \quad (7)$$

where q_ν is the valence distribution, $q_\nu = q - \bar{q}$, and $q_{2\nu s} = q_\nu + 2\bar{q}$. We assume a fitting ansatz of the form $q(\mu, x) = Nx^a(1-x)^b$. The valence distribution is normalized to $Q_E(0, \mu^2)$ and hence $N = Q_E(0, \mu^2)/B(a+1, b+1)$ for this case, where B is the Euler beta function. For $q_{\nu 2s}$, the normalization is not fixed and N is a fitting parameter, in addition to the parameters a, b fitted in both cases. The fitting parameters are found by minimizing the χ^2 function, $\chi^2 = \sum_{\nu=0}^{\nu_{\max}} ((Q_E(\nu, \mu^2) - Q_E^F(\nu, \mu^2))/\sigma_{Q_E(\nu, \mu^2)})^2$, where the $Q_E^F(\nu, \mu^2)$ is the cosine/sine Fourier transform of the x -dependent fitting ansatz and σ_{Q_E} is the error of Q_E . The sum over Ioffe times is taken to be $\nu_{\max} = P_{\max}^3 z_{\max}$ and we discuss the choice of z_{\max} below, while P_{\max}^3 is fixed by the maximum nucleon boost in our data.

3. Lattice Setup

We use one ensemble of $N_f = 2 + 1 + 1$ twisted mass fermions with a clover term, at a lattice spacing of $a = 0.093$ fm and a lattice volume of $T \cdot L^3/a^4 = 64 \cdot 32^3$. The pion mass is roughly twice the physical one, $m_\pi \approx 260$ MeV.

The pseudo-distribution approach is based on the same MEs as quasi-distributions. Hence, we reuse the asymmetric-frame data of [16], which are limited to a single value of $P^3 = 6\pi/L$. To fully realize the potential of the pseudo-distribution approach, data at other values of P^3 can be used and additionally, zero-momentum MEs are needed. Thus, we generated additional data corresponding to $P^3 = \{0, 1, 2, 4\} \cdot 2\pi/L$. The dominating computational effort concerned, obviously, the largest-momentum case, which allowed us to significantly enhance the covered range of Ioffe times. The details of our setup are shown in Table 1.

In these proceedings, we present results only for one value of the momentum transfer, $-t = 0.64 \text{ GeV}^2$, which is obtained from eight equivalent kinematic setups, i.e. momentum transfer 3-vectors $\vec{\Delta}$ taken as $(2, 0, 0)$, $(-2, 0, 0)$, $(0, 2, 0)$ and $(0, -2, 0)$ at both signs of P^3 . As hinted above, we present only the case of the isovector E GPD as an example.

4. Results

In Figure 1, we present the matched ITDs calculated from (5). A given Ioffe time ν can be obtained from different combinations of (P^3, z) and, in principle, ITDs from such combinations should be equivalent after they have been evolved to a common scale and matched. This holds,

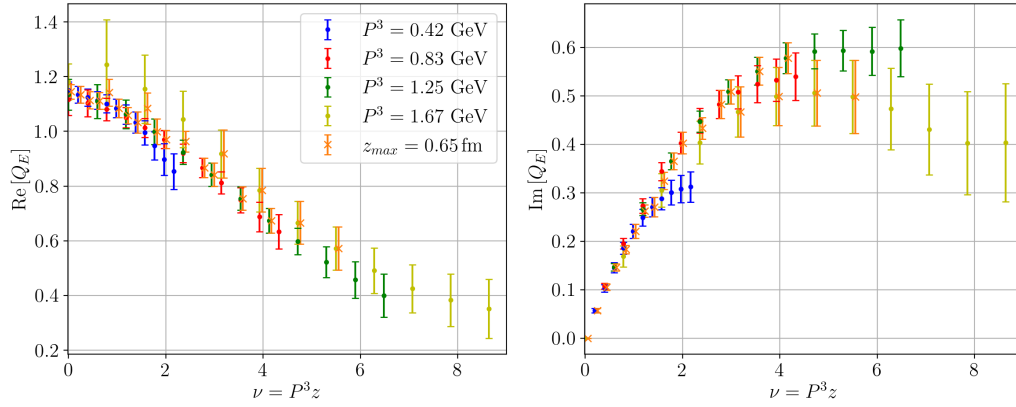


Figure 1: Matched E ITDs at four nucleon momentum values (circles) and the averaged distribution (crosses), slightly shifted for better visibility.

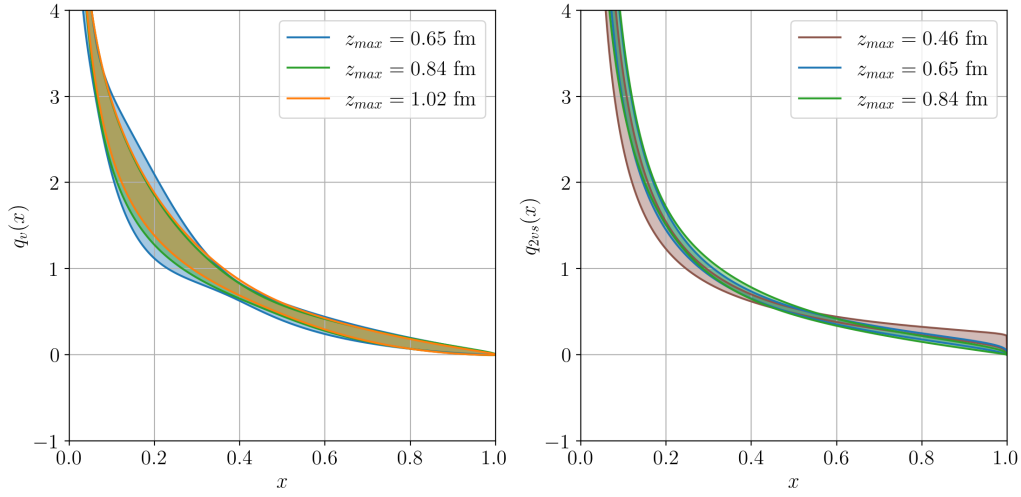


Figure 2: Left: $q_v(x)$, right: $q_{2vs}(x)$, reconstructed with different z_{max} values.

