

Constraining Higgs sectors of BSM models - the case of 95 GeV "Higgs"

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In view of lack of the direct experimental evidence for a Beyond the Standard Model (BSM) physics, accommodating a SM-like Higgs boson is one of the most important constraints that a BSM model must fulfill. Already for some time the FlexibleSUSY spectrum generator allowed for a reliable prediction of masses and decay patterns of the BSM Higgs boson in a large class of user defined supersymmetric and non-supersymmetric models. However, no easy way to compare those predictions with experimental data existed. To that end we present here an interface between FlexibleSUSY and HiggsTools, a computer program assessing in a statistically meaningful way consistency of a BSM Higgs sector with experiments.

Motivated by the recent ATLAS and CMS observation of the di-photon excess at a mass of ~ 95 GeV we demonstrate the capabilities of our framework by investigating whether the observed low mass excesses around 95 GeV seen in the data can be explained as the lightest scalar of the Minimal R-symmetric Supersymmetric Standard Model, without spoiling the SM-like properties of the second-to-lightest state. We also briefly comment on the light dark matter candidate which is a necessarily ingredient of such a setup.

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1. Introduction

The discovery of a 125 GeV Higgs boson at the LHC in 2012 and the subsequent precise measurement of its properties put a strong test on any potential Beyond the Standard Model (BSM) theory. However, to capitalize on this discovery one has to be able to predict those observables in a BSM model with an equally impressive accuracy. While precise computations of various observables existed for a selection of popular BSM models, the analysis of any other extension was troublesome.

To that end `FlexibleSUSY` [1, 2] was developed to enable automatized phenomenology in a large class of user defined BSM models. `FlexibleSUSY` capabilities include, among many other observables, precise predictions of Higgs boson masses. Furthermore, `FlexibleDecay` [3] extended `FlexibleSUSY` capabilities by allowing to compute decays of scalar particles. The special emphasis was put on the Higgs-like states, with the goal being to reach precision comparable with current experimental measurements. Here we describe a recent update to `FlexibleDecay` module. To streamline the application of experimental constraints from Higgs searches we created an interface to `HiggsTools` [4], a computer program which encodes searches for additional scalars and measurements of the detected Higgs boson at 125 GeV. This allows to derive a statistically meaningful information on the validity of a given parameter point in relation to currently available experimental data.

We illustrate the application of our work by investigating an intriguing question of whether there are additional Higgs-boson like scalar states, possibly with mass smaller than 125 GeV. Such scalar particles below the SM Higgs boson mass have been searched for in the past at LEP and Tevatron and now at the LHC, showing some promising deviations from the SM.

Such states occur frequently in BSM theories. For example, it has been show that in the Minimal R-symmetric Supersymmetric model (MRSSM) [5] the SM-like Higgs boson can be realised as a second-lightest scalar particle leaving room for a lighter state below 125 GeV [6]. The MRSSM by itself is an attractive alternative to the MSSM, elegantly solving some of its shortcomings while also significantly altering its phenomenology. The absence of left-right squark mixing eliminates many parameters responsible for flavour-violating interactions and thus alleviates the SUSY flavour problem. Its various aspects, including Higgs physics, electroweak precision observables, flavour physics, Dirac gauginos, color octet scalars and dark matter have been studied in a long series of papers [6–25].

In what follows we focus on LEP and LHC hints of extra scalar state at around 95 GeV [26–29]. As the LEP excess is very broad, both LEP and LHC cases can be addressed simultaneously by a single state which we will identify with a mostly singlet-like state from the MRSSM.

2. HiggsTools interface

A convenient way of parametrizing the BSM Higgs sector is via effective couplings κ [30],

$$\kappa_i^2 \equiv \frac{\Gamma(H_i \rightarrow AB)_{\text{BSM}}}{\Gamma(h \rightarrow AB)_{\text{SM}, m_h = m_{H_i}}}, \quad (1)$$

defined as ratios of BSM partial widths to SM widths for the SM Higgs of the same mass. Such couplings can be then used as input to `HiggsTools`, which then internally recomputes partial

widths and cross sections and quantifies the viability of a given parameter point in relation to current experimental data.

