

An Update on the Isospin-Breaking Effects in the Pion Decay Constant with Staggered Quarks

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We present an update on the ongoing computation of the isospin-breaking effects in the Pion Decay Constant from the BMW Collaboration. The calculation is carried out with $N_f=2+1+1$ staggered quarks with a near-physical pion mass and QED_L . We give an update on the isosymmetric value and the current determination used to compute the gradient-flow scale w_0 , then we present some preliminary results of the valence-valence contribution to the axial-pseudoscalar correlator for different volumes and lattice spacings. We also discuss the next steps and plans.

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1. Motivations

The unitarity of the first row of the CKM matrix is an open problem in precision flavour physics, as shown in Figure 1 from [1]. At a current level of precision, isospin-breaking effects become relevant.

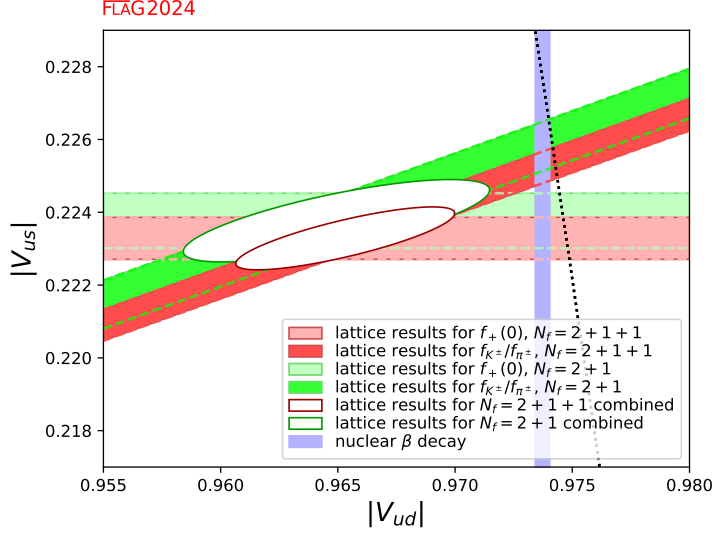


Figure 1: Status of the determination of the V_{ud} and V_{us} elements of the CKM matrix as reported by the Flavour Lattice Averaging Group [1].

In Isosymmetric QCD the pion decay constant f_π is defined by the hadronic matrix element:

$$\langle 0 | \bar{u} \gamma_4 \gamma_5 d | \pi(\vec{0}) \rangle = M_\pi f_\pi. \quad (1)$$

To compute the pion decay constant in QCD+QED F_π one needs to consider the leptonic decay rate $\pi^+ \rightarrow l^+ \nu_l(\gamma)$, in the PDG parametrization it is:

$$F_\pi^2 = \frac{\Gamma_{\pi^+ \rightarrow l^+ \nu_l(\gamma)}}{\frac{G_F^2}{8\pi} |V_{ud}|^2 M_\pi m_l^2 \left(1 - \frac{m_l^2}{M_\pi^2}\right)^2} = f_\pi^2 [1 + \delta R_\pi] \quad (2)$$

with δR_π the isospin-breaking correction to the decay constant. The full computation of δR_π requires the renormalization of the electroweak hamiltonian:

$$\mathcal{H}_W = \frac{G_F}{\sqrt{2}} V_{ud}^* [\bar{d} \gamma^\mu (1 - \gamma_5) u] [\bar{\nu}_l \gamma_\mu (1 - \gamma_5) l]. \quad (3)$$

There is one lattice computation of the isospin-breaking corrections in the pion and kaon decay constants by the RM123-Southampton collaboration [2]. In the ratio f_K/f_π the renormalization constant cancel (in a massless scheme) so there are more lattice computations of the correction to the latter, from the RM123-Southampton collaboration [3], from RBC-UKQCD [4] and a more recent from RBC-UKQCD with QED $_\infty$ [5].

The pion decay constant F_π is used for the scale setting in the BMW/DMZ determination of the anomalous magnetic moment of the muon [6].

