

A Data-driven Heavy-Metal Scenario for Ultra-High-Energy Cosmic Rays

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The mass composition of ultra-high-energy cosmic rays (UHECRs) is usually inferred from the depth of the shower maximum (X_{\max}) of cosmic-ray showers, which is only ambiguously determined by modern hadronic interaction models. We present a data-driven interpretation of UHECRs, the *heavy-metal scenario*, which assumes pure iron nuclei above $10^{19.6}$ eV (≈ 40 EeV) as the heaviest observed mass composition and introduces a global shift in the X_{\max} scale predicted by the two hadronic interaction models QGSJET II-04 and Sibyll 2.3d. We investigate the consequences of the proposed mass-composition model based on the obtained shifts in the X_{\max} values, which naturally lead to a heavier mass composition of UHECRs than conventionally assumed. We explore the consequences of our model on the energy evolution of relative fractions of primary species, consequently decomposed energy spectrum, hadronic-interaction studies and the arrival directions of UHECRs. We show that within this scenario, presented recently in [1], the cosmic-ray measurements can be interpreted in a more consistent way.

39th International Cosmic Ray Conference (ICRC2025)
15–24 July 2025
Geneva, Switzerland



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1. Introduction

Understanding the mass composition of ultra-high-energy cosmic rays (UHECRs) is an important step towards determining their sources. In addition, UHECRs provide a unique opportunity to study hadronic interactions at center-of-mass energies (100s TeV) far beyond the reach of current human-made accelerators (13 TeV). The mass composition of UHECRs is commonly estimated from the depth of the shower maximum (X_{\max}) and the number of muons produced during the development of extensive air showers that is detected at the ground level.

However, comparisons of these observables with model predictions reveal a tension. Recent results of the Pierre Auger Observatory show a discrepancy between the measured data and the models at more than 5σ [2]. This study suggests that a consistent description of the measured data requires a shift in the X_{\max} scale by $\approx 20 \text{ g/cm}^2 - 50 \text{ g/cm}^2$ together with rescaling of the hadronic component of the shower by $\approx 15\% - 25\%$ in the models EPOS-LHC [3], QGSJET II-04 [4] and SIBYLL 2.3d [5], which reduces the previous so-called muon puzzle approximately to its half. Moreover, machine-learning-based analyses of the Surface Detector data of the Pierre Auger Observatory [6] recently provided a precise estimation of the mean and fluctuation of X_{\max} , for the first time, for energies from 3 EeV up to 100 EeV. Interpreting these data with QGSJET II-04 leads to nonphysical negative values of the variance of logarithmic atomic mass $\ln A$, $\sigma^2(\ln A)$. For the other two models, SIBYLL 2.3d and EPOS-LHC, $\sigma^2(\ln A) \approx 0$ suggests a pure composition, while the first moment of $\ln A$ suggests a continuous increase of $\ln A$ with energy.

We present the *heavy-metal scenario*, very recently proposed in [1]. This scenario is built on two simple assumptions. First, we assume that the model predictions of X_{\max} can be shifted by an energy- and mass-independent value, while keeping the other predictions of the hadronic interaction models unchanged. Secondly, we assume pure iron nuclei above $10^{19.6} \text{ eV}$, where the fluctuations of X_{\max} are consistent with predictions for pure iron nuclei. We analyze publicly-available data from the Pierre Auger Observatory and investigate the implications of this scenario, which, from definition, leads to a heavier mass composition of UHECRs over a wide range of energies than conventionally interpreted. We explore how such an extreme mass-composition scenario affects the key aspects of UHECRs interpretations, including the inferred mass composition and spectral features, the muon puzzle, and arrival direction studies. For a more comprehensive analysis and discussions, we refer the reader to the full publication [1].

