

Probing TeV scale physics via ultra cold neutron decays and calculating non-standard baryon matrix elements

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We motivate undertaking precision analyses of neutron decays to look for signatures of new scalar and tensor interactions that can arise in extensions of the Standard Model at the TeV scale. The key ingredient needed to connect experimental data with theoretical analysis are high-precision calculations of matrix elements of isovector bilinear operators between the decaying neutron and final state proton. We describe the status of our Lattice QCD program of using valence clover fermions on dynamical $N_f = 2 + 1 + 1$ HISQ configurations generated by the MILC Collaboration. On the theoretical side we use the effective field theory method and provide both model independent and dependent analyses to obtain bounds on possible scalar and tensor interactions, both from low energy experiments and LHC data.

*The XXIX International Symposium on Lattice Field Theory
July 10-16, 2011
Lake Tahoe, California*

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†LA-UR-12-10238

1. Introduction

New physics at a given high energy scale, such as that generated by the exchange of any heavy boson, gives rise to measurable effects at the hadronic scale. The most well known example is weak interactions (WI) mediated by the exchange of W and Z vector bosons that are responsible, for example, for neutron β -decay. In the 1930s, Fermi parameterized the charged current decay in terms of an effective four-fermion interaction. This effective theory was also able to explain muon decay in terms of a single coupling constant G_F . Elucidation of the Lorentz structure of the interaction, $V - A$, came in 1957 with the work of Marshak and Sudershan, and Feynman and Gell-Mann. Until then, physicists had considered scalar and tensor interactions as viable candidates. Analysis of low energy decays revealed a number of properties of WI, such as quark-lepton universality, and the weak scale, ~ 100 GeV, related to the mass of the W and Z vector bosons that are now part of the electroweak theory. In formulating a consistent electroweak theory, it was also clear that associated with the weak scale there has to be a second scale, $O(\text{TeV})$, characterizing the breaking of electro-weak symmetry since at the hadronic scale WI are 10^{-5} times smaller than the electric force.

To break the electro-weak symmetry, explain the origin of masses of quarks, leptons and weak bosons (Higgs phenomena) and to stabilize the Standard Model against radiative corrections motivates the need for new physics at the TeV scale. The LHC has, in fact, been designed to probe new physics at this scale. To explore the $1 - 10$ TeV range, and re-examine the possibility of scalar and tensor interactions, we follow the complementary approach to direct detection at the LHC by looking for deviations from the SM in precision experiments at the hadronic scale. This talk gives an overview of the combined experimental and theoretical precision analyses of ultracold neutron (UCN) decays required to probe novel interactions at the TeV Scale. In particular, we would like to address the following questions:

- Are there novel scalar or tensor interactions at the TeV scale?
- What are the respective coupling constants and what bounds can be put on them?
- What are their manifestations at the hadronic scale and what decays/processes are the best candidates for investigating them or for setting upper bounds on their strengths?
- What beyond the SM candidate theories can accomodate these new interactions and satisfy all other known constraints?

Decays of UCN are interesting for probing new physics because there are terms in the neutron decay distribution that are particularly sensitive to scalar and tensor interactions at the 10^{-3} level. In these terms in the decay distribution, the SM contributions are suppressed due to the helicity flip factor m_e/E_e , so they are of the same order as the signal. The details of the theoretical framework, analysis and preliminary results are given in Ref. [1], and we will follow the notation established there. In this talk we summarize the key points of the formalism and give some preliminary results. A discussion of the Lattice QCD part of the project will be presented by Lin [2] and of the existing bounds from other low energy experiments and from the LHC by Bhattacharya [3].