2. Mixed Determination

The isospin-breaking effects in the pion decay constant are computed by combining the work from the RM123-Southampton collaboration [2] and the BMW results for isosymmetric QCD and the sea quark IBE.

The dependency of the pion decay constant in units of w_0 on the quark electric charge and bare quark masses is parametrized as follows (up to higher order effects):

$$w_0 F_\pi = A + B F_\pi^{-2} M_{\pi^+}^2 + C F_\pi^{-2} \left(M_{K^\pm}^2 + M_{K^0}^2 - M_{\pi^\pm}^2 \right) / 2 + E e_v^2 + F e_v e_s + G e_s^2 \quad (4)$$

where e_v and e_s are, respectively, the valence and sea quark charges. The coefficients B and C are mass derivatives and are computed in QCD. The coefficients E , F and G are the valence-valence, sea-valence and sea-sea derivatives.

Once all the coefficients in Eq. (4) are determined, the parametrization can be split into the QCD+seaQED and the valQED part:

$$[w_0 F_\pi]_{\text{QCD+QED}} = [w_0 F_\pi]_{\text{QCD+seaQED}} + [w_0 F_\pi]_{\text{valQED}}, \quad (5)$$

both of the terms on the right-hand side are dependent on the scheme used to define QCD. It is fundamental to remember that w_0 has not valence quark isospin-breaking effects so $[w_0 F_\pi]_{\text{valQED}}$ includes only the valence quark IBE on F_π . The quantity $[w_0 F_\pi]_{\text{QCD+seaQED}}$ is determined with the BMW data, but the valence contribution is computed from the RM123-Southampton result¹.

To combine the two results, the same scheme has to be employed. The work [2] is defined in the GRS scheme by the quantities [7]:

$$\begin{aligned} [M_\pi]^{\text{GRS}} &= 135.0(2) \text{ MeV} \\ \left[\frac{1}{2} (M_{K^\pm} + M_{K^0}) \right]^{\text{GRS}} &= 494.6(1) \text{ MeV} \\ [f_\pi]^{\text{GRS}} &= 130.65(12) \text{ MeV} \\ [\delta R_\pi]^{\text{GRS}} &= 0.0153(19)^2. \end{aligned} \quad (6)$$

Since they were working in the electroquenched approximation, which neglects the sea quark IBE, the QCD result includes the sea quark isospin-breaking effects following:

$$[*]^{\text{GRS}} = [*]_{\text{QCD+QED}}^{\text{phys}} - [*]_{\text{valQED}}^{\text{GRS}} = [*]_{\text{QCD+seaQED}}^{\text{GRS}}. \quad (7)$$

¹The valence QED contribution to the pion decay rate from Ref. [2] is given in a GRS scheme in which the quark masses, renormalized at $\mu = 2 \text{ GeV}$ in $\overline{\text{MS}}$, remain constant as the electromagnetic coupling is turned on or off and the overall lattice scale is fixed using f_π . In the electroquenched approximation adopted there valence QED corrections do not affect the renormalization of the strong coupling – *i.e.* the GRS condition $[\hat{g}_s]_{\text{QCD}}(\mu) = [\hat{g}_s]_{\text{QCD+valQED}}(\mu)$ (*), where $\hat{}$ indicates the quantities renormalized in the respective theories, holds trivially – and thus do not change the lattice spacing from its pure QCD value at fixed β . We stress that in the full QCD+QED calculation presented here the lattice scale is consistently set using F_π and, since the purely gluonic quantity w_0 is free of valence QED contributions – *viz.* a relation similar to (*) occurs for w_0 as well –, it is justified to use the value of δR_π from Ref. [2] as input in our computation to evaluate the term $[w_0 F_\pi]_{\text{QCD+valQED}} / [w_0 F_\pi]_{\text{QCD}}$ and Eq. (8) without introducing scheme ambiguities.

²The value used for our determination of $w_0 F_\pi$ is $\delta R_\pi = 0.0150(18)$, it does not include the original χ PT estimate of the sea-quark IBE.