2. Adjustment of the X_{\max} scale

Assuming pure iron nuclei above $10^{19.6} \text{ eV}$ ($\approx 40 \text{ EeV}$) while keeping the elongation rate unchanged, the shift of the predicted X_{\max} scale, ΔX_{\max} , is fitted using the Auger DNN data [6]. The obtained shifts are $\Delta X_{\max} = 52 \pm 1_{-8}^{+11} \text{ g/cm}^2$ and $\Delta X_{\max} = 29 \pm 1_{-7}^{+12} \text{ g/cm}^2$ for QGSJET II-04 and SIBYLL 2.3d, respectively. Note that these values are consistent with the findings from [2] at 3 EeV-10 EeV. The mean and standard deviation of the X_{\max} from the Fluorescence Detector measurements [7] and from the Surface Detector measurements, the DNN method [6] and the Δ -method [8], are shown in the left panel of Fig. 1. The thin lines represent the predictions of the unmodified models of hadronic interactions and the thick lines the predictions corrected for ΔX_{\max} . Interpreting the X_{\max} moments using the moments of $\ln A$, the umbrella plot for the unmodified

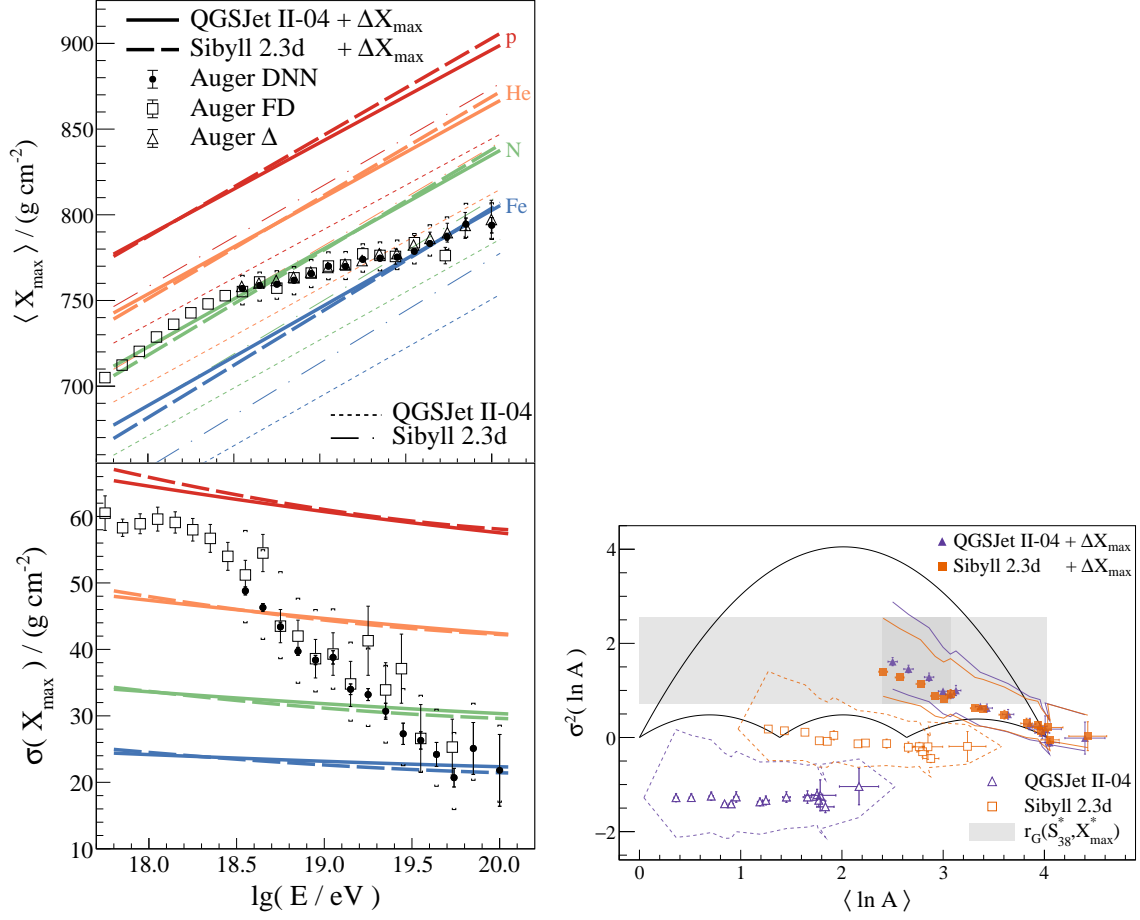


Figure 1: *Left panel:* The energy evolution of the mean and standard deviation of X_{\max} for data [7, 8, 10] measured by the Pierre Auger Observatory (black). Thin lines show the original model predictions for four primary species, thick lines represent the adjusted model predictions. *Right panel:* The interpreted mean and variance of $\ln A$ from the Auger DNN measurement for unmodified (open markers) and modified (full markers) model predictions. The region including the possible combinations of p, He, N and Fe nuclei is indicated by black curves. Taken from [1].