In this section we describe the construction of such effective couplings in FlexibleSUSY and the process of linking it with HiggsTools. This interface is planned to be part of 2.9.0 release of FlexibleSUSY (the proper paper documenting it is in preparation) while a preliminary version can be already obtained from [github](#).

2.1 Construction of effective couplings

For the calculation of effective couplings in Eq. (1) one has to construct a hypothetical SM with the Higgs boson of mass equal to the BSM Higgs state. FlexibleSUSY already contains a built-in SM in addition to the BSM model created by the user. For the construction of such a SM equivalent state, a quartic coupling parameter λ is tuned to provide the SM Higgs mass m_h equal to a BSM Higgs mass m_{H_i} (for CP-even Higgs) or m_{A_i} (for CP-odd Higgs). The case with undefined CP properties (as happens in CP-violating models) is also correctly handled. This procedure is done separately for every state H_i or A_i . Since for $m_h \gtrsim 650$ GeV the quartic coupling λ enters a non-perturbative regime, this procedure limits the application of our interface to BSM Higgs masses in the range $[1, 650]$ GeV. BSM states beyond that mass range are not passed to HiggsTools.

2.2 HiggsTools interface

HiggsTools allows for two ways to compare theory prediction with experimental data, by either specifying cross-sections and branching ratios directly or providing effective coupling description. As calculation of cross-sections is beyond the scope of codes like spectrum generators we choose to use the effective coupling input.

The procedure of calling HiggsTools is transparent to the user. The user only has to point to the location of HiggsTools during configuration of FlexibleSUSY using two new configure flags

```
$ ./configure --with-higgstools-include=... --with-higgstools-libdir=...
```

which are used to specify the location of HiggsTools header files and library, respectively, which we mark with dots in the above code snippet. Those should be replaced with actual paths depending on HiggsTools installation on user's computer.

At runtime, the user has to specify the location of HiggsSignals and HiggsBounds databases using two new command line options (we use the single extended SM (SSM) in this example call):

```
$ ./run_SSM.x --higgsbounds-dataset=... --higgsignals-dataset=...
```

We refer to HiggsTools documentation [4] as to how to obtain and install both the library and respective databases.

The output of HiggsTools is then shown in two new SLHA-like [31, 32] blocks printed out by FlexibleSUSY. The new HiggsSignals block looks for example as follows

```
Block HIGGSIGNALS
  1      1.59000000E+02  # number of degrees of freedom
```

2	1.68975284E+02	# chi^2
3	1.52489682E+02	# SM chi^2 for mh = 125.090000 GeV
4	2.63146088E-04	# p-value

and contains in the first two lines the number of degrees of freedom from the number of observables in HiggsTools database and the χ_{BSM}^2 for the BSM model reported by HiggsTools. To allow for a quick interpretation of results we also provide as a reference a χ_{SM}^2 for a SM and p -value computed for $\chi_{\text{BSM}}^2 - \chi_{\text{SM}}^2$ difference assuming a $2d$ distribution. This can be enough in many cases to determine the validity of a concrete parameter point. If more sophisticated statistical treatment is needed, one can refer directly to χ_{BSM}^2 .

For the HiggsBounds we provide the ratios r to the 95% excluded cross sections in the second new block, which may look like

Block HIGGSBOUNDS			
25	1	4.03838082E-01	# LEP [eeHZ>bb] from hep-ex/0306033 (LEPComb 2.5fb-1, M=(12, 120))
25	2	8.36015679E-01	# expRatio
35	1	6.60215252E-01	# LHC8 [vbfH,HW,Htt,H,HZ]>[bb,tautau,WW,ZZ,gamgam] from CMS-PAS-HIG-12-045 (CMS 17.3fb-1, M=(110, 600))
35	2	3.32131274E+00	# expRatio

The first column contains the PDG ID of a BSM Higgs-like state while the second distinguishes between the observed (1) and expected (2) limits. The analysis providing the strongest exclusion is given in the SLHA comment of entry 1. Points with $r_{\text{obs}} > 1$ should be considered by the user as excluded at 95% C.L..