Using the parametrization (4), the identity (7) can be used to evaluate $[w_0 F_\pi]_{\text{QCD+seaQED}}^{\text{GRS}}$ which can then be combined with the RM123S results using the relation, up to negligible $O(\alpha_{\text{em}}^2)$ corrections:

$$[w_0 F_\pi]_{\text{QCD+QED}} = [w_0 F_\pi]_{\text{QCD+seaQED}}^{\text{GRS}} \sqrt{1 + [\delta R_\pi]^{\text{GRS}}}. \quad (8)$$

3. Isosymmetric QCD value in the FLAG scheme

The isosymmetric QCD value of $w_0 f_\pi$ is computed using the ensembles of the BMW collaboration shown in Figure 2 with QCD tree-level Symanzik action for the gauge field and $N_f=2+1+1$ Staggered fermions coupled with 4 levels of stout smearing with radius $\rho = 0.125$. The decay constant is computed from the pseudoscalar two-point function through the chiral WTI.

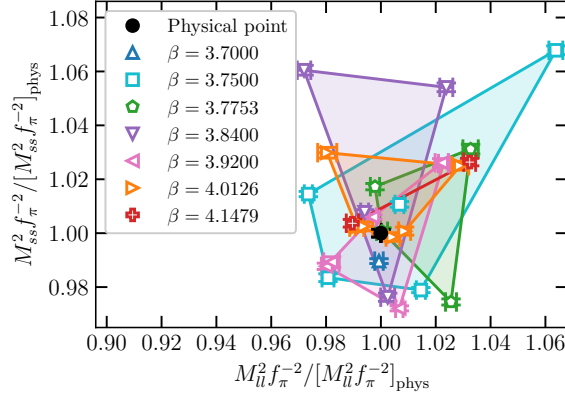


Figure 2: Landscape of the ensembles used for the isoQCD determination of $w_0 f_\pi$ around the physical point from [6].

The analysis follows the strategy presented in [8]. To have a better control over the continuum extrapolation we add the determination of w_0 using the Zeuthen-flow [9] in addition to the Wilson-flow [10], Figure 3 shows the continuum extrapolation of $w_0 f_\pi$ in the lattice spacing and in the taste-breaking term and the relative probability distribution of both the statistical and systematic effect. The result is presented in the FLAG scheme [11].

Table 1 shows the breakdown of the total error on $w_0 f_\pi$. The uncertainty is dominated by the systematic error and, in particular, the biggest contribution is represented by the continuum extrapolation under the entries: *Type of gradient flow*, *Lattice spacing cuts* and *Order of fit polynomials*. To improve the determination, we are generating data at a finer value of the lattice spacing. Other sources of systematic error listed are: *Pseudoscalar fits* (effects of the fit range of the effective pion decay constant) and *Finite volume χ PT order* (the difference in the finite volume effects between NLO and NNLO χ PT).

The isosymmetric part of $w_0 f_\pi$ is determined with a 0.23% precision, thus making it necessary to include IBE.

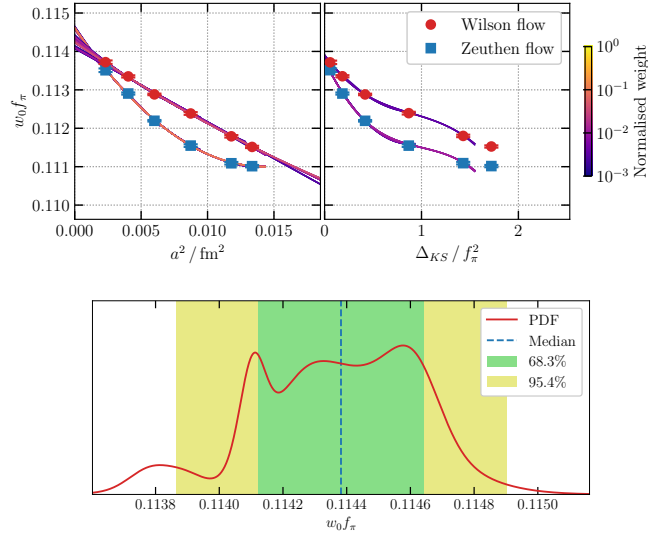


Figure 3: Continuum limit extrapolation of the iso QCD value of $w_0 f_\pi$. The upper plot represents the continuum extrapolation both in the lattice spacing and in the taste-symmetry breaking terms Δ_{KS} (computed by comparing the masses of different meson tastes on our ensembles); in the lower plot the probability distribution function (PDF) including both statistical and systematic variations is shown in red.