and modified model predictions is depicted in the right panel of Figure 1. While the original model predictions lie mostly outside of the allowed region for all possibilities when mixing four different primary species (protons, and helium, nitrogen and iron nuclei), the shifted model predictions are within the allowed region of the umbrella plot. Interestingly, also the value of $\sigma(\ln A)$ inferred from the correlation coefficient r_G between the ground signal and X_{\max} [9] is consistent with the shifted models.

3. Mass Composition and Energy Spectrum

To obtain the energy evolution of the relative abundance of the four particle species with the adjusted model predictions of X_{\max} for QGSJET II-04 and SIBYLL 2.3d, we fit the X_{\max} distributions from [11]. The fitted primary fractions for SIBYLL 2.3d are shown in the left panel of Figure 2 together with two parameterizations of the energy evolution of the primary fractions; a smoothed

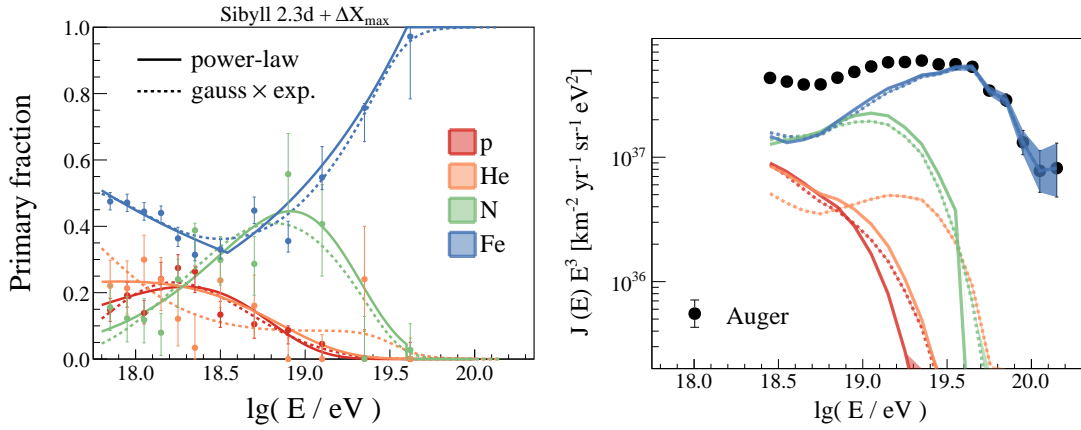


Figure 2: *Left panel:* Energy evolutions of the primary fractions for SIBYLL 2.3d + ΔX_{max} from fits to the X_{max} distributions [11]. The full and dotted lines represent two parameterizations of the fraction evolutions. *Right panel:* Differential fluxes of individual primary species as a function of energy for SIBYLL 2.3d + ΔX_{max} obtained from the all-particle spectrum (black markers) from [13]. Taken from [1].

Gaussian multiplied by an exponential function and power-law functions with a simple exponential cutoff (see [1] for details). The energy evolution of the primary fractions shows a transition from lighter to heavier nuclei with energy, similarly to [12] but with a smaller contribution of light elements in the mix.

The energy spectrum of individual primaries is depicted in the right panel of Figure 2 for SIBYLL 2.3d using the all-particle spectrum from [13]. The instep feature, around ≈ 15 EeV, is then connected to the transition between nitrogen and iron nuclei within the heavy-metal scenario. The flux suppression is by assumption dominated by iron nuclei. The rigidity cutoff for both nitrogen and iron nuclei coincides at $\approx 10^{18.2}$ V, suggesting a common origin of these primaries.

4. Hadronic Interactions

The mass composition of cosmic rays is closely related to the number of muons in the extensive air showers. With the mass composition obtained within the heavy-metal scenario, we show the reduction of the so-called muon puzzle to approximately half in the left panel of Fig. 3 in the direct muon measurements at the Pierre Auger Observatory when the ΔX_{max} is applied to the model predictions. This result is consistent with the finding in [2].