2. Neutron decay

The differential neutron decay distribution function $D(E_e, \mathbf{p}_e, \mathbf{p}_\nu, \boldsymbol{\sigma}_n)$, keeping only terms relevant to the discussion here, can be written as [4, 1]

$$D(E_e, \mathbf{p}_e, \mathbf{p}_\nu, \boldsymbol{\sigma}_n) = 1 + c_0 + c_1 \frac{E_e}{M_N} + \frac{m_e}{E_e} \bar{b} + \bar{B}(E_e) \frac{\boldsymbol{\sigma}_n \cdot \mathbf{p}_\nu}{E_\nu} + \dots, \quad (2.1)$$

where \mathbf{p}_e and \mathbf{p}_ν denote the electron and neutrino three-momenta, and $\boldsymbol{\sigma}_n$ denotes the neutron polarization. Here $\bar{b} \equiv b^{\text{SM}} + b^{\text{BSM}}$ is an effective Fierz interference term and $\bar{B}(E_e)$ is an effective energy-dependent correlation coefficient given by

$$\bar{B}(E_e) = B_{\text{LO}}(\tilde{\lambda}) + c_0^{(B)} + c_1^{(B)} \frac{E_e}{M_N} + \frac{m_e}{E_e} (b_\nu^{\text{SM}} + b_\nu^{\text{BSM}}) \quad (2.2)$$

The terms most sensitive to novel scalar and tensor interactions are proportional to the helicity-flip factor $\frac{m_e}{E_e}$, and characterized by the parameters \bar{b} and $b_\nu^{\text{SM}} + b_\nu^{\text{BSM}}$. The beyond the Standard Model (BSM) contributions to these parameters arise from novel scalar and tensor interactions (with coupling constants ε_S and ε_T) generated at the TeV scale. To extract \bar{b} requires precision measurements of the electron spectrum, while $(b_\nu^{\text{SM}} + b_\nu^{\text{BSM}})$ requires measuring the asymmetry between the neutron polarization and the anti-neutrino momentum.

Current bounds on b^{BSM} and b_ν^{BSM} leave open the possibility that these could be as large as 10^{-3} , which implies that the scale for these novel interactions is TeV. Standard Model contributions, b^{SM} and b_ν^{SM} , are $O(10^{-3})$ and known to $O(10^{-5})$. So the sought after signal is the same size as the well-understood SM background. Two experiments at the Los Alamos ultracold neutron source, called UCNb [5] and UCNB [6], are being carried out by our experimental colleagues to measure these parameters at the 10^{-3} level. In a UCN decay the recoil effects, $(p_n - p_p)/M_N \sim 10^{-3}$, are small and of the same size as the signal while there are no nuclear effects. Even though the signal in UCN decays is not swamped by the background, nevertheless, the experiments are very challenging.

3. UCNb and UCNB experiments

The two separate experiments, UCNb and UCNB, will use the same UCN source. The setup is shown in Fig. 1. The proposed plan is to make 10^{-3} measurements of \bar{b} and $\bar{b} - (b_\nu^{\text{SM}} + b_\nu^{\text{BSM}})$ by the end of 2013 and based on the results and detector performance and sensitivity achieved propose a 10^{-4} level experiment to DOE.

The goal of the UCNb experiment is a high precision measurement of the electron decay spectrum. In Fig. 2 we show the full SM expected spectrum $F_0(E_e)$ of the electron in the limit of zero nuclear recoil (left panel) and the deviation from it assuming $b^{\text{BSM}} = 10^{-3}$ (right panel). To achieve the design sensitivity requires building a very sensitive and high resolution calorimeter. The UCNB experiment measures the asymmetry between the momentum direction of the neutrino and the spin of the decaying neutron and yields $\bar{b} - (b_\nu^{\text{SM}} + b_\nu^{\text{BSM}})$. The neutrino momentum is determined indirectly by measuring the electron and proton momenta in coincidence. (In practice, the precision possible on these momenta is not sufficient to fully reconstruct the neutrino momentum so a reduced-detail observable is used from which $\bar{b} - (b_\nu^{\text{SM}} + b_\nu^{\text{BSM}})$ will be extracted [1].)