3. Application: The case of 95 GeV Higgs in the MRSSM

We now briefly discuss an example application of the interface described in the previous section. During recent years, a collection of evidence hinting at a low mass, Higgs-like resonance between 95–100 GeV emerged. These excesses can be accommodated for example in models where the SM-like Higgs mixes with an extra gauge-singlet field, as is the case in the MRSSM.

In this section we summarize the aforementioned hints and show how they can be realised in the MRSSM. Thanks to the work described in the previous section, we prove that we still obtain a SM-like second-to-lightest Higgs boson. We also comment on light dark matter, a necessary prediction of such a setup and general collider constraint. This analysis is described in details in Ref. [33] to which we refer the reader for more details.

3.1 Hints for a 95 GeV Higgs

There is a long standing LEP anomaly observed in the $Zb\bar{b}$ final state in the $95 \lesssim m_{b\bar{b}} \lesssim 100$ invariant mass window [26]. This anomaly can be explained for example by a scalar state s with a mass of 98 GeV whose combined production and branching ratio is roughly an order of magnitude

smaller than that of a SM-like Higgs h of the same mass [34]

$$\mu_{Zb\bar{b}} = \frac{\sigma(e^+e^- \rightarrow Z^*s \rightarrow Zb\bar{b})}{\sigma(e^+e^- \rightarrow Z^*h_{98}^{\text{SM}} \rightarrow Zb\bar{b})} = 0.117 \pm 0.057. \quad (2)$$

The data collected by CMS in Run 1 and the first year of Run 2 showed a di-photon excess, recently confirmed based on full Run 2 data set together with a small deviation seen in ATLAS [27–29, 35]

$$\mu_{\gamma\gamma}^{\text{ATLAS+CMS}} = \frac{\sigma(gg \rightarrow s \rightarrow \gamma\gamma)}{\sigma(gg \rightarrow h_{95.4}^{\text{SM}} \rightarrow \gamma\gamma)} = 0.24^{+0.09}_{-0.08}. \quad (3)$$

These observations have triggered speculations that they could arise from production of a new particle (see e.g. [36, 37] and references therein).

3.2 A light CP-even scalar in the MRSSM

In the MRSSM the CP-conserving Higgs boson is a mixture of four states: H_u , H_d , S and T . Apart from the known MSSM-like $SU(2)_L$ doublets H_u , H_d it features additionally an $SU(2)_L$ singlet (S) and a triplet (T). As the triplet should be rather heavy due to the electroweak precision observables ([8]) and SM-like Higgs state should be H_u dominated, it is enough to focus here on the 2×2 mass submatrix in $H_u - S$ space (we refer to [33] for notation and details)

$$\mathcal{M}_{u,S}^{\phi} = \begin{pmatrix} m_Z^2 + \Delta m_{rad}^2 & v_u \left(\sqrt{2}\lambda_u \mu_u^{\text{eff},-} + g_1 M_B^D \right) \\ v_u \left(\sqrt{2}\lambda_u \mu_u^{\text{eff},-} + g_1 M_B^D \right) & 4(M_B^D)^2 + m_S^2 + \frac{\lambda_u^2 v_u^2}{2} \end{pmatrix}, \quad (4)$$

A light, mostly signet scalar with small mixing to the SM-like Higgs requires

$$M_B^D, m_S \lesssim m_Z \quad (5)$$

The coupling of singlet to SM fermions is generated via aforementioned small mixing, controlled by the (2,1) entry of matrix in Eq. 4. The relative increase of $\mu_{\gamma\gamma}^{\text{ATLAS+CMS}}$ compared $\mu_{Zb\bar{b}}$ is achieved via the interplay between scaling of partial widths and total decay width of the singlet-like state with the mixing. Care has to be also taken to ensure that the mixing does not destabilize the SM-like properties of the H_u state. This is checked with `HiggsTools`.