The description of measured X_{max} distributions from [11] is fair enough even for the QGSJET II-04 model, when ΔX_{max} is applied. We show in the right panel of Fig. 3 such an example for $10^{18.0-18.5}$ eV in the case of SIBYLL 2.3d + ΔX_{max} . Estimating the slope of the X_{max} distribution as in [14] ($\Lambda_\eta = (55.8 \pm 2.3 \pm 1.6) \text{ g/cm}^2$), we obtain $\Lambda_\eta = (51.9 \pm 0.4) \text{ g/cm}^2$ and $\Lambda_\eta = (50.0 \pm 0.4) \text{ g/cm}^2$ in case of QGSJET II-04 + ΔX_{max} and SIBYLL 2.3d + ΔX_{max} , respectively. These somewhat lower values in case of the heavy-metal scenario might be a consequence of a higher helium contribution or a higher p-p cross-section or a lower elasticity that are extrapolated in the models from the accelerator measurements.

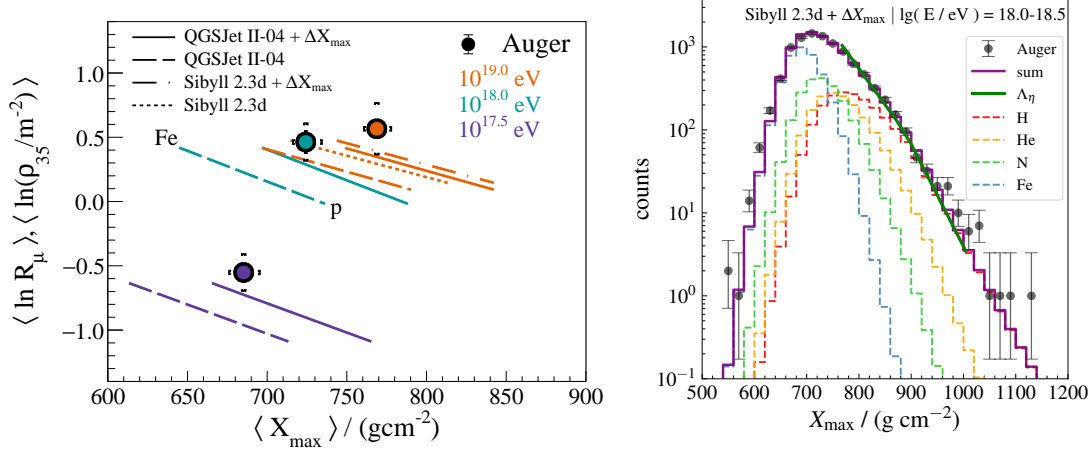


Figure 3: *Left panel:* The muon number obtained by direct measurements at Auger [15, 16] compared to the predictions with and without the application of ΔX_{\max} . *Right panel:* The X_{\max} distribution of Auger data from [11] compared to the prediction for the heavy-metal scenario using SIBYLL 2.3d + ΔX_{\max} in the energy bin $10^{18.0-18.5}$ eV. We illustrate the Λ_η fit by the green line. Taken from [1].

5. Arrival directions

We also investigate the implications of the heavy-metal scenario on the arrival directions of UHECRs, taking into account their deflections in the Galactic magnetic field (GMF). We use the eight UF23 models of the GMF [17] together with the turbulent component of the Planck-tuned JF12 model [18, 19].

We first study the possible features of the extragalactic dipole that would result in the observed one on the Earth above 8 EeV [20]. Following the same procedure as in [21], we use the evolution of the mass composition of cosmic rays above 8 EeV obtained for SIBYLL 2.3d + ΔX_{\max} . By simulating an ideal dipole flux outside the Galaxy with varying directions and amplitudes, we find regions of possible extragalactic dipole directions compatible with the Pierre Auger Observatory measurements [20] at 1σ and 2σ levels. The 1σ and 2σ regions of the possible extragalactic directions of the dipole are visualized in the left panel of Figure 4. Due to the heavy mass composition, the extragalactic dipole amplitude must be relatively large ($\geq 12\%$) to match the observed amplitude on the Earth.