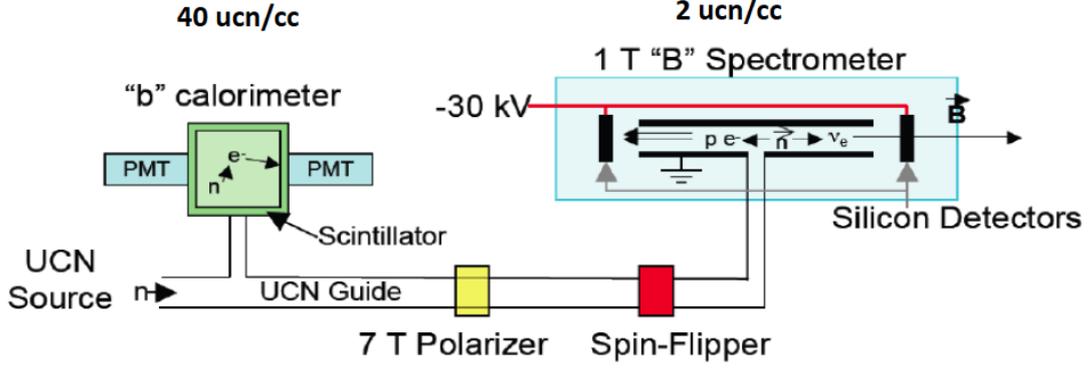


Figure 1: The experimental set up for UCNb (left panel) and UCNB (right panel). The UCNb experiment is a calorimeter that measures the spectrum of the electron. In the UCNB experiment the UCN enter the spectrometer after being polarized to over 99.5%. The asymmetry between the momentum direction of the neutrino and the spin of the decaying neutron is measured indirectly by detecting coincidences between the proton and the electron. The analysis is more complex and what one extracts in the end is $\bar{b} - (b_V^{\text{SM}} + b_V^{\text{BSM}})$.

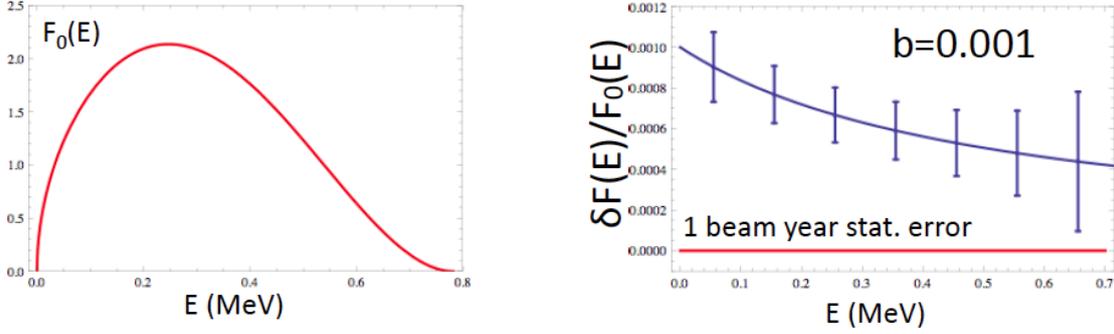


Figure 2: The left panel shows a plot of the electron spectrum $F_0(E)$. The right panel shows the expected deviations assuming $\bar{b} = 0.001$, and the statistical errors with one year of running.

4. Existing constraints

In Fig. 3, we illustrate how precision measurements of b and b_V at the 10^{-3} level can reveal the existence of new physics with mass scale Λ_i in the multi-TeV range, which will be explored directly at the LHC. Furthermore, the two panels highlight the difference in bounds between using previous phenomenological estimates for g_S and g_T [7] and current LQCD estimates [1, 2]. In these figures, b_{0^+} represents the existing constraint from $0^+ \rightarrow 0^+$ nuclear beta decays. While the LHC will be able to produce W^{BSM} with masses in the 0.2–3 TeV range and may provide precise measurements of their masses, the proposed neutron decay experiments will provide the most precise measurement of their spin and couplings, needed to reconstruct the TeV-scale theory that will replace the SM. Moreover, even if W^{BSM} is out of the kinematic reach of the LHC (couplings are $O(10^{-4})$ or smaller) or swamped by background, the interactions that it generates may lead to observable effects in the proposed neutron-decay experiments.

To connect the new theory (the theory replacing the SM at the TeV scale) to low-energy experiments requires the calculation of the matrix elements of the new interactions between the decaying

neutron and the final state proton. Even though the new theory is unknown, its consequences at the hadronic scale can be evaluated in terms of an effective theory. In this model-independent effective theory, the leading operators encapsulating novel scalar and tensor interactions are iso-vector quark bilinear operators $\bar{u}d$ and $\bar{u}\sigma^{\mu\nu}d$. Knowing their matrix elements at $q \equiv p_n - p_p = 0$, *i.e.* the scalar and tensor charges g_S and g_T , will allow us to put bounds on the couplings $\varepsilon_{S,T}$. The best method for calculating these matrix elements with control over all systematic errors is Lattice QCD. An overview of our ongoing lattice calculations is given next.