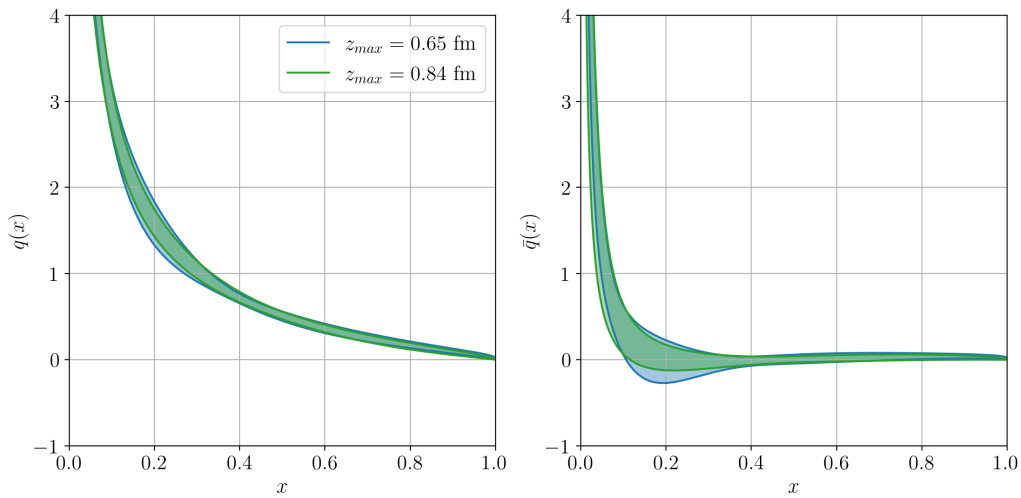


Figure 3: Left: $q(x)$, right: $\bar{q}(x)$, reconstructed with different z_{max} values.

P^3 [GeV]	N_{confs}	$N_{\text{kinematics}}$	$N_{\text{sourcepos.}}$	$N_{\text{measurements}}$
0.0	404	1	8	3,232
0.42	100	8	8	6,400
0.83	100	8	8	6,400
1.25	269	8	8	17,216
1.67	404	8	32	103,424

Table 1: Details of our lattice setup. Given is the nucleon boost, the number of used gauge field configurations, the number of kinematic setups, the number of source positions and the total number of measurements included in our statistics. New data generated for this work are highlighted in red and the ones in black correspond to data used previously for quasi-GPDs in [16].

provided that the evolution/matching procedure is performed at a sufficiently small value of z . Thus, the challenge is to find a suitable maximum value for z . We adopt the pragmatic criterion that z_{max} is the maximum value for which different (P^3, z) pairs still lead to consistent values. Then, we average ITDs corresponding to the same Ioffe times and an example of this process is depicted in Figure 1, with circles pertaining to actual lattice data at a given boost and crosses to the averages over (P^3, z) combinations (with data up to $z_{\text{max}} = 0.65$ fm entering the averages), where applicable. The agreement between (P^3, z) pairs persists until $z/a \approx 10/8$ for the real/imaginary part. Overall, we explore cases with $z_{\text{max}} \in \{0.65, 0.84, 1.02\}$ fm for the real part and for the imaginary part, displaying a tendency for larger effects, we look at more conservative set of values $z_{\text{max}} \in \{0.47, 0.65, 0.84\}$ fm.

The final x -reconstructed distributions are presented in Figures 2 (left: q_v , right: $q_{v2s} = q_v + 2\bar{q}$) and 3 (left: $q = q_v + \bar{q}$, right: \bar{q}). For the latter, we show results from common z_{max} values for the real/imaginary part data, although in principle, we could combine results from different z_{max} (we leave the analysis of the mixed cases for future work). In any case, the obtained z_{max} -dependence is relatively small, with all z_{max} values giving consistent x -dependent GPDs. The obvious effect of decreasing z_{max} is the inflation of the error from missing parts of the Ioffe time range that constrain the behavior in x -space.

The results for the “full” distribution, $q = q_v + \bar{q}$, are qualitatively and to a large extent also quantitatively similar to the ones from the quasi-distribution approach [16]. Larger differences are observed in the sea distribution, with quasi-GPDs leading to its significantly non-zero values already at $x \approx 0.5$. This points to the likely possibility that systematic effects evinced by both approaches are similar in the valence sector and enhanced for sea quarks. However, the final conclusions can be reached only after a more systematic analysis, which is our current work in progress. In our upcoming publication, we will report results for a wide range of momentum transfers for both E and H GPDs, with assessment of various systematic effects, such as the choice of the fitting ansatz. Meanwhile, the presented example is the first exploration of pseudo-GPDs, with unambiguously encouraging results.

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