Median	0.11438	
Total error	0.00026	0.22736 %
Statistical error	0.00012	0.10786 %
Systematic error	0.00023	0.20015 %
Type of gradient flow	0.00013	0.10966 %
Pseudoscalar fits	0.00004	0.03144 %
Finite volume χ PT order	0.00001	0.01023 %
Lattice spacing cuts	0.00006	0.04956 %
Order of fit polynomials	0.00019	0.16215 %

Table 1: Error budget of the isoQCD value of $w_0 f_\pi$. The first row contains the median value, from the second on the errors are presented with the value and the percentage relative value. The last five rows present a breakdown of the systematic error on $w_0 f_\pi$.

4. Sea Quark Isospin-Breaking Effects

As mentioned in Section 2, the sea quark isospin-breaking effects to $w_0 F_\pi$ are determined using the BMW ensemble. The sea-sea diagrams contributing to the electromagnetic derivative of the pion decay constant are:

$$(9)$$

The renormalized sea quark electromagnetic derivative is computed by cancelling the UV-divergence with the one coming from the scheme observables in the parametrization (4) following the convention from [12].

5. Total Error Budget

Combining the results for the isoQCD and the IBE as mentioned in Section 2 we obtain the following preliminary QCD+QED value of $w_0 F_\pi$:

$$[w_0 F_\pi]_{\text{QCD+QED}} = 0.11527(15)(26)[31] \text{ (Preliminary)}, \quad (11)$$

where the first error is the statistical, the second is systematic, and the latter is the combined total error. The systematic error consists of the following elements:

$$(26) = (23)_{\text{QCD}}(11)_{\text{QED-Valence}}(5)_{\text{QED-Sea}} \text{ (Preliminary)}. \quad (12)$$

The QCD error dominates the final uncertainty, while the QED error is dominated by the valence QED value.

Combining $w_0 F_\pi$ with the QCD+QED value of $F_\pi = 131.711(45)$ MeV (obtained using $V_{ud} = 0.97367(32)$ and $\Gamma = 3.8408(7) \times 10^7 \text{ s}^{-1}$ from [14]) we get the following value of w_0 :

$$[w_0]_{\text{QCD+QED}} = 0.17270(22)(40)[46] \text{ fm (Preliminary)}. \quad (13)$$

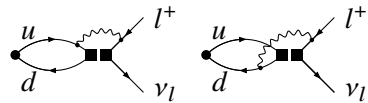
The systematic error is made by the following contributions:

$$(40) = (34)_{\text{QCD}}(17)_{\text{QED-Valence}}(8)_{\text{QED-Sea}}(7)_{\text{exp}} \text{ (Preliminary)}, \quad (14)$$

where the additional systematic error *exp* accounts for the experimental value of F_π . We are working on improving the QCD error and improving the valence QED determination by providing an independent determination of δR_π and by checking the power-law finite volume effects in light of the work from [4].

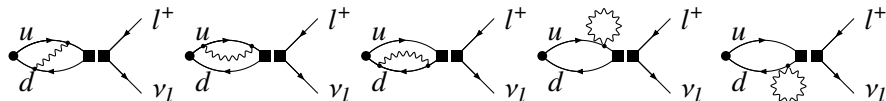
6. Ongoing Work on the Valence Quark Isospin-Breaking Effects

The valence-quark IBE in the pion decay constant requires the computation of the non-factorizable diagrams:



$$(15)$$

(actively involving the lepton) and the factorizable ones; we present some preliminary results of the factorizable diagrams. The factorizable valence contributions to the electromagnetic derivative of the pion decay constant are given by the following diagrams:



$$(16)$$

The diagrams are computed using sequential propagators and 16 photon sources in Coulomb gauge and 3 U(1) random sources. The zeroth mode of the photon is removed following the QED_L prescription. We use the ensembles listed in Table 3. We tested the program by comparing it with numerical derivatives at a fixed random source and photon source.