5. Lattice QCD Calculations

In Ref. [1], we show that to interpret experimental results at the 10^{-3} level requires estimating g_S and g_T with 10–20% precision. In Fig. 4, taken from Ref. [1] and discussed further in the talk by Bhattacharya [3], we show the allowed region in the $\varepsilon_{S,T}$ plane as a function of the uncertainty in g_S and g_T assuming an experimental accuracy of 10^{-3} . Need for higher precision beyond 10–20% in g_S and g_T will require higher precision in experiments. Thus, the anticipated experimental precision sets the goal of our LQCD calculations for the 2013-2014 timeframe.

As discussed further in the talk by Lin [2], 10–20% precision in g_S and g_T requires control over the extrapolations in quark mass to the physical pion mass and in lattice spacing to the continuum limit, and non-perturbative determination of the renormalization constants of the operators. To control these extrapolations requires high statistics data at three or more values of the lattice spacing and light quark masses. To satisfy these requirements within the time frame of the experiments, three years, motivated us to use the $2 + 1 + 1$ flavor dynamical gauge configurations generated by

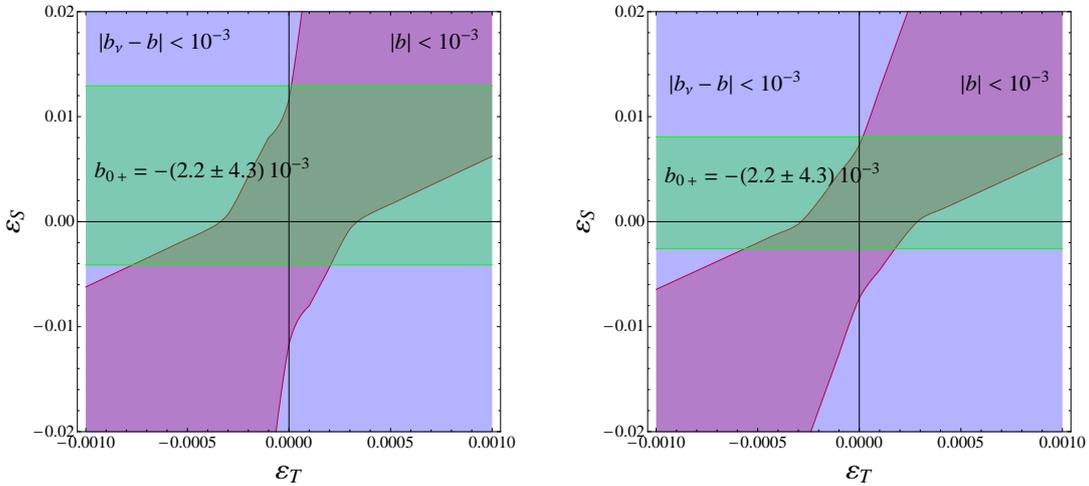


Figure 3: Left panel: 90% Confidence Level (C.L.) allowed regions in the ε_S - ε_T plane implied by (i) the existing bound on b_{0^+} characterizing the $0^+ \rightarrow 0^+$ nuclear beta decays (green horizontal band); (ii) projected measurements of \bar{b} and $(b_V^{\text{SM}} + b_V^{\text{BSM}}) - \bar{b}$ in neutron decay (red bow-tie shapes and blue regions) at the 10^{-3} level; (iii) hadronic matrix elements taken in the ranges $0.25 < g_S < 1.0$, $0.6 < g_T < 2.3$ [7]. Right panel: same as left panel but with scalar and tensor charges taken from Lattice QCD: $g_S = 0.8(4)$ and $g_T = 1.05(35)$ [1]. Note that the constraint from b_{0^+} also improves with higher precision in g_S . The effective couplings $\varepsilon_{S,T}$ are defined in the $\overline{\text{MS}}$ scheme at 2 GeV.