In Tab. 1 we present two benchmark points consistent with both LEP and LHC excesses simultaneously. The points differ in the way how a proper dark matter relic density is achieved, on which we comment briefly at the end of this section. In Fig. 1 we show $2d$ slices of the parameter space around BMP7 allowed by SM-like Higgs measurements at 95% C.L. as reported by `HiggsSignals`. Similarly, Fig. 2 shows regions excluded by searches of non SM-like Higgses at 95% C.L. as reported by `HiggsBounds`.

Finally, as shown in Ref. [33], the points in Tab. 1 are also characterized by a Bino-Singlino dark matter with a mass of around 45 GeV. Both BMP7 and 8 give relic density of $\Omega h^2 = 0.121$ (either via the Z -boson resonance in the case of BMP7 or via the stau t -channel exchange in the case of BMP8) while remaining allowed by all dark matter direct detection experiments, including LUX-ZEPPELIN [38] (as checked by `micrOMEGAs v6` [39]). Constraints from direct production of SUSY particles at colliders were checked using `SModelS v2.3` [40].

	BMP7	BMP8
$\tan \beta$	49.5	49.8
B_μ	176^2	142^2
λ_d, λ_u	-0.193, -0.00658	0.161, -0.0135
Λ_d, Λ_u	1.49, -1.03	1.49, -0.722
M_B^D	45.2	42.1
m_S^2	27.4^2	54.1^2
$m_{R_u}^2, m_{R_d}^2$	$1292^2, 522^2$	$1033^2, 788^2$
μ_d, μ_u	1536, 658	1500, 1282
M_W^D	1458	1490
M_O^D		3000
m_T^2, m_O^2		$3000^2, 1500^2$
$m_{Q;1,2}^2, m_{Q;3}^2$	$3803^2, 3900^2$	$1465^2, 3477^2$
$m_{D;1,2}^2, m_{D;3}^2$	$3148^2, 3728^2$	$1456^2, 1990^2$
$m_{U;1,2}^2, m_{U;3}^2$	$1271^2, 2452^2$	$3285^2, 3967^2$
$m_{L;1,2}^2, m_{E;1,2}^2$	$1000^2, 1000^2$	$1680^2, 1022^2$
$m_{L;3,3}^2, m_{E;3,3}^2$	$1000^2, 1000^2$	$803^2, 185^2$
m_{H_d}	-1884^2	-1711^2
m_{H_u}	-1063^2	-1534^2
v_S	-3087	2004
v_T	0.35	0.0142
m_{h_1}	95.4	95.4
m_{h_2}	125.25	124.72
m_{W^\pm}	80.375	80.371
m_{χ_1}	44.98	42.65
$m_{\tilde{\tau}_R}$	1000	124.7
ρ_1^\pm	717	1310
m_a	24.85	54.20

Table 1: Benchmark points for the scenario discussed here: input parameters, parameters determined via tadpole equations and selected predicted, phenomenologically relevant, pole masses. Dimensionful parameters are given in GeV or GeV^2 , as appropriate. Input values are listed in the upper part of the table, while derived masses of some light physical states are in the lower part.

4. Summary and conclusions

Higgs physics remains one of the main goals of the LHC as well as a target for future e^+e^- colliders. Calculation of Higgs boson related observables and their comparison with experimental data is therefore instrumental in guiding our search for the theory of the Beyond the Standard Model physics. In this note we described the recent improvement to FlexibleSUSY spectrum generator which allows for an easy assessment of the validity of BSM Higgs sectors by interfacing