The correction is obtained by expanding the axial-pseudoscalar (using the conserved axial current) and pseudoscalar-pseudoscalar two-point function:

$$\frac{G_{ud}^{PA_4}(x_4)^2}{G_{ud}^{PP}(x_4)} = M_\pi Z_A^{-2} F_\pi^2 e^{-M_\pi T/2} \frac{\sinh^2 [M_\pi (T/2 - x_4 - 1/2)]}{\cosh [M_\pi (T/2 - x_4)]} \quad (17)$$

where T is the temporal extension and Z_A is the renormalization constant which is not considered in these results and, thus, it is left for future work.

action	β	a [fm]	$L/a \times T/a$	tag	am_s	m_s/m_l	#confs
4stout	3.7000	0.1315	24×48	volume/24	0.057291	27.899	48
			32×64	volume/32	0.057291	27.899	48
			48×64	volume/48	0.057291	27.899	48
			64×64	volume/64	0.057291	27.899	48
	3.8400	0.0952	64×96	dir00	0.043194	28.500	48

Table 3: Set of ensembles used for the computation of the valence quarks isospin-breaking effects.

Figure 5 shows the preliminary result for the effective factorizable valence derivative for the four different volumes considered.

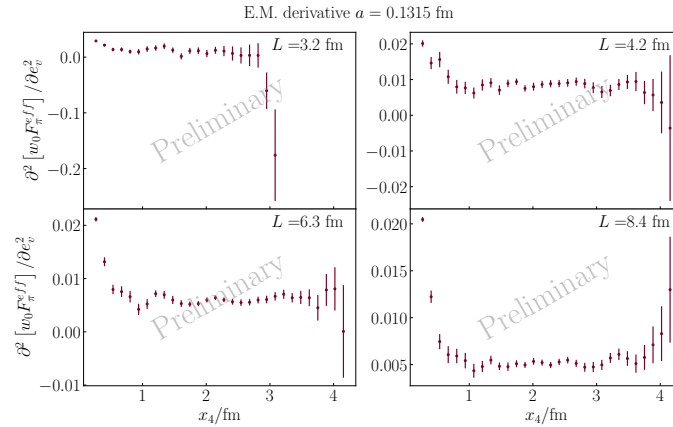


Figure 5: Factorizable valence electromagnetic derivative of $w_0 F_\pi$ for $a = 0.1315$ fm. The preliminary results for the four volumes are shown.

Figure 6 shows the preliminary result for the effective factorizable valence derivative for the finer value of the lattice spacing. The effective derivative presents a clear plateau at large time separation, showing little excited states contamination and little oscillating contribution.

7. Conclusions & Outlook

In this Proceedings, we have presented the results of the current determination of the pion decay constant in QCD+QED by the BMW collaboration for the determination of the scale setting quantity w_0 , the preliminary value is:

$$[w_0]_{\text{QCD+QED}} = 0.17270(22)(40)[46] \text{ fm (Preliminary)}, \quad (18)$$

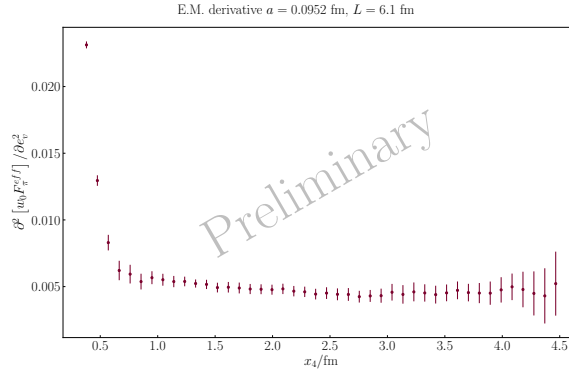


Figure 6: Factorizable valence electromagnetic derivative of $w_0 F_\pi$ for $a = 0.0952$ fm.

where the precision is dominated by the continuum extrapolation of the isosymmetric part. The determination of isospin-breaking effects was done by combining sea quark effects from BMW ensembles and the valence-quark effects from the RM123 work.