Secondly, we backtrack the 100 most energetic Auger events above 78 EeV [22], assuming pure iron nuclei. The simulated directions of the backtracked particles at the edge of the Galaxy are shown in the right panel of Figure 4. Due to the heavy mass composition, these particles experience strong deflections in the GMF, and the distribution of the backtracked cosmic rays at the edge of the Galaxy is the most dense in the direction close to the Galactic anti-center.

6. Summary

We present a data-driven mass composition scenario of ultra-high-energy cosmic rays, assuming a pure iron nuclei above the energy $10^{19.6}$ eV (approximately the region of steep flux suppression) and global shifts in the X_{\max} scale of two hadronic interaction models QGSJET II-04 and SIBYLL 2.3d resulting in a consistent mass interpretation of the mean and variance of $\ln A$.

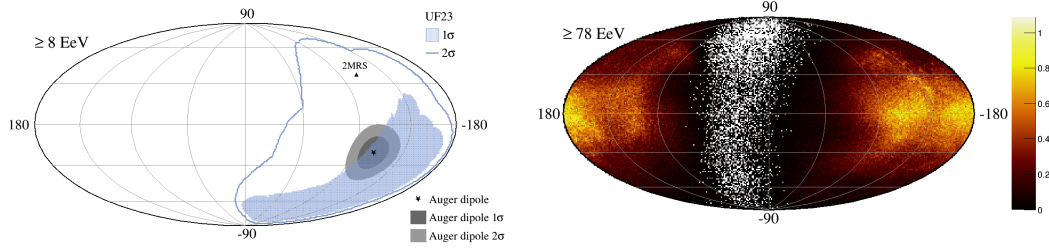


Figure 4: *Left panel:* The 1σ and 2σ regions of allowed extragalactic dipole directions compatible with the measurement of the Pierre Auger Observatory [20] after propagation in the GMF using UF23 models. *Right panel:* The distribution of the backtracked directions at the edge of the Galaxy of the 100 most energetic Auger events above 78 EeV [22] using UF23 models of the GMF. Taken from [1].

The fitted X_{\max} shifts are consistent with the results from [2] obtained from the two-dimensional fits of X_{\max} and the ground signal distributions at energies 3 EeV - 10 EeV, including the heavier mass composition than obtained for unmodified model predictions. Additionally, we show the consequences of the proposed heavy-metal scenario on the consistency of the decomposed energy spectrum, hadronic-interaction studies, and backtracked arrival directions in the Galactic magnetic field (GMF).

Within the heavy-metal scenario, the decomposed energy fluxes of four primaries (protons and He, N and Fe nuclei) show the flux suppression of iron and nitrogen nuclei at rigidity $\approx 10^{18.2}$ V and explain the instep feature in the energy spectrum as a result of the fading of the nitrogen component. The agreement of the rigidity cutoff for nitrogen and iron nuclei also suggests their common origin. With the adjusted X_{\max} scale and, therefore, heavier mass composition over the whole energy range, the muon puzzle is reduced by approximately half. The measured X_{\max} distributions can be well described using the new mass-composition model. However, we obtain slightly lower values of the slope of the X_{\max} distribution, that might be caused by a higher fraction of helium nuclei in our scenario or the need for a lower extrapolated p-p cross section or a higher elasticity in the modeled hadronic interactions. The observed dipole anisotropy above 8 EeV is compatible with an extragalactic origin under our mass-composition scenario, requiring extragalactic dipole amplitudes $\geq 12\%$. Assuming only iron nuclei, the arrival directions of the most energetic Auger events above 78 EeV, when backtracked through the GMF, point towards the Galactic anti-center region, which is consistent with the expectations from isotropic arrival directions at the Earth. The estimated source luminosity points to only the most powerful astrophysical objects, such as hard X-ray AGNs, as viable candidates.

In the next steps, we plan to explore the possibility of including even heavier nuclei than iron at the highest energies, as suggested by recent studies, see e.g. [23, 24]. Moreover, we intend to test the proposed scenario using next-generation hadronic interaction models that have been recently developed and made publicly available.

Acknowledgments

The work was supported by the Czech Academy of Sciences: LQ100102401, Czech Science Foundation: 21-02226M, Ministry of Education, Youth and Sports, Czech Republic: FORTE CZ.02.01.01/00/22_008/0004632, German Academic Exchange service (DAAD PRIME). The authors are very grateful to the Pierre Auger Collaboration for discussions about this work.

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