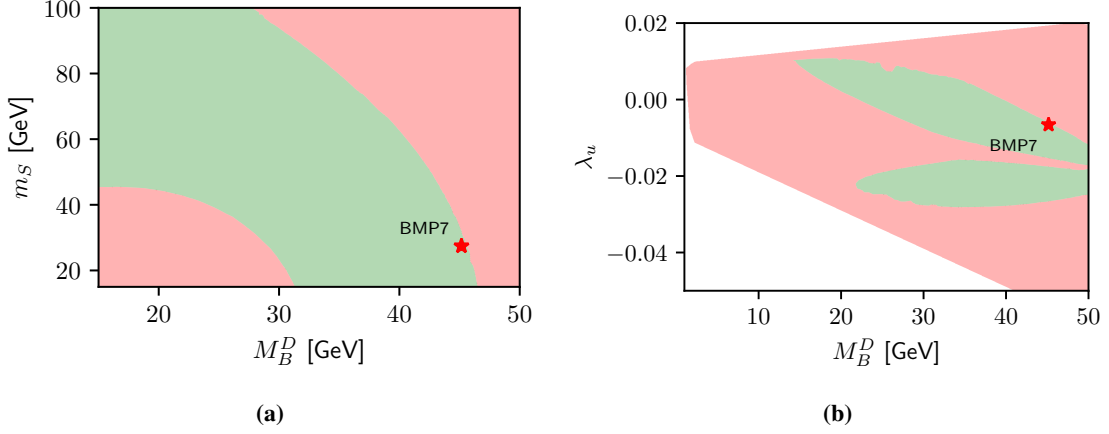


Figure 1: Parameter regions around BMP7 allowed (green) and excluded (red) by SM-like Higgs data at 95% C.L. as reported by HiggsSignals. White regions is where no spectrum could be generated.

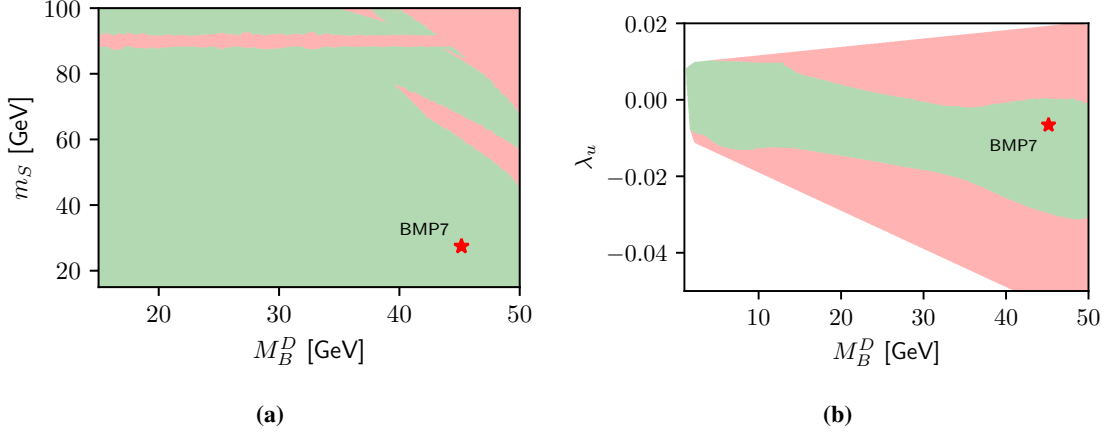


Figure 2: Parameter regions around BMP7 allowed (green) and excluded (red) by searches of non SM-like Higgses at 95% C.L. as reported by HiggsBounds. White regions is where no spectrum could be generated.

it with a popular program HiggsTools. We showcased it by analyzing excesses observed at LHC and LEP, interpreted as potential Higgs-like states at around 95 GeV in the context of the Minimal R-symmetric Supersymmetric Standard Model. Both the LEP excess in the $b\bar{b}$ channel and the LHC one in the $\gamma\gamma$ channel can be simultaneously accommodated, giving a SM-like 125 GeV Higgs boson in agreement with experimental measurements (as reported by HiggsTools). We provide two benchmark points illustrating these features. Those points also predict a correct dark matter relic density realised with a light bino-singlino particle (whose existence in the MRSSM is a necessary consequence of the existence of a 95 GeV Higgs-like state), which is also allowed by current direct detection experiments. Both benchmark points are allowed by current collider limits. While most states in BMP8 are fairly heavy (apart from light staus), BMP7 features an interesting prospect for future EW searches at the LHC thanks to the production and subsequent decay of $\rho^- \rightarrow W^- \chi_1$ (with a 100% branching ratio).

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