We are working on improving the continuum extrapolation (by adding a finer lattice spacing) and on a completely independent determination of the valence effects, which will include the sea-valence (that have been estimated to be suppressed). This requires the determination of the non-factorizable diagrams involving the leptons and the renormalization of the electroweak Hamiltonian. The authors of [4, 15] showed the poor convergence of the series of finite volume effects to the pion decay constant in QED_L, for this reason, we plan a detailed study of the finite volume effects using different volumes (up to 10.8 fm).

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References

- [1] (FLAVOUR LATTICE AVERAGING GROUP (FLAG)), Y. Aoki et al., *FLAG Review 2024*, [2411.04268](https://arxiv.org/abs/2411.04268).

- [2] M. Di Carlo, D. Giusti, V. Lubicz, G. Martinelli, C. T. Sachrajda, F. Sanfilippo et al., *Light-meson leptonic decay rates in lattice QCD+QED*, *Phys. Rev. D* **100** (2019) 034514, [[1904.08731](#)].
- [3] D. Giusti, V. Lubicz, G. Martinelli, C. T. Sachrajda, F. Sanfilippo, S. Simula et al., *First lattice calculation of the QED corrections to leptonic decay rates*, *Phys. Rev. Lett.* **120** (2018) 072001, [[1711.06537](#)].
- [4] M. Di Carlo, M. T. Hansen, A. Portelli and N. Hermansson-Truedsson, *Relativistic, model-independent determination of electromagnetic finite-size effects beyond the pointlike approximation*, *Phys. Rev. D* **105** (2022) 074509, [[2109.05002](#)].
- [5] N. H. Christ, X. Feng, T. Izubuchi, L. Jin, C. T. Sachrajda and X.-Y. Tuo, *Lattice Calculation of Light Meson Radiative Leptonic Decays*, [2510.26993](#).
- [6] (BMW, DMZ), F. M. Stokes et al., *BMW/DMZ calculation of the hadronic vacuum polarisation for the muon magnetic moment*, 3, 2026. [2603.03835](#).
- [7] D. Giusti, V. Lubicz, C. Tarantino, G. Martinelli, F. Sanfilippo, S. Simula et al., *Leading isospin-breaking corrections to pion, kaon and charmed-meson masses with Twisted-Mass fermions*, *Phys. Rev. D* **95** (2017) 114504, [[1704.06561](#)].
- [8] A. Boccaletti et al., *High precision calculation of the hadronic vacuum polarisation contribution to the muon anomaly*, [2407.10913](#).
- [9] A. Ramos and S. Sint, *Symanzik improvement of the gradient flow in lattice gauge theories*, *Eur. Phys. J. C* **76** (2016) 15, [[1508.05552](#)].
- [10] M. Lüscher, *Properties and uses of the Wilson flow in lattice QCD*, *JHEP* **08** (2010) 071, [[1006.4518](#)]. [Erratum: *JHEP* 03, 092 (2014)].
- [11] (FLAVOUR LATTICE AVERAGING GROUP (FLAG)), Y. Aoki et al., *FLAG Review 2021*, *Eur. Phys. J. C* **82** (2022) 869, [[2111.09849](#)].
- [12] S. Borsanyi et al., *Leading hadronic contribution to the muon magnetic moment from lattice QCD*, *Nature* **593** (2021) 51–55, [[2002.12347](#)].
- [13] V. Lubicz, G. Martinelli, C. T. Sachrajda, F. Sanfilippo, S. Simula and N. Tantalo, *Finite-Volume QED Corrections to Decay Amplitudes in Lattice QCD*, *Phys. Rev. D* **95** (2017) 034504, [[1611.08497](#)].
- [14] (PARTICLE DATA GROUP), R. L. Workman et al., *Review of Particle Physics*, *PTEP* **2022** (2022) 083C01.
- [15] P. Boyle et al., *Isospin-breaking corrections to light-meson leptonic decays from lattice simulations at physical quark masses*, *JHEP* **02** (2023) 242, [[2211.12865](#)].