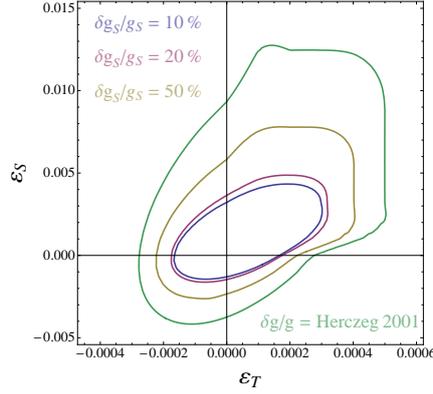


Figure 4: Constraints over the allowed regions in the $\varepsilon_S, \varepsilon_T$ plane as a function of the uncertainty in g_S and g_T for different values of $\delta g_S/g_S$ assuming $\delta g_T = 2/3\delta g_S$. The outer contour is based on the uncertainties in g_S and g_T used in the phenomenological analysis by Herczeg [7]. The convergence of the contours show that for an experimental uncertainty of 10^{-3} , the goal of LQCD calculations should be to calculate g_S and g_T to 10-20% accuracy.

the MILC Collaboration using the HISQ action. At the end of 2011, nine ensembles, each of about 1000 lattices, at the three lattice spacings 0.12, 0.09 and 0.06 fm and each with the light quark mass corresponding to a pion mass of 310, 220 and 140 MeV had already been generated. Furthermore, the generation of these lattices is continuing and supported by the USQCD national program.

For valence quarks we chose to use the improved clover fermion action since the signal in baryon correlators with staggered fermions is poor. Using clover valence quarks on HISQ lattices, however, introduces the well-known problem of exceptional configurations. These are configurations in which the clover Dirac operator has near zero modes. When the valence and sea quark actions are the same, such configurations are suppressed by a concomitant vanishing of their weight in the path integral. Unchecked, exceptional configurations would make anomalously large contributions and degrade the reliability of the ensemble average. In our calculations with clover-on-HISQ, such a cancellation is not guaranteed and checks are required for each ensemble. Our tests, to date, on ensembles with pion masses of 310 and 220 MeV, do not show evidence of exceptional configurations. Preliminary simulations at 140 MeV suggest that this fortuitous circumstance will most likely not hold at the physical pion mass, so we have restricted our calculations to quark masses with 310 and 220 MeV pions. The analysis of the data presented by Lin [2] shows that data at 310 and 220 MeV already reduces the uncertainty due to the chiral extrapolation very significantly. Nevertheless, in light of the problem with exceptional configurations, our long-term strategy for simulations at the physical quark mass is to perform clover-on-clover calculations.

Another potential source of systematic errors in LQCD calculations is guaranteeing that the matrix elements are calculated between ground states of the proton and the neutron. Mixing with excited states arises because lattice operators automatically couple to the ground state and all its excited states with the same quantum numbers. There are two ways to reduce excited state contamination—design better interpolating operators to enhance overlap with ground state, and increase the separation between the source and sink. We are addressing this problem by using

smearred sources, doing the calculation with multiple source-sink separations and by explicitly including possible contributions of excited states in our analysis.

Our existing calculations on the 0.12 and 0.09 fm lattices, with pion masses of 310 and 220 MeV and 400–500 configurations, show that the statistical errors in g_S are 5–6 times those in g_T and g_A . The uncertainty in the latter is already at the 5–10% level. These preliminary studies also show that there is very small dependence of the estimates on the lattice spacing. This gives us confidence that a reliable continuum extrapolation will be possible using data obtained at 0.12, 0.09 and 0.06 fm. We are therefore confident that analyzing the full set of 1000 lattices in each ensemble will yield estimates with 10–20% accuracy in g_S and g_T . We are also on schedule to reach this milestone over the next two years, in time when the first experimental results are expected.

Acknowledgments

We thank our collaborators in the theory effort (V. Cirigliano, A. Filipuzzi, M. Gonzalez-Alonso and M. Graesser) and our experimental colleagues. We thank the MILC Collaboration for sharing the HISQ lattices. These calculations were performed using the Chroma software suite [8]. Numerical simulations were carried out in part on facilities of the USQCD Collaboration, which are funded by the Office of Science of the U.S. Department of Energy, and the Extreme Science and Engineering Discovery Environment (XSEDE), which is supported by National Science Foundation grant number OCI-1053575. The speaker is supported by the DOE grant DE-KA-1401020.

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